



# $ar{B} ightarrow X_u \, I \, ar{ u}$ theory

#### Gil Paz

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### Charge

"We would like to invite you to give a review talk on the  $B->Xu \ l v$  theory, discussing (very quickly) the available methods, [and] the (very little) recent progress and the open problems."

## Outline

- Reminders
- Available methods
- Recent progress
- Open problems and future progress

• For *B* decays:  $5 \text{ GeV} \sim m_b \gg \Lambda_{\text{QCD}} \sim 0.5 \text{ GeV}$ Observables expandable in  $\Lambda_{\text{QCD}}/m_b \sim 0.1$ 

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- If we could measure total  $\Gamma(\bar{B} \to X_u \, | \, \bar{\nu})$  we could use a local OPE

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 $c_n$  perturbative,  $\langle O_n \rangle$  non-perturbative numbers

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$$\begin{split} M_X^2 &\sim m_b^2 & \text{local OPE} & (\text{``OPE region''}) \\ M_X^2 &\sim m_b \Lambda_{\text{QCD}} & \text{Non local OPE} & (\text{``end point region''}) \\ M_X^2 &\sim \Lambda_{\text{QCD}}^2 & \text{No inclusive description} & (\text{``resonance region''}) \end{split}$$

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- At subleading power in  $\Lambda_{QCD}/m_b$ :
- Several subleading shape functions (SSF) appear
- Different linear combinations for  $ar{B} o X_u \, I \, ar{
  u}$  and  $ar{B} o X_s \, \gamma$
- $\bar{B} \rightarrow X_s \gamma$  has unique SSF ("resolved photon contributions")
- Shape functions moments are related to HQET parameters: E.g. leading shape function:  $1^{st}$  moment  $\leftrightarrow m_b$ ,  $2^{nd}$  moment  $\leftrightarrow \mu_{\pi}^2$

• BLL use  $q^2 - m_X$  cut to reduce shape function dependance [Bauer, Ligeti, Luke, PRD **64**, 113004, (2001)]

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- Methods with flexible cuts must include shape functions effects:
- BLNP

[Lange, Neubert, GP, PRD 72, 073006, (2005)]

- GGOU

[Gambino, Giordano, Ossola, Uraltsev, JHEP 10, 058, (2007)]

- DGE

[Andersen, Gardi, JHEP 01, 097, (2006)]

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• In the following only discuss BLNP, GGOU, DGE

#### BLNP

- Based on  $d\Gamma \sim H \cdot J \otimes S + \frac{1}{m_b} \sum_i H \cdot J \otimes s_i + ...$
- Factorize perturbative coefficient into hard  ${\sf H}$  and jet  ${\sf J}$  function
- Leading power  $H \cdot J \otimes S$  at  $\mathcal{O}(\alpha_s)$
- Subleading shape functions:  $H \cdot J \otimes s_i$  at  $\mathcal{O}(\alpha_s^0)$
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- Many BLNP NNLO calculations are known:
   H, J at O(α<sub>s</sub><sup>2</sup>), j<sub>i</sub>/m<sub>b</sub> at O(α<sub>s</sub>), resolved photon contributions Not fully combined yet

# GGOU

Based on

$$W_i \sim F_i \otimes W_i^{pert}$$

- $W_i$  structure functions that give  $d\Gamma$
- $W_i^{pert}$  known perturbative quantities
- $F_i(k_+, q^2, \mu)$  OPE-constrained non-perturbative distribution functions
- uses kinetic scheme, Wilsonian cutoff  $\mu \sim 1~{
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- Improvements to GGOU discussed later

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- $m_b(\overline{MS})$  used as an input
- See also Gardi's talk at MITP 2015 workshop

# Available methods: $|V_{ub}|$

- Latest HFLAV summary [arXiv:1612.07233]
- Table 91: Summary of inclusive determinations of  $|V_{ub}|$ The errors quoted on  $|V_{ub}|$  correspond to experimental and theoretical uncertainties

Framework	$ V_{ub} [10^{-3}]$
BLNP	$4.44 \pm 0.15^{+0.21}_{-0.22}$
DGE	$4.52 \pm 0.16 ^{+0.15}_{-0.16}$
GGOU	$4.52 \pm 0.15 \substack{+0.11 \\ -0.14}$
ADFR	$4.08 \pm 0.13^{+0.18}_{-0.12}$
BLL $(m_X/q^2 \text{ only})$	$4.62 \pm 0.20 \pm 0.29$

• This doesn't include latest BaBar analysis [BaBar PRD **95** 072001 (2017)]

# Recent progress

• GGOU: neural network approach

[Gambino, Healey, Mondino PRD 94 014031 (2016)]

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[Hashimoto, Colquhoun, Izubuchi, Kaneko, Ohki, EPJ Web Conf. **175** 13006 (2018)

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I apologize if I missed other progress

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#### GGOU: neural network approach

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- See also Gambino's talk later today
- Motivation: About 100 forms considered in GGOU but each parameterized by simple 2-parameter functional forms is that good enough?



• Use Neural Networks to parameterize shape functions without bias Extract  $|V_{ub}|$  from theoretical constraints and data Similar to NNPDF

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• Selection of NN replicas of  $F_2(k_+, 0)$  trained on first 3 moments only



Demonstrates NN capability to properly sample the functional space

• After further pruning, e.g. keep only one dominant peak



- The results are used to extract  $|V_{ub}|$  in the GGOU framework
- Good agreement is found with original GGOU and 2014 HFLAV

# Recent progress: Power corrections

[Gambino, Healey, Turczyk PLB 763 60 (2016)]

#### Power corrections

- Inclusive  $\bar{B} \rightarrow X_c \ell \bar{\nu}_\ell$  decays allow to extract HQET parameters
- Moments of shape function(s) are related to these parameters
- Dimension 7 and 8 HQET operators contribution to  $\bar{B} \rightarrow X_c \ell \bar{\nu}_\ell$ [Mannel, Turczyk, Uraltsev JHEP **1011**, 109 (2010)]
- |V<sub>cb</sub>| extraction from inclusive B decays uses dimension 7 and 8 HQET operators [Gambino, Healey, Turczyk PLB 763, 60 (2016)]

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#### Table 2

Default fit results: the second and third columns give the central values and standard deviations.

$m_{b}^{kin}$	4.546	0.021	$r_1$	0.032	0.024
$\overline{m}_c$ (3 GeV)	0.987	0.013	$r_2$	-0.063	0.037
$\mu_{\pi}^2$	0.432	0.068	$r_3$	-0.017	0.025
$\mu_G^2$	0.355	0.060	$r_4$	-0.002	0.025
$\rho_D^3$	0.145	0.061	$r_5$	0.001	0.025
$\rho_{LS}^3$	-0.169	0.097	$r_6$	0.016	0.025
$\overline{m}_1$	0.084	0.059	$r_7$	0.002	0.025
$\overline{m}_2$	-0.019	0.036	$r_8$	-0.026	0.025
$\overline{m}_3$	-0.011	0.045	$r_9$	0.072	0.044
$\overline{m}_4$	0.048	0.043	r <sub>10</sub>	0.043	0.030
$\overline{m}_5$	0.072	0.045	$r_{11}$	0.003	0.025
$\overline{m}_6$	0.015	0.041	r <sub>12</sub>	0.018	0.025
$\overline{m}_7$	-0.059	0.043	r <sub>13</sub>	-0.052	0.031
$\overline{m}_8$	-0.178	0.073	$r_{14}$	0.003	0.025
$\overline{m}_9$	-0.035	0.044	$r_{15}$	0.001	0.025
$\chi^2/dof$	0.46		$r_{16}$	0.001	0.025
BR(%)	10.652	0.156	$r_{17}$	-0.028	0.025
10 <sup>3</sup>  V <sub>cb</sub>	42.11	0.74	$r_{18}$	-0.001	0.025

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[Gunawardana, GP, JHEP 1707 137 (2017)]

#### Motivation

Original motivation: How to express moments of shape function(s) in terms of m<sub>1</sub>,...m<sub>9</sub> and r<sub>1</sub>,...r<sub>18</sub>? (See also appendix A of [Heinonen, Mannel, arXiv:1609.01334])

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- Method of [Gunawardana, GP, JHEP 1707 137 (2017)] allows to
- 1) Find such relations
- 2) List HQET parameters, in principle, to arbitrary dimension
- Construct NRQED and NRQCD bilinear operators, in principle, to *arbitrary* dimension

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- 1) Find such relations
- 2) List HQET parameters, in principle, to arbitrary dimension
- 3) Construct NRQED and NRQCD bilinear operators, in principle, to *arbitrary* dimension
  - Structure of effective field theories is simpler than we think:
     SM EFT

[Henning, Lu, Melia, Murayama, JHEP 1708, 016 (2017)]

- NRQED/NRQCD/HQET

[Gunawardna, GP JHEP **1707** 137 (2017)] [Kobach, Pal PLB **772** 225 (2017)]

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- Paz et al. arXiv:1702.0890 v1  $\rightarrow$  v2

"discussion of operators with multiple color structures was added"

• [Gunawardana, GP, JHEP **1707** 137 (2017)] method Consider matrix elements of the form  $\langle H|\bar{h} iD^{\mu_1} \dots iD^{\mu_n}h|H \rangle$  $\langle H|\bar{h} iD^{\mu_1} \dots iD^{\mu_n}s^{\lambda}h|H \rangle$ 

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- Orthogonality:  $v_{\mu_1}=v_{\mu_n}=v_\lambda=0$
- P,T, and Hermitian conjugation:
   SI (SD) matrix elements are sym. (anti-sym.) under inversion
- Four dimensions:

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- Published and  $v \ge 2$  of [Gunawardana, GP, JHEP 1707 137 (2017)]
- Checking possible multiple color structures

### New Result: Dimension 9 HQET operators

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New Result: Dimension 9 HQET operators • Using the general method: SI Dimension 9 HQET operators  $\frac{1}{2M_{H}}\langle H|\bar{h}\,iD^{\mu_{1}}\,iD^{\mu_{2}}\,iD^{\mu_{3}}\,iD^{\mu_{4}}\,iD^{\mu_{5}}\,iD^{\mu_{6}}\,h|H\rangle = a^{(9)}_{12,34}\,\Pi^{\mu_{1}\mu_{2}}\,\Pi^{\mu_{3}\mu_{4}}\,\Pi^{\mu_{5}\mu_{6}} + a^{(9)}_{12,34}\,\Pi^{\mu_{1}\mu_{2}}\,\Pi^{\mu_{1}\mu_{5}}\,\Pi^{\mu_{5}\mu_{6}} + a^{(9)}_{12,34}\,\Pi^{\mu_{5}\mu_{6}}\,\Pi^{\mu_{5}\mu_{6}} + a^{(9)}_{12,34}\,\Pi^{\mu_{1}\mu_{2}}\,\Pi^{\mu_{5}\mu_{6}}\,\Pi^{\mu_{5}\mu_{6}} + a^{(9)}_{12,34}\,\Pi^{\mu_{5}\mu_{6}}\,\Pi^{\mu_{5}\mu_{6}}\,\Pi^{\mu_{5}\mu_{6}} + a^{(9)}_{12,34}\,\Pi^{\mu_{5}\mu_{6}}\,\Pi^{\mu_{5}\mu_{6}}\,\Pi^{\mu_{5}\mu_{6}} + a^{(9)}_{12,34}\,\Pi^{\mu_{5}\mu_{6}}\,\Pi^{\mu_{5}$  $+a_{12,35}^{(9)}\left(\Pi^{\mu_1\mu_2}\Pi^{\mu_3\mu_5}\Pi^{\mu_4\mu_6}+\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_4}\Pi^{\mu_5\mu_6}\right)+a_{12,36}^{(9)}\left(\Pi^{\mu_1\mu_2}\Pi^{\mu_3\mu_6}\Pi^{\mu_4\mu_5}+\Pi^{\mu_1\mu_4}\Pi^{\mu_2\mu_3}\Pi^{\mu_5\mu_6}\right)+$  $+a_{13,25}^{(9)}\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_5}\Pi^{\mu_4\mu_6}+a_{13,26}^{(9)}(\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_5}+\Pi^{\mu_1\mu_5}\Pi^{\mu_2\mu_3}\Pi^{\mu_4\mu_6})+a_{14,25}^{(9)}\Pi^{\mu_1\mu_4}\Pi^{\mu_2\mu_5}\Pi^{\mu_3\mu_6}+A_{13,26}^{(9)}(\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_5}+\Pi^{\mu_1\mu_5}\Pi^{\mu_2\mu_5}\Pi^{\mu_4\mu_6})+a_{14,25}^{(9)}\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_5}\Pi^{\mu_4\mu_6}+A_{13,26}^{(9)}(\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_5}+\Pi^{\mu_1\mu_5}\Pi^{\mu_2\mu_5}\Pi^{\mu_4\mu_6})+a_{14,25}^{(9)}\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_5}\Pi^{\mu_4\mu_6}+A_{13,26}^{(9)}(\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_5}+\Pi^{\mu_1\mu_5}\Pi^{\mu_2\mu_5}\Pi^{\mu_4\mu_6})+A_{14,25}^{(9)}(\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_5}+\Pi^{\mu_1\mu_5}\Pi^{\mu_2\mu_5}\Pi^{\mu_4\mu_6})+A_{14,25}^{(9)}(\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_5}+\Pi^{\mu_1\mu_5}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_6})+A_{14,25}^{(9)}(\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_6})+A_{14,25}^{(9)}(\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_6})+A_{14,25}^{(9)}(\Pi^{\mu_1\mu_4}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_6})+A_{14,25}^{(9)}(\Pi^{\mu_1\mu_4}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_6})+A_{14,25}^{(9)}(\Pi^{\mu_1\mu_4}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_6})+A_{14,25}^{(9)}(\Pi^{\mu_1\mu_4}\Pi^{\mu_2\mu_6}\Pi^{\mu_4\mu_6})+A_{14,25}^{(9)}(\Pi^{\mu_1\mu_4}\Pi^{\mu_4\mu_6})+A_{14,25}^{(9)}(\Pi^{\mu_1\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4}\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4})+A_{14,25}^{(9)}(\Pi^{\mu_4})+A_$  $+a_{14.26}^{(9)}\left(\Pi^{\mu_{1}\mu_{4}}\Pi^{\mu_{2}\mu_{6}}\Pi^{\mu_{3}\mu_{5}}+\Pi^{\mu_{1}\mu_{5}}\Pi^{\mu_{2}\mu_{4}}\Pi^{\mu_{3}\mu_{6}}\right)+a_{15.26}^{(9)}\Pi^{\mu_{1}\mu_{5}}\Pi^{\mu_{2}\mu_{6}}\Pi^{\mu_{3}\mu_{4}}+a_{16.23}^{(9)}\Pi^{\mu_{1}\mu_{6}}\Pi^{\mu_{2}\mu_{3}}\Pi^{\mu_{4}\mu_{5}}+$  $+a_{16,24}^{(9)}\Pi^{\mu_1\mu_6}\Pi^{\mu_2\mu_4}\Pi^{\mu_3\mu_5}+a_{16,25}^{(9)}\Pi^{\mu_1\mu_6}\Pi^{\mu_2\mu_5}\Pi^{\mu_3\mu_4}+b_{12,36}^{(9)}\left(\Pi^{\mu_1\mu_2}\Pi^{\mu_3\mu_6}v^{\mu_4}v^{\mu_5}+\Pi^{\mu_1\mu_4}\Pi^{\mu_5\mu_6}v^{\mu_2}v^{\mu_3}\right)+$  $+b^{(9)}_{12,46} \left(\Pi^{\mu_1\mu_2}\Pi^{\mu_4\mu_6}v^{\mu_3}v^{\mu_5} + \Pi^{\mu_1\mu_3}\Pi^{\mu_5\mu_6}v^{\mu_2}v^{\mu_4}\right) + b^{(9)}_{12,56} \Pi^{\mu_1\mu_2}\Pi^{\mu_5\mu_6}v^{\mu_3}v^{\mu_4} +$  $+b^{(9)}_{13,26} \left(\Pi^{\mu_1\mu_3}\Pi^{\mu_2\mu_6}v^{\mu_4}v^{\mu_5}+\Pi^{\mu_1\mu_5}\Pi^{\mu_4\mu_6}v^{\mu_2}v^{\mu_3}\right)+b^{(9)}_{13,46}\Pi^{\mu_1\mu_3}\Pi^{\mu_4\mu_6}v^{\mu_2}v^{\mu_5}+$  $+b_{14,26}^{(9)}\left(\Pi^{\mu_1\mu_4}\Pi^{\mu_2\mu_6}v^{\mu_3}v^{\mu_5}+\Pi^{\mu_1\mu_5}\Pi^{\mu_3\mu_6}v^{\mu_2}v^{\mu_4}\right)+b_{14,36}^{(9)}\Pi^{\mu_1\mu_4}\Pi^{\mu_3\mu_6}v^{\mu_2}v^{\mu_5}+b_{15,26}^{(9)}\Pi^{\mu_1\mu_5}\Pi^{\mu_2\mu_6}v^{\mu_3}v^{\mu_4}+b_{14,36}^{(9)}\Pi^{\mu_1\mu_4}\Pi^{\mu_2\mu_6}v^{\mu_3}v^{\mu_5}+\Pi^{\mu_1\mu_5}\Pi^{\mu_3\mu_6}v^{\mu_2}v^{\mu_4}\right)+b_{14,36}^{(9)}\Pi^{\mu_1\mu_4}\Pi^{\mu_3\mu_6}v^{\mu_2}v^{\mu_5}+h_{15,26}^{(9)}\Pi^{\mu_1\mu_5}\Pi^{\mu_2\mu_6}v^{\mu_4}v^{\mu_4}+h_{14,36}^{(9)}\Pi^{\mu_1\mu_5}\Pi^{\mu_2\mu_6}v^{\mu_5}v^{\mu_5}+h_{14,36}^{(9)}\eta^{\mu_1\mu_5}\Pi^{\mu_2\mu_6}v^{\mu_5}v^{\mu_5}+h_{14,36}^{(9)}\eta^{\mu_1\mu_5}\Pi^{\mu_2\mu_6}v^{\mu_5}v^{\mu_5}+h_{14,36}^{(9)}\eta^{\mu_5}v^{\mu_5}+h_{14,36}^{(9)}\eta^{\mu_5}v^{\mu_5}v^{\mu_5}+h_{14,36}^{(9)}\eta^{\mu_5}v^{\mu_5}+h_{14,36}^{(9)}v^{\mu_5}v^{\mu_5}+h_{14,36}^{(9)}\eta^{\mu_5}v^{\mu_5}+h_{14,36}^{(9)}\eta^{\mu_5}v^{\mu_5}+h_{14,36}^{(9)}v^{\mu$  $+b_{16,23}^{(9)}\left(\Pi^{\mu_{1}\mu_{6}}\Pi^{\mu_{2}\mu_{3}}v^{\mu_{4}}v^{\mu_{5}}+\Pi^{\mu_{1}\mu_{6}}\Pi^{\mu_{4}\mu_{5}}v^{\mu_{2}}v^{\mu_{3}}\right)+b_{16,24}^{(9)}\left(\Pi^{\mu_{1}\mu_{6}}\Pi^{\mu_{2}\mu_{4}}v^{\mu_{3}}v^{\mu_{5}}+\Pi^{\mu_{1}\mu_{6}}\Pi^{\mu_{3}\mu_{5}}v^{\mu_{2}}v^{\mu_{4}}\right)+$  $+b_{16,25}^{(9)}\,\Pi^{\mu_{1}\mu_{6}}\Pi^{\mu_{2}\mu_{5}}v^{\mu_{3}}v^{\mu_{4}}+b_{16,34}^{(9)}\,\Pi^{\mu_{1}\mu_{6}}\Pi^{\mu_{3}\mu_{4}}v^{\mu_{2}}v^{\mu_{5}}+c^{(9)}\,\Pi^{\mu_{1}\mu_{6}}v^{\mu_{2}}v^{\mu_{3}}v^{\mu_{4}}v^{\mu_{5}}$ 

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- Multiple color structures arise from combining pure color octets: [*iD<sup>μ<sub>i</sub></sup>*, *iD<sup>μ<sub>j</sub>*], [*iD<sup>μ<sub>i</sub></sup>*, [*iD<sup>μ<sub>j</sub></sup>*, *iD<sup>μ<sub>k</sub>*]], [*iD<sup>μ<sub>i</sub></sup>*, [*iD<sup>μ<sub>k</sub></sup>*, *iD<sup>μ<sub>l</sub></sup>*]]]
  </sup></sup>
- For  $|V_{ub}|$  and  $|V_{cb}|$  with tree level dimension  $\geq 7$  power corrections Only  $T^aT^b$  color structure is needed

# New Result: Moments of the leading power shape function

• Moments of the shape function are related to HQET parameters The matrix elements decomposition makes their calculation easy

$$2M_B \int d\omega \, \omega^k \, S(\omega) = n_{\mu_1} ... n_{\mu_k} \langle \bar{B}(v) | \bar{h} \, i D^{\mu_1} ... i D^{\mu_k} \, h | \bar{B}(v) \rangle$$

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$$\int d\omega \, S(\omega) = 1, \qquad \int d\omega \, \omega \, S(\omega) = 0, \qquad \int d\omega \, \omega^2 \, S(\omega) = -a^{(5)} = -\lambda_1/3,$$

$$\int d\omega \, \omega^3 \, S(\omega) = -a^{(6)} = -\rho_1/3,$$

$$\int d\omega \, \omega^4 \, S(\omega) = a_{12}^{(7)} + a_{13}^{(7)} + a_{14}^{(7)} - b^{(7)} = m_1/5 - m_2/3,$$

$$\int d\omega \, \omega^5 \, S(\omega) = 2a_{12}^{(8)} + 2a_{13}^{(8)} + 2a_{15}^{(8)} + b_{12}^{(8)} + b_{14}^{(8)} + b_{15}^{(8)} - c^{(8)} =$$

$$= (-8r_1 + 2r_2 + 2r_3 + 2r_4 + r_5 + r_6 + r_7) / 15,$$

$$\int d\omega \, \omega^6 \, S(\omega) = -a_{12,34}^{(9)} - 2a_{12,35}^{(9)} - 2a_{12,36}^{(9)} - a_{13,25}^{(9)} - 2a_{13,26}^{(9)} - a_{14,25}^{(9)} - 2a_{14,26}^{(9)} - a_{15,26}^{(9)} + a_{16,23}^{(9)} - a_{16,23}^{(9)} - a_{16,24}^{(9)} - a_{16,25}^{(9)} + 2b_{12,36}^{(9)} + 2b_{12,36}^{(9)} + b_{12,56}^{(9)} + 2b_{13,46}^{(9)} + 2b_{14,26}^{(9)} + b_{13,46}^{(9)} + 2b_{14,26}^{(9)} + b_{14,36}^{(9)} + b_{15,26}^{(9)} + 2b_{16,23}^{(9)} + 2b_{16,24}^{(9)} + b_{16,34}^{(9)} - c^{(9)}$$

• Future: moments of other SSF [Gunawardana, GP, in progress)]

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The future looks promising for  $\bar{B} \rightarrow X_u \, I \, \bar{\nu}$  and inclusive  $|V_{ub}|!$