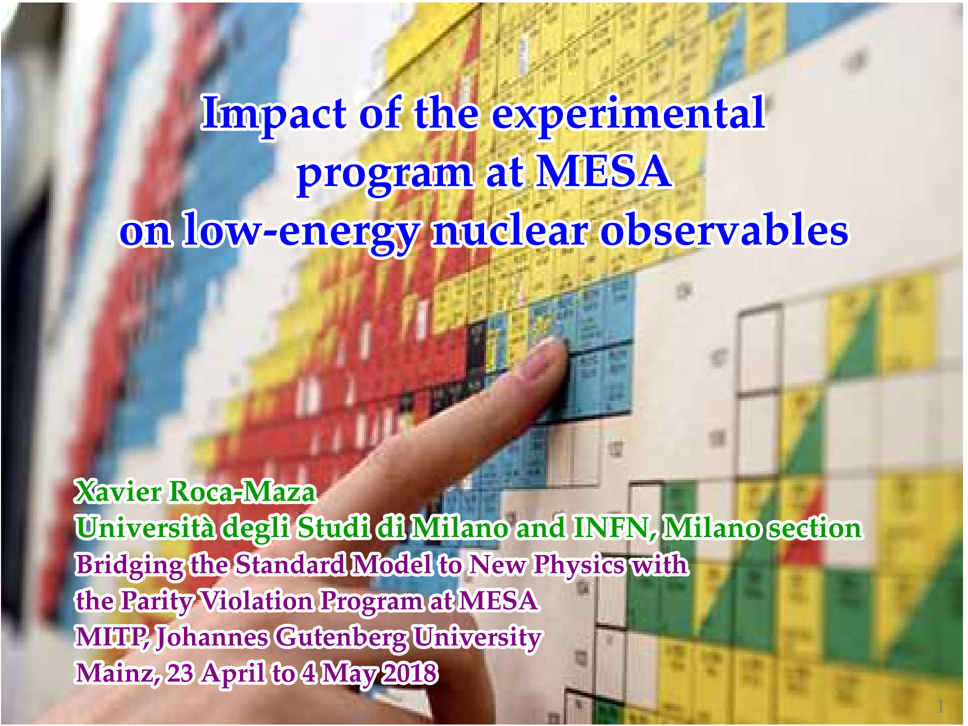

Impact of the experimental program at MESA on low-energy nuclear observables

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Bridging the Standard Model to New Physics with the Parity Violation Program at MESA
MITP, Johannes Gutenberg University
Mainz, 23 April to 4 May 2018

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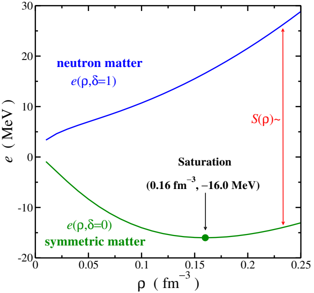
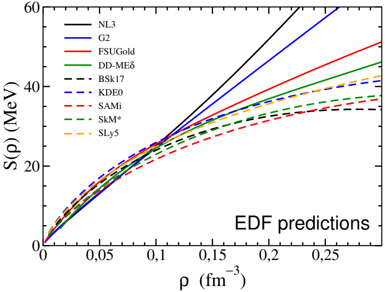
The Nuclear Equation of State: Infinite System





- Expansion for small asymmetries $\delta = \dfrac{\rho_n - \rho_p}{\rho_n + \rho_p}$:

$$e(\rho, \delta) = e(\rho, \delta = 0) + S(\rho)\delta^2 + \mathcal{O}[\delta^4]$$

- Expansion on the density around saturation $x \equiv \dfrac{\rho - \rho_0}{3\rho}$:

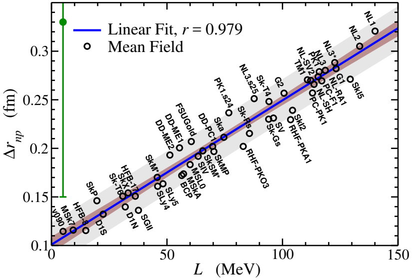
$$e(\rho, \delta) = \left[e(\rho_0, \delta = 0) + \frac{1}{2}K_0 x^2\right] + \left[J + Lx + \frac{1}{2}K_{sym} x^2\right]\delta^2 + \mathcal{O}[\delta^2, \rho^3]$$

Uncertainties on $S(\rho)$ around saturation (mainly due to L) **impact** on many nuclear physics and astrophysics **observables**.

Example: L and the neutron skin in ^{208}Pb

$$\Delta r_{np} \equiv \langle r_n^2 \rangle^{1/2} - \langle r_p^2 \rangle^{1/2}$$

$$\text{Macroscopic model: } \Delta r_{np} \sim \frac{1}{12} \frac{(N-Z)}{A} \frac{R}{J} L \qquad (L \propto p_0^{neut})$$

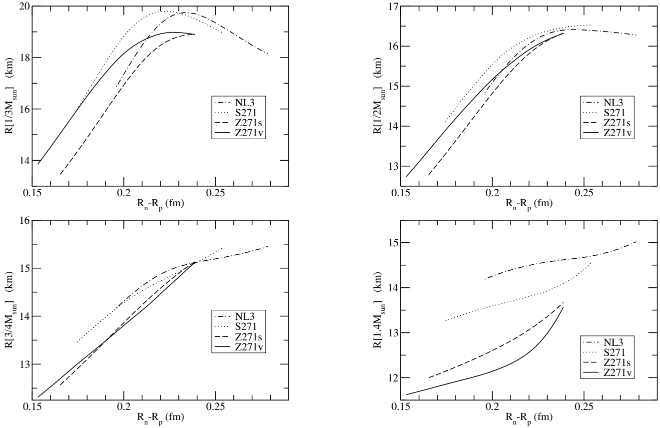


Physical Review Letters **106**, 252501 (2011)

The faster the symmetry energy increases with density (L), the largest the size of the neutron skin in (heavy) nuclei.

[Exp. from strongly interacting probes: $\sim 0.15 - 0.22$ fm (*Physical Review C* 86 *015803 (2012)*)].

Example: L (or Δr_{np}) and the radius of a neutron star



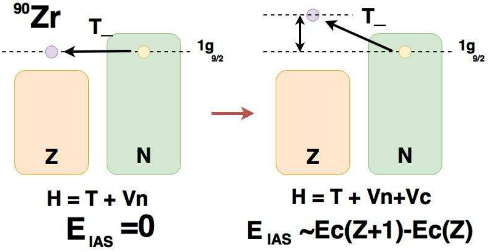
The Astrophysical Journal, 593 (2003) 463

Measurement of Δr_{np} $\rightarrow$ constraint **low-mass neutron star R**.
Why? 0.5 $M_{\odot}$ have central densities near normal saturation density ρ_0.

Are there other cases in which a measurment of Δr_{np} can improve our understanding of the nuclear EoS?

(Not exhaustive, other cases have been previously discussed)

The isobaric analog state energy: E_{IAS}



- **Analog state** can be defined: $|A\rangle = \dfrac{T_-|0\rangle}{\langle 0|T_+T_-|0\rangle}$
- **Displacement energy or** E_{IAS}

$$E_{IAS} = E_A - E_0 = \langle A|\mathcal{H}|A\rangle - \langle 0|\mathcal{H}|0\rangle = \frac{\langle 0|[T_+[\mathcal{H}, T_-]|0\rangle}{\langle 0|T_+T_-|0\rangle}$$

$E_{IAS} \neq 0$ **only due to Isospin Symmtry Breaking terms** $\mathcal{H}$
E_{IAS}^{exp} **usually accuratelly measured !**

Contributions

$[\mathcal{H}, T_-] \neq 0$? essentially **Coulomb potential** but not only

Table: Estimate of the different effects on E_{IAS} in ^{208}Pb.

	E_{IAS} Correction
Coumb direct	$\sim$ 20 MeV
Coulomb exchange	$\sim$ -300 keV
n-p mass difference	$\sim$ tens keV
Electromagnetic spin-orbit	$\sim$ - tens keV
Finite size effects	$\sim$ - 100 keV
Short range correlations	$\sim$ 100 keV
sospin impurity	$\sim$ -100 keV
Isospin symmetry breaking	$\sim$ - 250 keV
	$\sim$ 19 MeV

Physical Review Letters **23**, 484 (1969).

$E_{IAS}^{exp} = 18.83 \pm 0.01$ MeV. *Nuclear Data Sheets 108, 1583 (2007).*

Coulomb direct contribution: very simple model

• Assuming indepentent particle model and good isospin for $|0\rangle$ ($\langle 0|T_+T_-|0\rangle = 2T_0 = N - Z$)

$$E_{IAS} \approx E_{IAS}^{C,direct} = \frac{1}{N-Z}\int [\rho_n(\vec{r}) - \rho_p(\vec{r})]\, U_C^{direct}(\vec{r})d\vec{r}$$

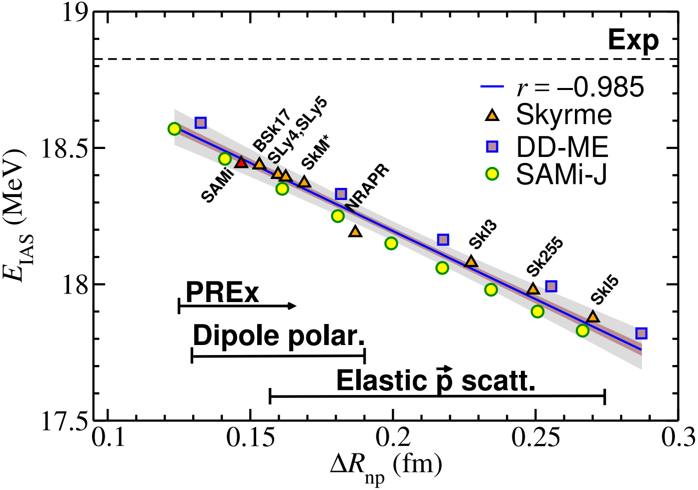
where $U_C^{direct}(\vec{r}) = \int \frac{e^2}{|\vec{r}_1 - \vec{r}|}\rho_{ch}(\vec{r}_1)d\vec{r}_1$

• Assuming also a uniform neutron and proton distributions of radius R_n and R_p respectively, and $\rho_{ch} \approx \rho_p$ one can find

$$E_{IAS} \approx E_{IAS}^{C,direct} \approx \frac{6}{5}\frac{Ze^2}{R_p}\left(1 - \sqrt{\frac{5}{12}}\frac{N}{N-Z}\frac{\Delta r_{np}}{R_p}\right)$$

One may expect: **the larger the Δr_{np} the smallest E_{IAS}**

E_{IAS} in Energy Density Functionals (No Corr.)



arXiv:1803.09120 [nucl-th]

Nuclear models (EDFs) where the nuclear part is isospin symmetric and U_{ch} is calculated from the ρ_p

Corrections: within self-consistent HF+RPA

Within the **HF+RPA** one can **estimate** the E_{IAS} accounting (in an effective way) for **short-range correlations, isospin impurities and effects of the continuum** (if a large sp base is adopted).

- **Coulomb exchange** exact (usually Slater approx.):

$$U_C^{x,\text{exact}}\varphi_i(\vec{r}) = -\frac{e^2}{2}\int d^3r' \, \frac{\varphi_j^*(\vec{r}')\varphi_j(\vec{r})}{|\vec{r}-\vec{r}'|}\varphi_i(\vec{r}')$$

- The **electromagnetic spin-orbit** correction to the nucleon single-particle energy (non-relativistic),

$$\varepsilon_i^{\text{emso}} = \frac{\hbar^2 c^2}{2m_i^2 c^4}\langle \vec{l}_i \cdot \vec{s}_i \rangle x_i \int \frac{1}{r}\frac{dU_C}{dr}|R_i(r)|^2$$

where x_i: $g_p - 1$ for Z and g_n for N; $g_n = -3.82608545(90)$ and $g_p = 5.585694702(17)$, $R_i \to R_{nl}$ radial wf.

Corrections:

- **Finite size** effects (assuming spherical symmetry):

$$\rho_{ch}(q) = \left(1 - \frac{q^2}{8m^2}\right)\left[G_{E,p}(q^2)\rho_p(q) + G_{E,n}(q^2)\rho_n(q)\right]$$
$$- \frac{\pi q^2}{2m^2}\sum_{l,t}\left[2G_{M,t}(q^2) - G_{E,t}(q^2)\right]\langle\vec{l}\cdot\vec{s}\rangle\int_0^\infty dx\frac{j_1(qx)}{qx}|R_{nl}(x)x^2|^2$$

- **Vacuum polarization:** lowest order correction in the fine-structure constant to the Coulomb potential $\frac{eZ}{r}$:

$$V_{vp}(\vec{r}) = -\frac{2}{3}\frac{\alpha e^2}{\pi}\int d\vec{r}'\frac{\rho(\vec{r}')}{|\vec{r}-\vec{r}'|}\mathcal{K}_1\left(\frac{2}{\lambda\!\!\!^-_e}|\vec{r}-\vec{r}'|\right)$$

where e is the fundamental electric charge, α the fine-structure constrant, $\lambda\!\!\!^-_e$ the reduced Compton electron wavelength and

$$\mathcal{K}_1(x) \equiv \int_1^\infty dt e^{-xt}\left(\frac{1}{t^2} + \frac{1}{2t^4}\right)\sqrt{t^2-1}$$

Corrections:

- **Isospin symmetry breaking** (Skyrme-like): **two parts** (contact interaction)

charge symmetry breaking +

$$V_{\mathrm{CSB}} = V_{nn} - V_{pp}$$

$V_{\mathrm{CSB}}(\vec{r}_1, \vec{r}_2) \equiv \frac{1}{4}\left[\tau_z(1) + \tau_z(2)\right] s_0 (1 - P_\sigma)$

τ_z Pauli in isospin space; P_σ are the usual projector operators in spin space.

charge independence breaking*

$$V_{\mathrm{CIB}} = \frac{1}{2}(V_{nn} + V_{pp}) - V_{pn}$$

$V_{\mathrm{CIB}}(\vec{r}_1, \vec{r}_2) \equiv \frac{1}{2}\tau_z(1)\tau_z(2) u_0 (1 - P_\sigma)$

* general operator form $\tau_z(1)\tau_z(2) - \frac{1}{3}\vec{\tau}(1)\cdot\vec{\tau}(2)$. Our prescription $\tau_z(1)\tau_z(2)$ not change structure of HF+RPA.

- Opposite to the other corrections, **ISB contributions depends on new parameters that need to be determined!**

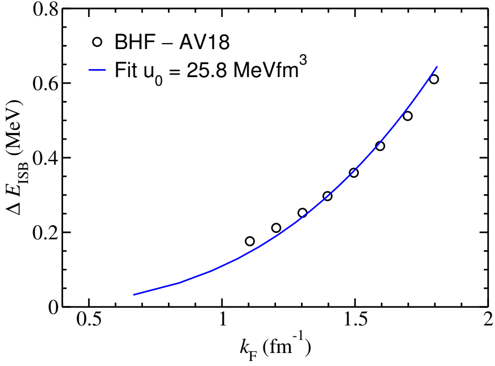
Isospin symmetry breaking in the medium:

- **keeping** things **simple**: **CSB and CIB** interaction just **delta function** depending on s_0 **and** u_0. **Different possibilities:**

→ **Fitting** to (two) experimentally known **IAS energies**

→ **Derive from theory**

→ **our option:** u_0 **to reproduce BHF** (symmetric nuclear matter) and s_0 **to reproduce** E_{IAS} **in** ^{208}**Pb**



Physics Letters B 445, 259 (1999)

Re-fit of SAMi: SAMi-ISB

• All these **corrections** are relatively **small** but **modify binding energies, neutron and proton distributions, etc.**
⇒ a **re-fit of the interaction is needed**.
• Use **SAMi fitting protocol** (special care for spin-isospin resonances) including all corrections and **find SAMi-ISB**

Table: Saturation properties

	SAMi	SAMi-ISB	
ρ_∞	0.159(1)	0.1613(6)	fm^{-3}
e_∞	−15.93(9)	−16.03(2)	MeV
m^*_{IS}	0.6752(3)	0.730(19)	
m^*_{IV}	0.664(13)	0.667(120)	
J	28(1)	30.8(4)	MeV
L	44(7)	50(4)	MeV
K_∞	245(1)	235(4)	MeV

SAMi-ISB finite nuclei properties

El.	N	B [MeV]	B^{exp} [MeV]	r_c [fm]	r_c^{exp} [fm]	ΔR_{np} [fm]
Ca	28	417.67	415.99	3.49	3.47	0.214
Zr	50	783.60	783.89	4.26	4.27	0.097
Sn	82	1102.75	1102.85	4.73	–	0.217
Pb	126	1635.78	1636.43	5.50	5.50	0.151

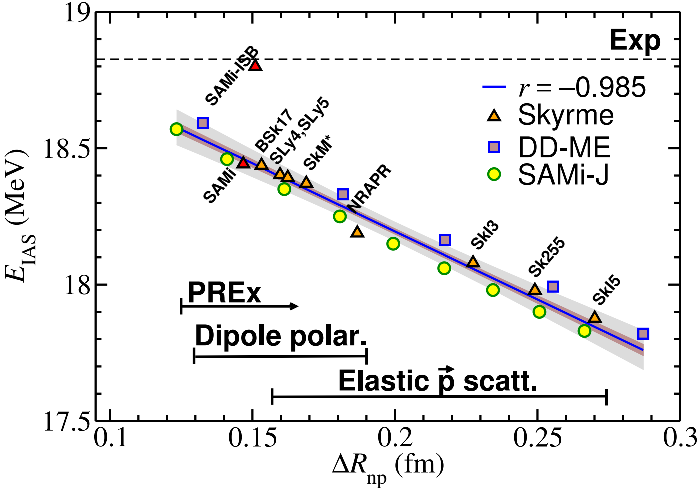
Corrections on E_{IAS} for ^{208}Pb one by one

	E_{IAS} [MeV]	Correction [keV]
No corrections[a]	18.31	
Exact Coulomb exchange	18.41	+100
n/p mass difference	18.44	+30
Electromagnetic spin-orbit	18.45	+10
Finite size effects	18.40	-50
Vacuum polarization (V_{ch})	18.53	+130
Isospin symmetry breaking	**18.80(5)**	+270

[a] From Skyrme Hamiltonian where the nuclear part is isospin symmetric and V_{ch} is calculated from the ρ_p

$E_{IAS}^{exp} = 18.83 \pm 0.01$ MeV. *Nuclear Data Sheets 108, 1583 (2007).*

E_{IAS} with SAMi-ISB



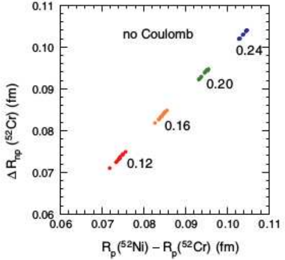
arXiv:1803.09120 [nucl-th]

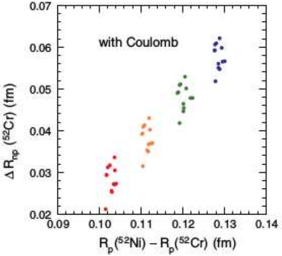
Measurement of Δr_{np} → determine ISB in the nuclear medium

Differences in the charge radii of mirror nuclei

If isospin symmetry is conserved $\Rightarrow r_n(N,Z) = r_p(Z,N)$
and, therefore,
$\Delta r_{np} \equiv r_n(N,Z) - r_p(N,Z) = r_p(Z,N) - r_p(N,Z)$

- In real nuclei, where isospin symmtry is broken:



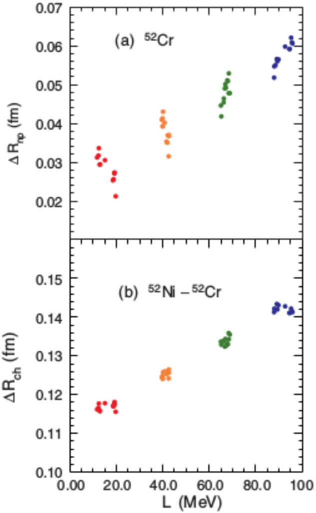


Phys. Rev. Lett. 119 122502 (2017)

Regarding Coulomb: "This is due to the self-consistent competition between the Coulomb interaction and the symmetry potential in the EDF calculations."

Differences in the charge radii of mirror nuclei

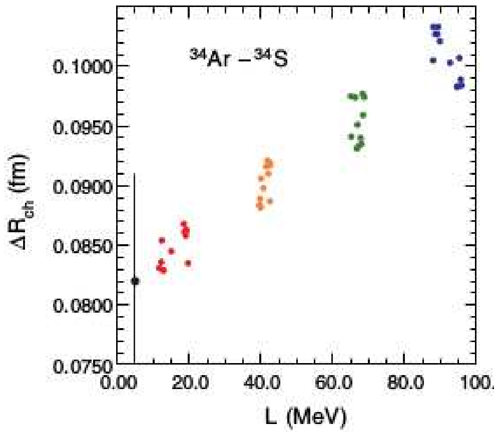
$$r_p(Z,N) - r_p(N,Z) \rightarrow r_{ch}(Z,N) - r_{ch}(N,Z)$$



- r_{ch} obtained from the **proton radii** by making corrections for the **finite charge size** of the proton and neutron and for the **relativistic effects**
- **No nuclear ISB taken into account** explicitely (No Coulomb exchange simulates nuclear ISB - SkX)
- **Why correlation improves?** (Needs to be further investigated)
- **Do other models agree?** [seems to be the case Phys. Rev. C 93 014314 (2018)].

Differences in the charge radii of mirror nuclei

From experimental r_{ch} in mirror nuclei $L < 60$ MeV



- **some dispersion from models not as nice as for ^{52}Cr ...**

Conclusions in very short

- **An accurate knowledge of neutron skin thickness will lead to an accurate determination of ISB contributions to the EoS in the medium via E_{IAS} (and V_{ISB} in the medium)**

- **Differences in the charge radii of mirror nuclei** may constitute a complementary **(model independent)** tool in **constraining the nuclear EoS**.

Thank you for your attention!