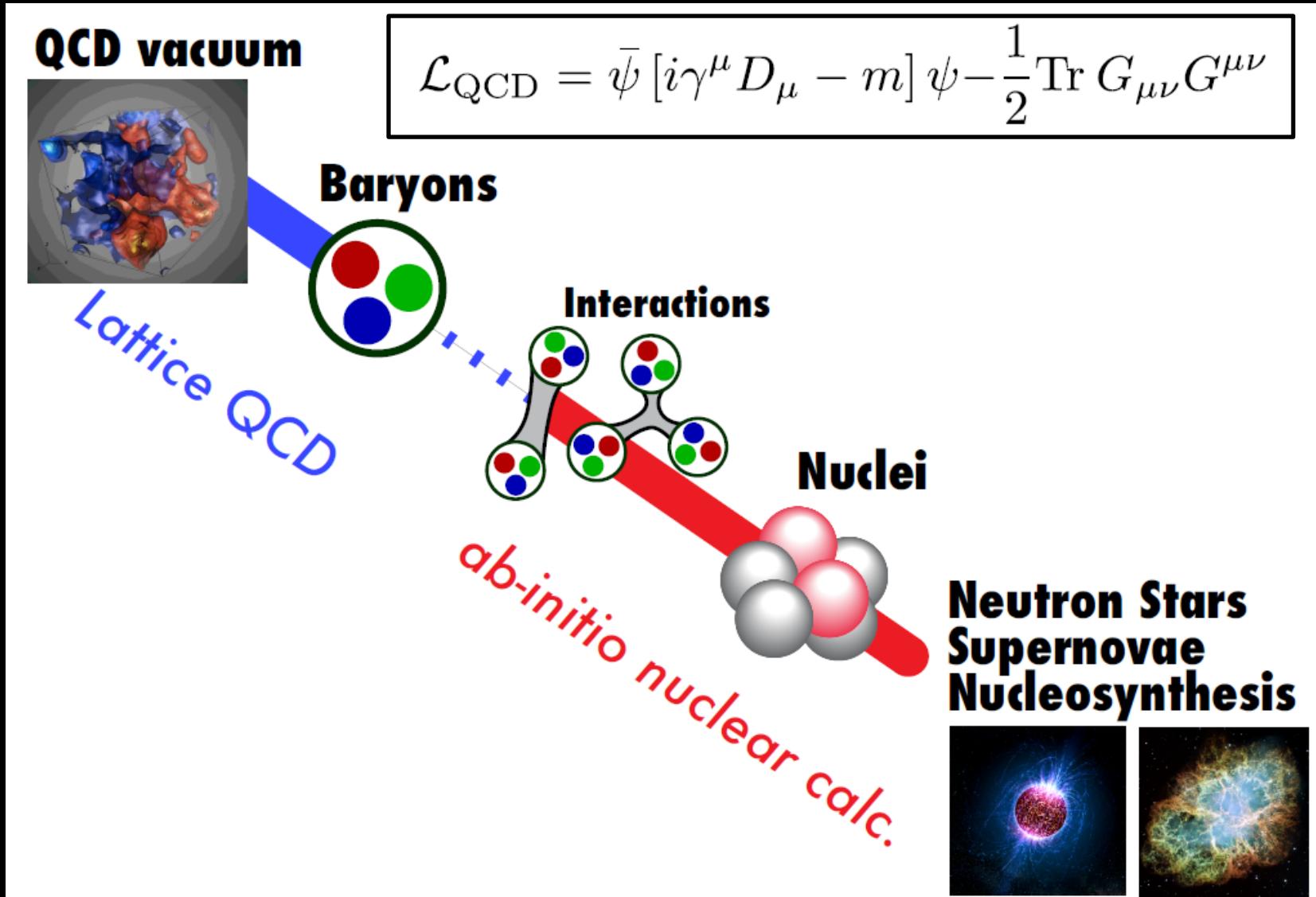


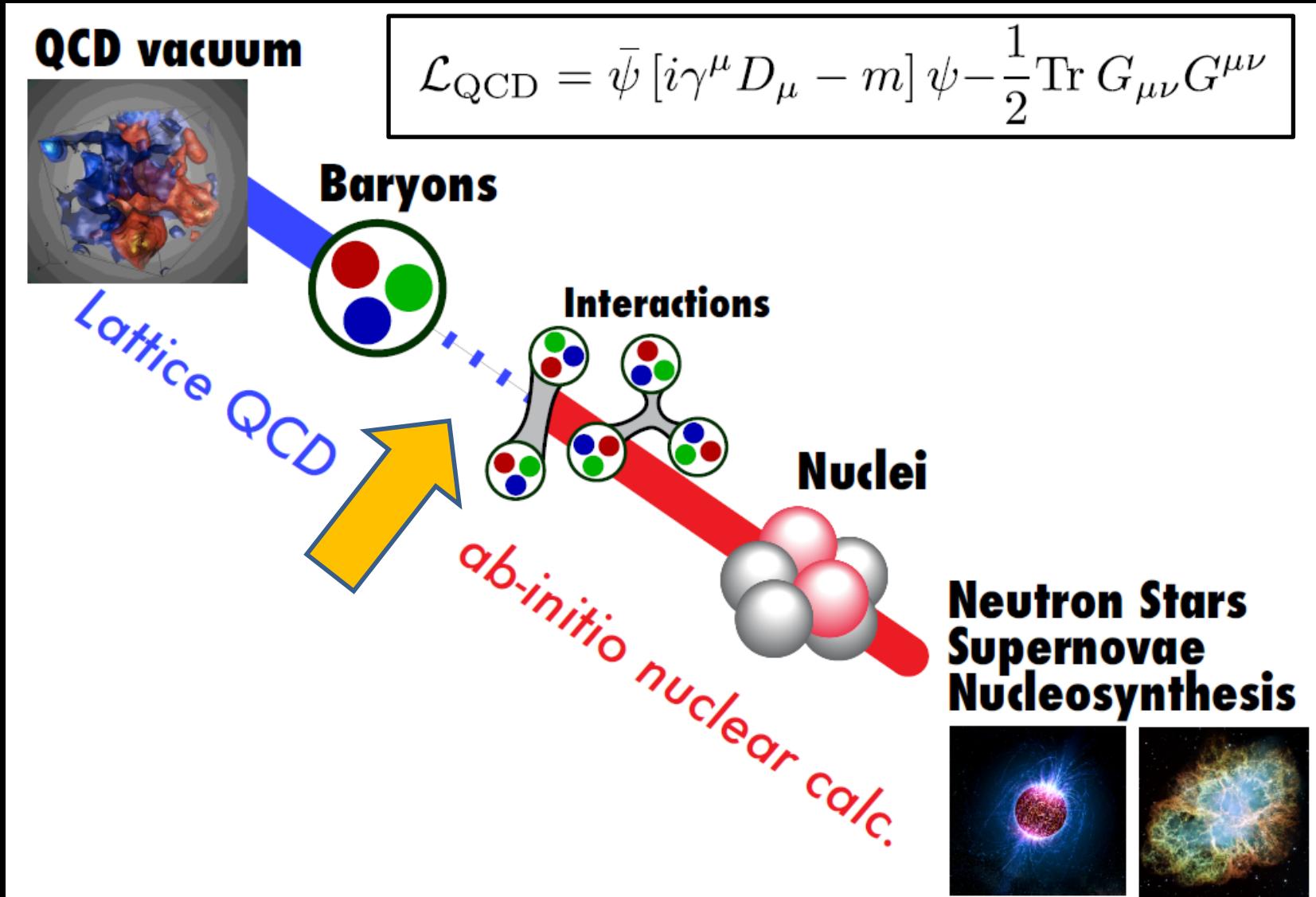
Baryon Interactions from Lattice QCD



Tetsuo Hatsuda (iTHEMS, RIKEN)

“56 International Winter Meeting on Nuclear Physics” (Jan. 26, 2018)

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Contents

I. Precision QCD (5min)

II. Lattice QCD simulations
(10min)

III. Baryon interactions
from LQCD (10min)

IV. Some examples
(10min)

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Lecture Notes in Physics 936

Morten Hjorth-Jensen
Maria Paola Lombardo
Ubirajara van Kolck *Editors*

An Advanced Course in Computational Nuclear Physics

Bridging the Scales from Quarks to Neutron Stars

 Springer

Chapt.3, p.55-91
Lattice QCD (T. Hatsuda)

I. Precision QCD

$$\mathcal{L} = -\frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu} + \bar{q}\gamma^\mu(i\partial_\mu - g t^a A_\mu^a)q - m\bar{q}q$$

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Running masses: $m_q(Q)$

quark masses (from lattice QCD)	[MeV] (MS-bar @ 2GeV)
m_u	2.16 (9)(7)
m_d	4.68 (14)(7)
m_s	93.8 (1.5)(1.9)

FLAG Coll.(2015) <http://itpwiki.unibe.ch/flag/>

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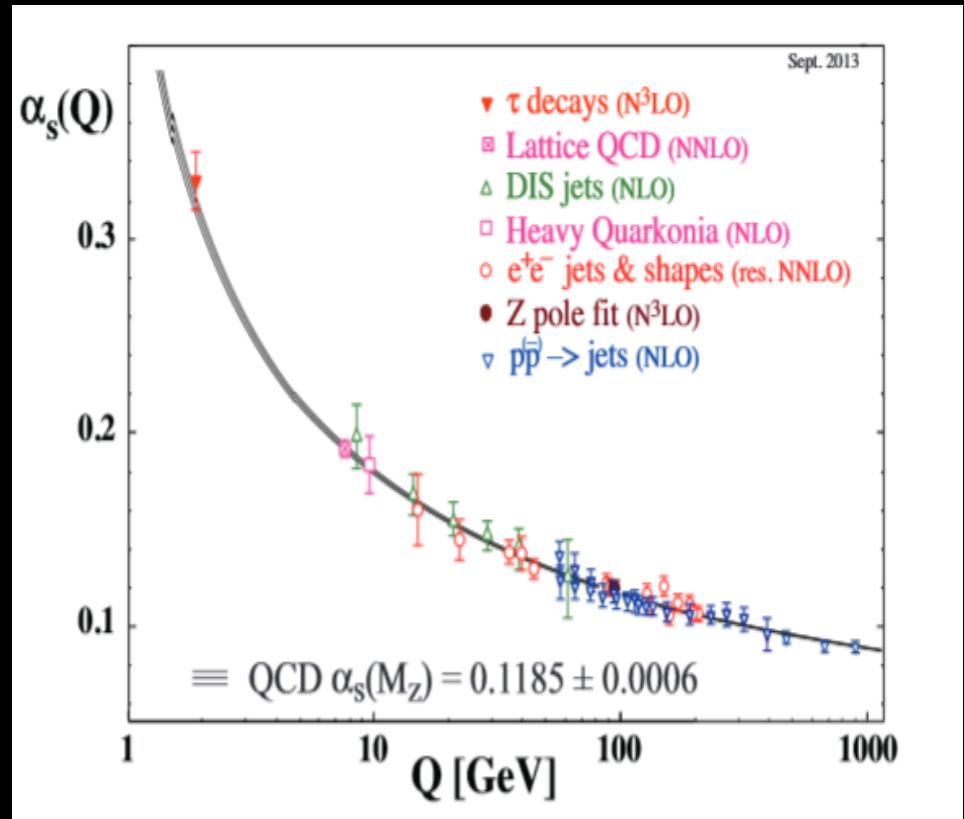
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PDG (2014) <http://pdg.lbl.gov/>

Running coupling: $\alpha_s(Q)=g^2/4\pi$



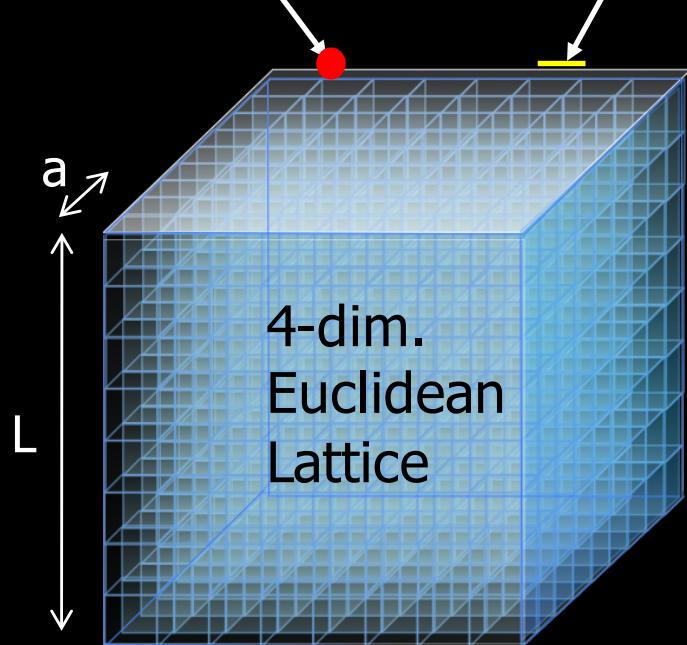
II. Lattice QCD Simulations

$$Z = \int [dU][dq d\bar{q}] \exp \left[- \int d\tau d^3x \mathcal{L}_E \right]$$

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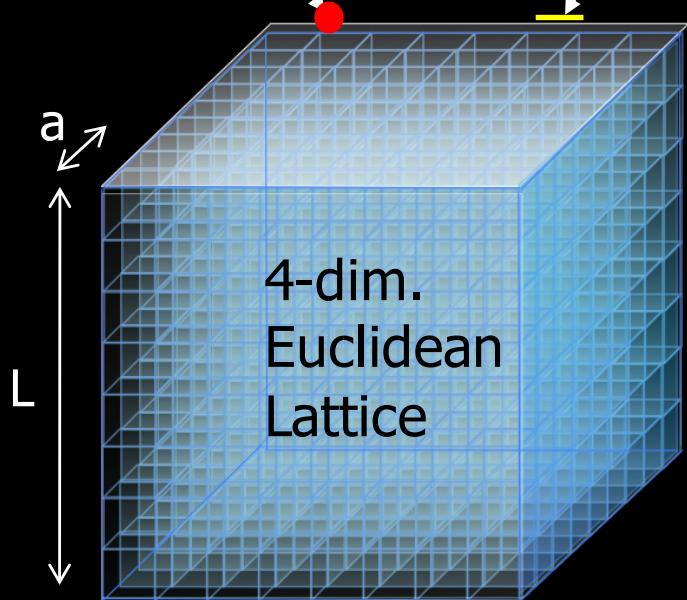
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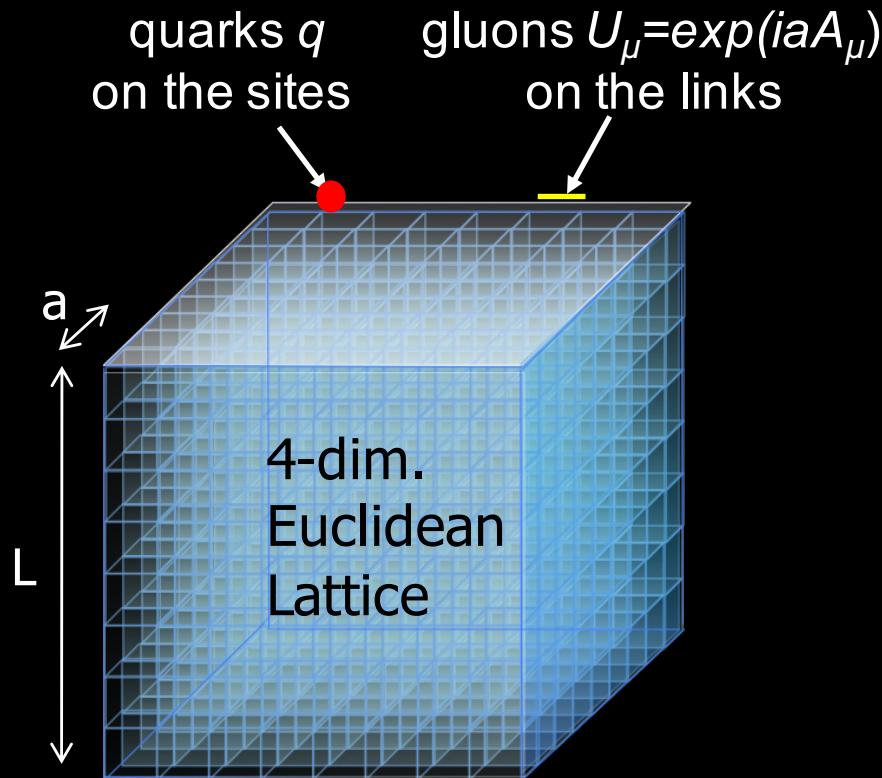
Integration variables

$$N_s^3 \times N_T \times (4_L \times 8_C + 4_S \times 3_C \times N_F)$$

$\sim 10^8$ for $N_s = N_T = 32$ and $N_F = 3$

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Importance Sampling

(In practice; Hybrid MC = MD + Metropolis)

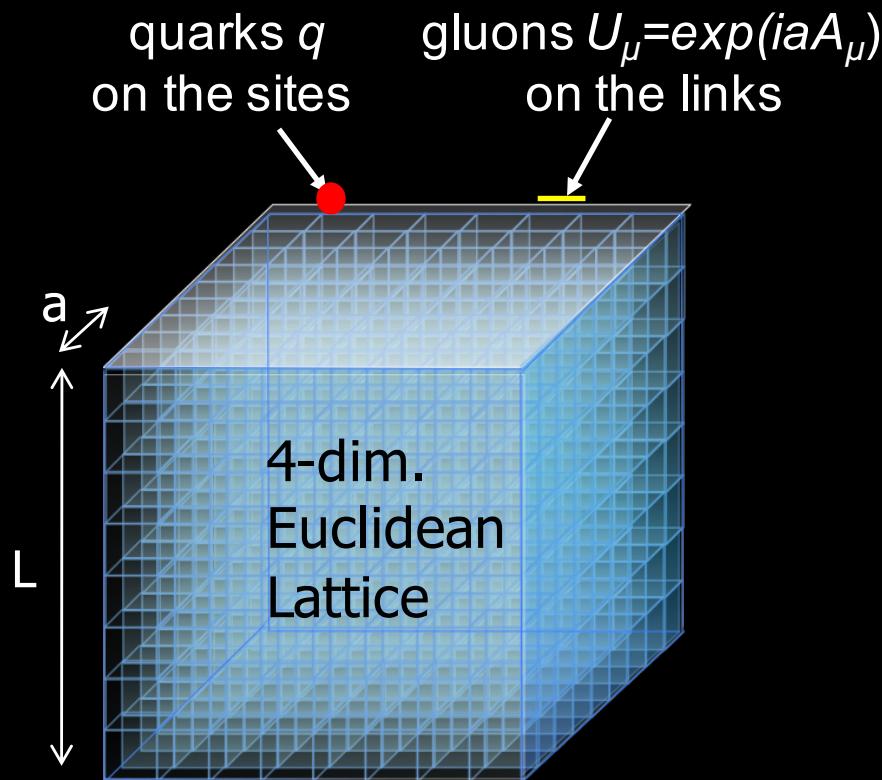
$$\begin{aligned} \langle \mathcal{O} \rangle &= \frac{1}{\mathcal{Z}} \int [d\phi] \mathcal{O}(\phi) e^{-S(\phi)} \\ &= \frac{1}{N} \sum_{n=1}^N \mathcal{O}^{(n)} \pm \sqrt{\frac{\sigma^2}{N}} \end{aligned}$$

Signal Noise

$$\sigma^2 = \frac{1}{N-1} \sum_{n=1}^N (\mathcal{O}^{(n)} - \langle \mathcal{O} \rangle)^2$$

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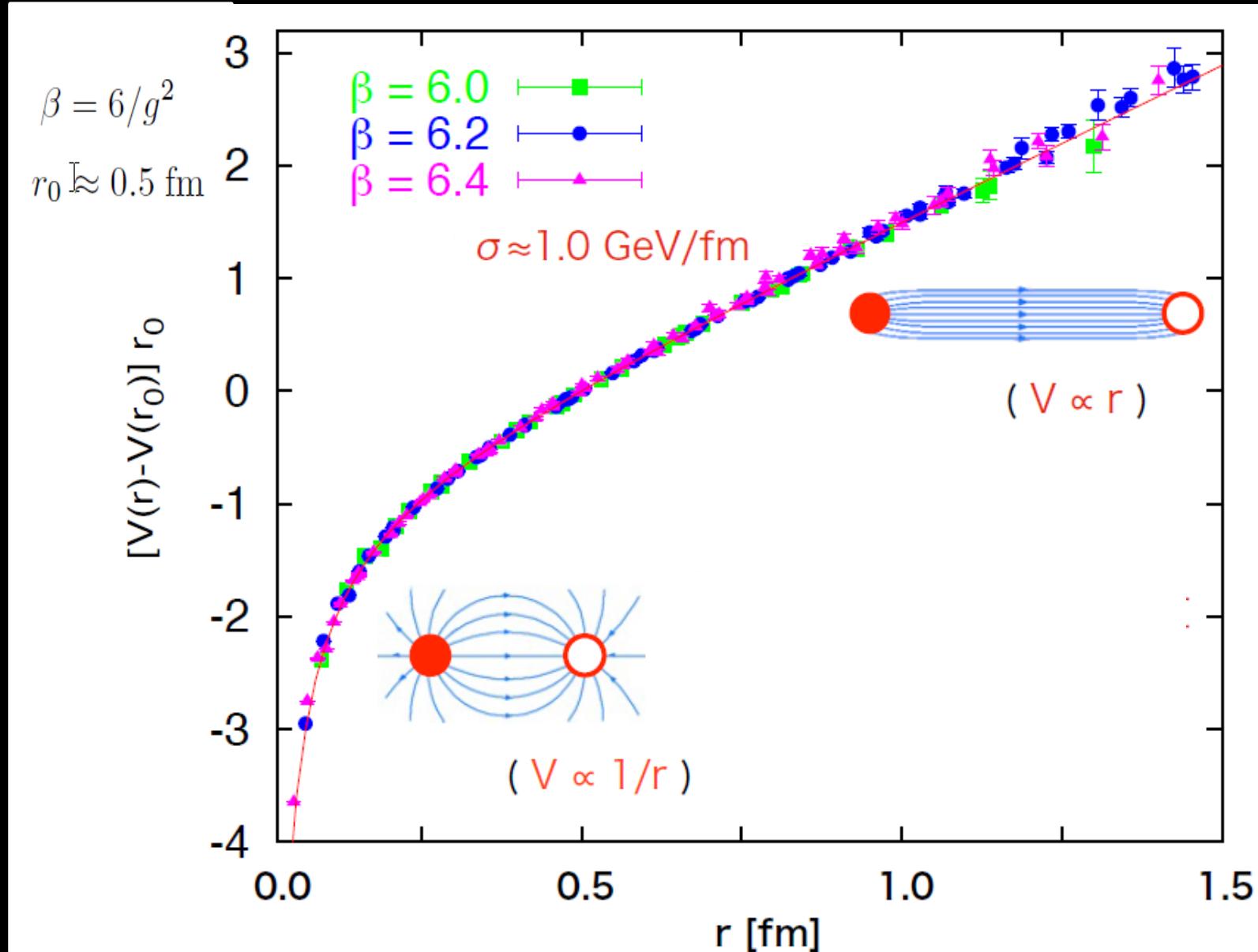
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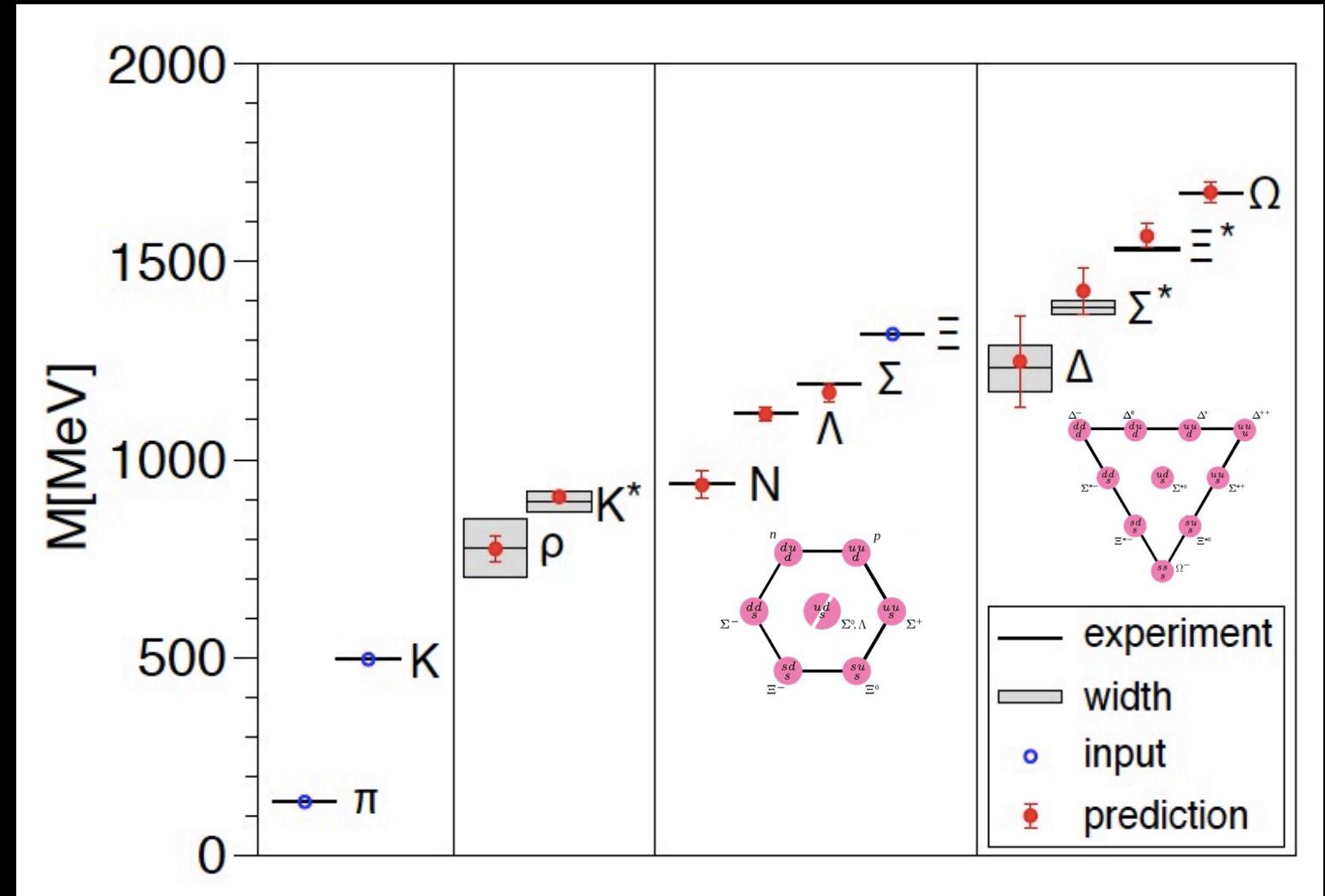
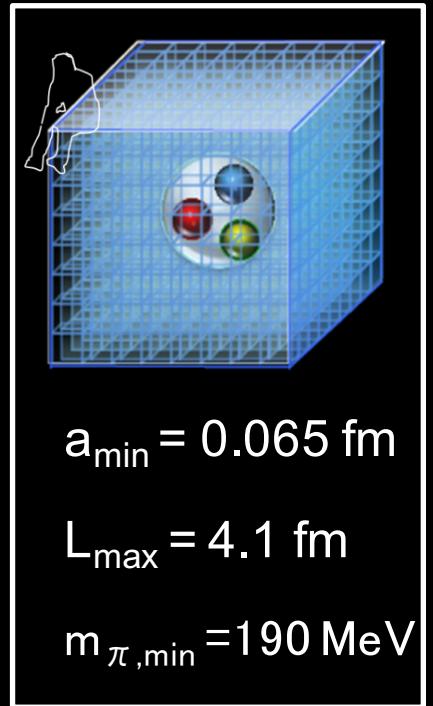
Continuum & Thermodynamic Limits
 $(a \rightarrow 0 \text{ & } L \rightarrow \infty)$

Linear Confinement in LQCD with $N_F=0$

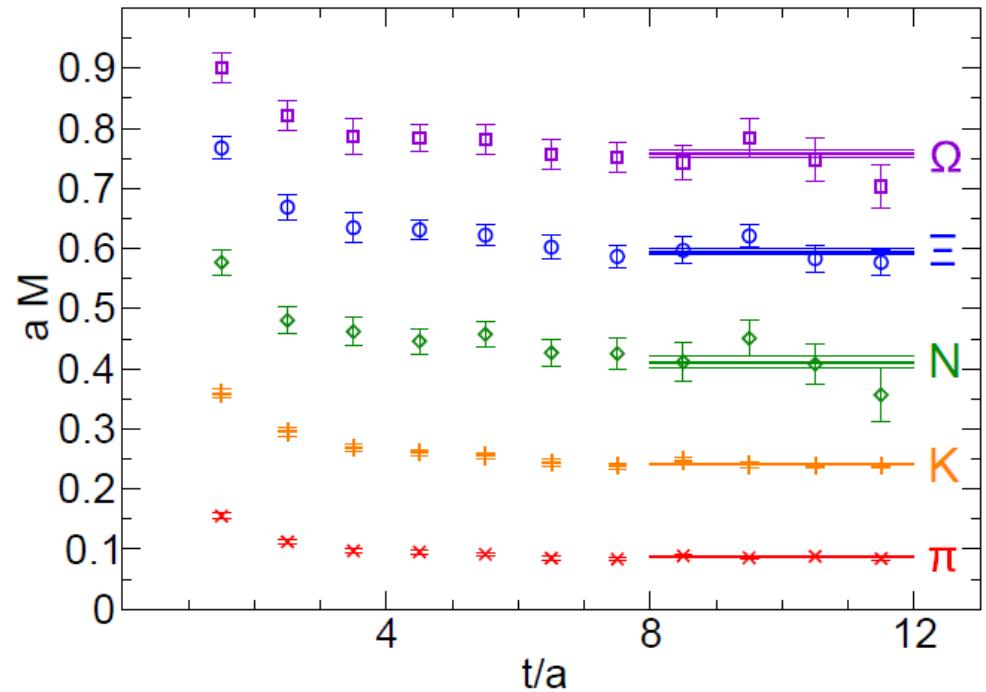


adapted from G. Bali, Phys. Rep. 343 (2001) 1

Light Hadron Masses in LQCD with $N_F=2+1$



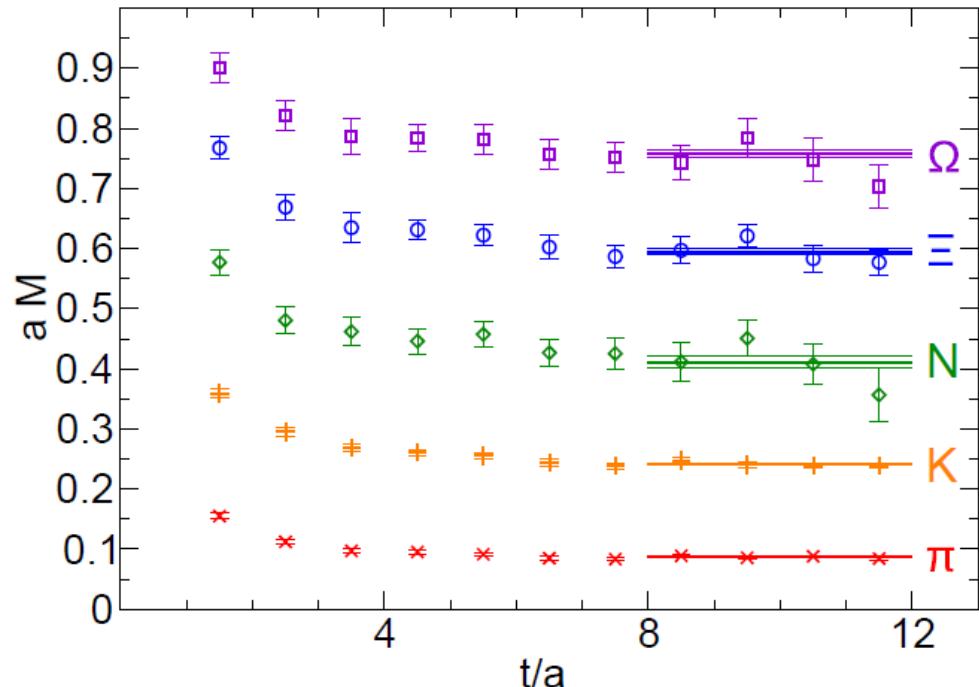
taken from Fodor and Hoelbling, Rev. Mod. Phys. 84 (2012) 449



Effective Mass plot

$$C_H(\tau) \xrightarrow{\tau \rightarrow \infty} |Z_H|^2 e^{-M_H \tau}$$

$$aM_H^{\text{eff}}(\tau) = \ln \left(\frac{C_H(\tau)}{C_H(\tau + a)} \right)$$

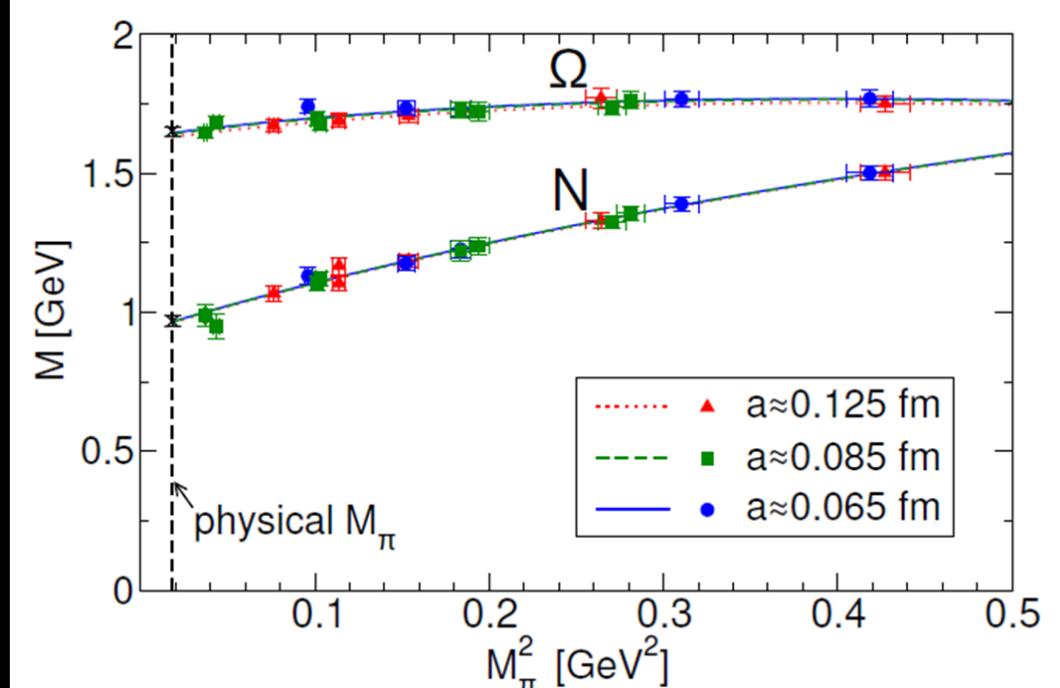


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Chiral
Extrapolation



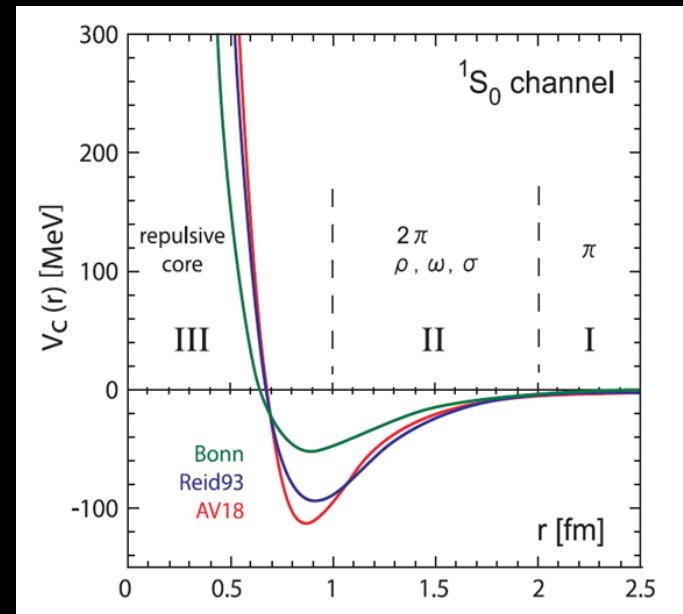
III. Baryon interactions from LQCD ?

- NN int.: about 4500 np and pp scatt. data

“high precision” NN interactions		# of parameters
CD Bonn	(p space)	38
AV18	(r space)	40
EFT in N ³ LO	(nπ+contact)	24

R. Machleidt, arXiv:0704.0807 [nucl-th]

- NNN, YN, YY : data very limited
- YYN, YNN, YYY : data none



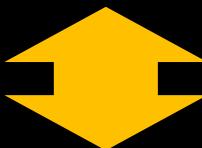
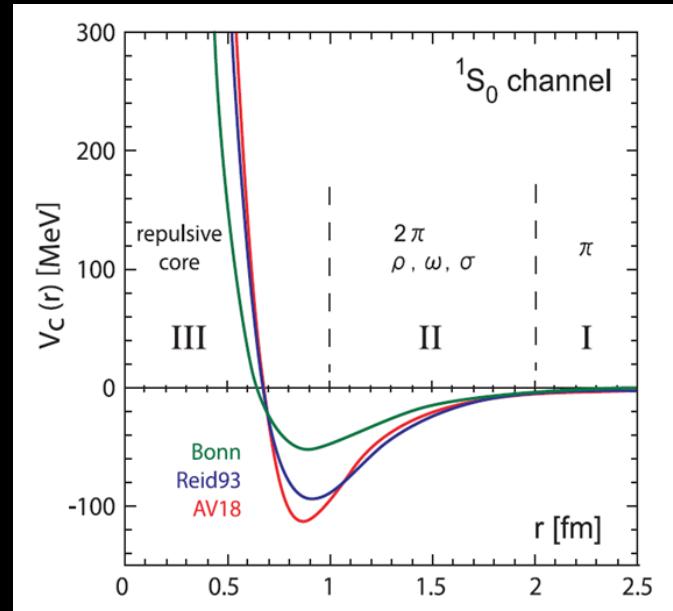
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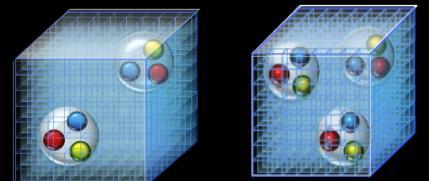
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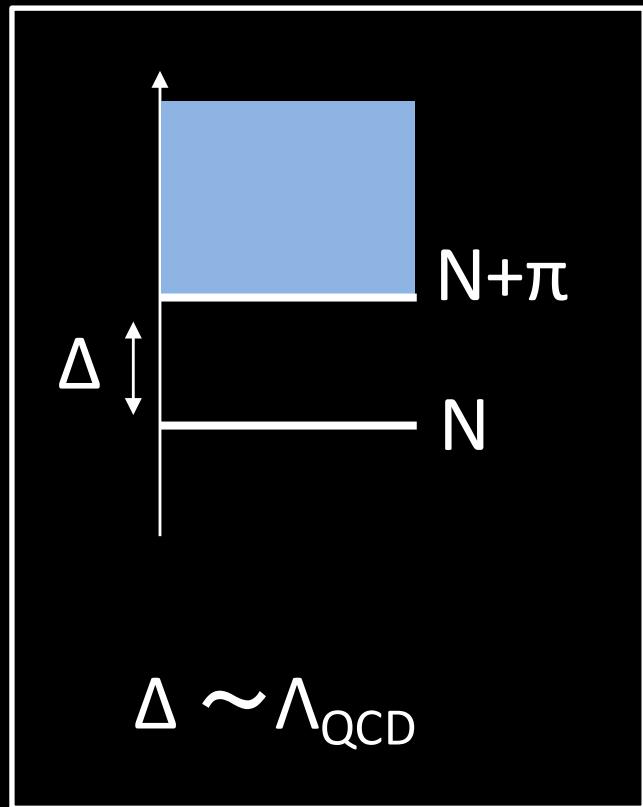
QCD has only 4 parameters: g , $m_{u,d,s}$

→ Derivation of the BB and BBB interactions
from LQCD ?

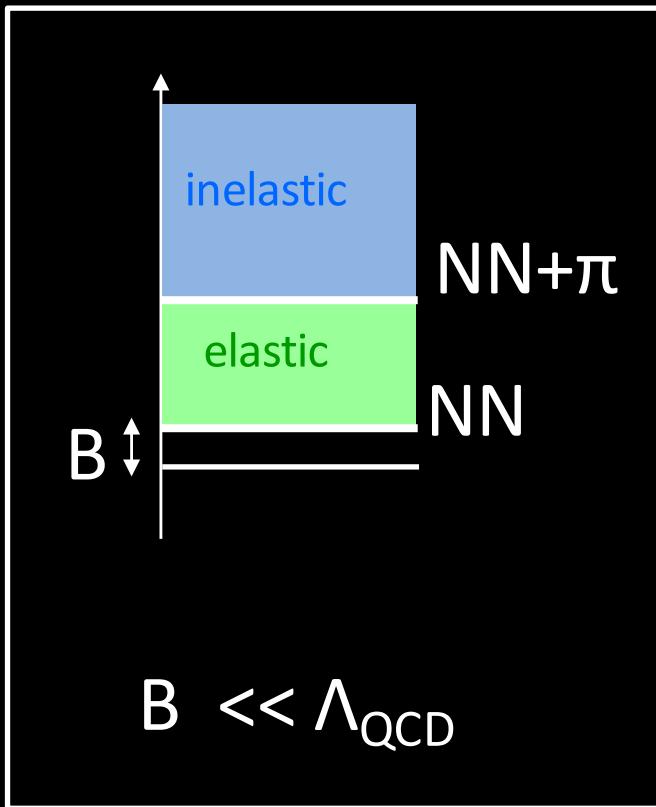


Difference between single-baryon and multi-baryon in QCD

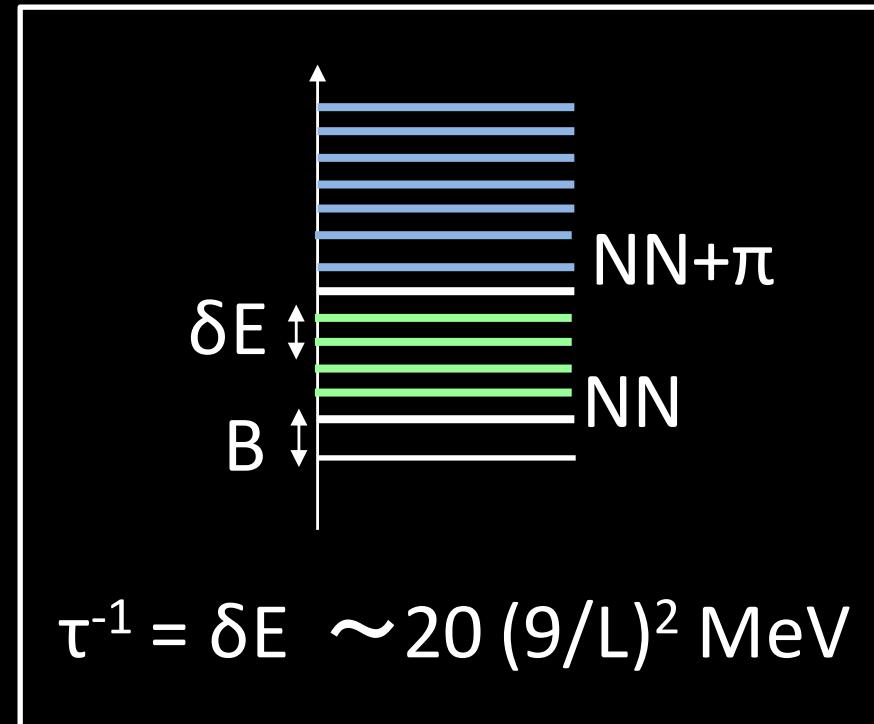
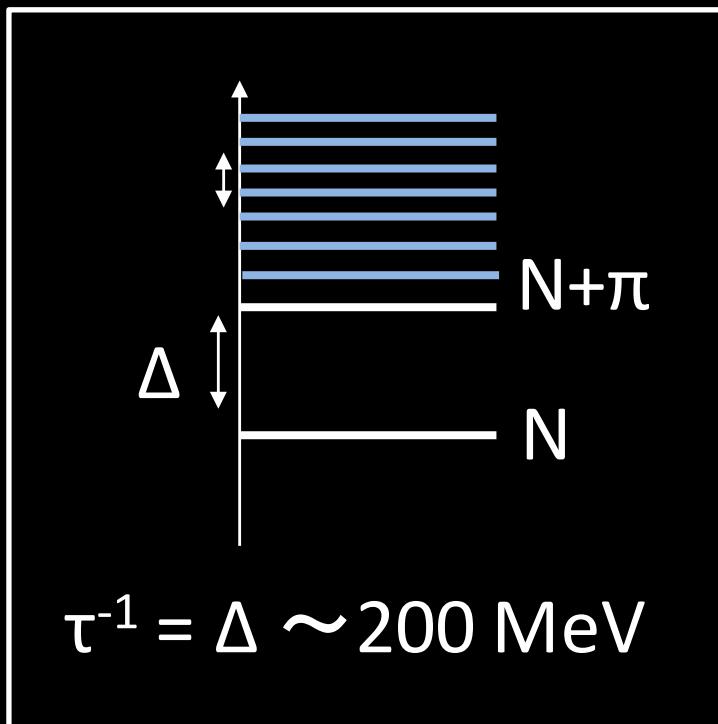
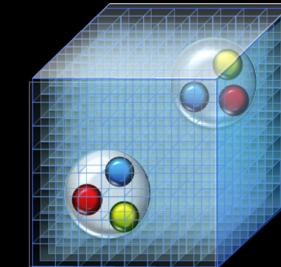
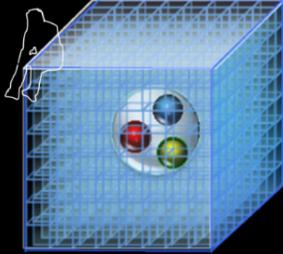
Single nucleon



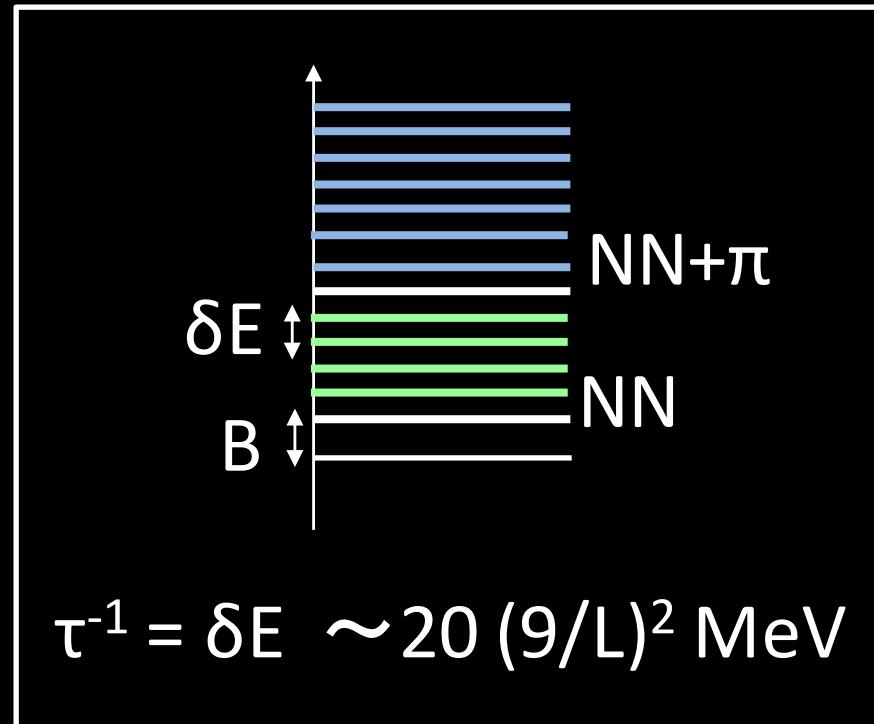
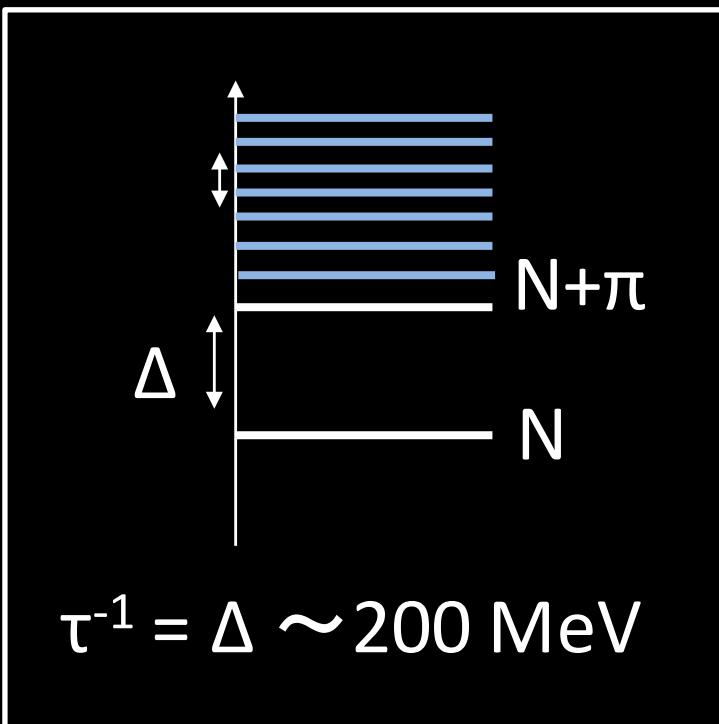
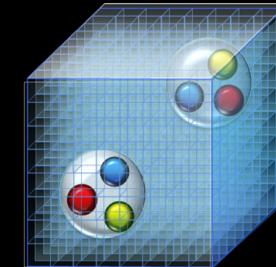
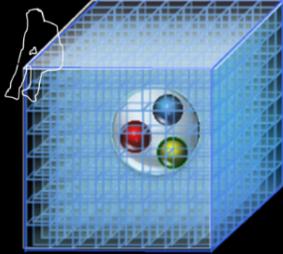
Two nucleons



Difference between single-baryon and multi-baryon in LQCD



Difference between single-baryon and multi-baryon in LQCD



$$S/N \sim \exp(-m_N \tau) \times \sqrt{N}$$

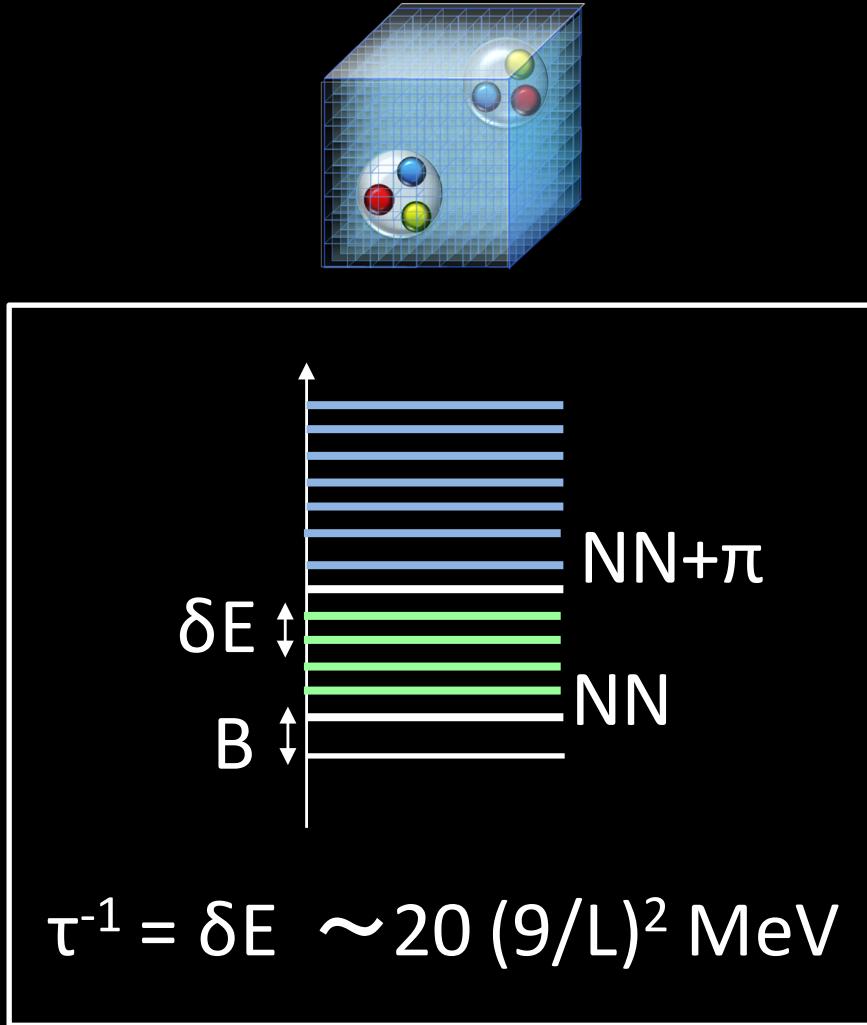
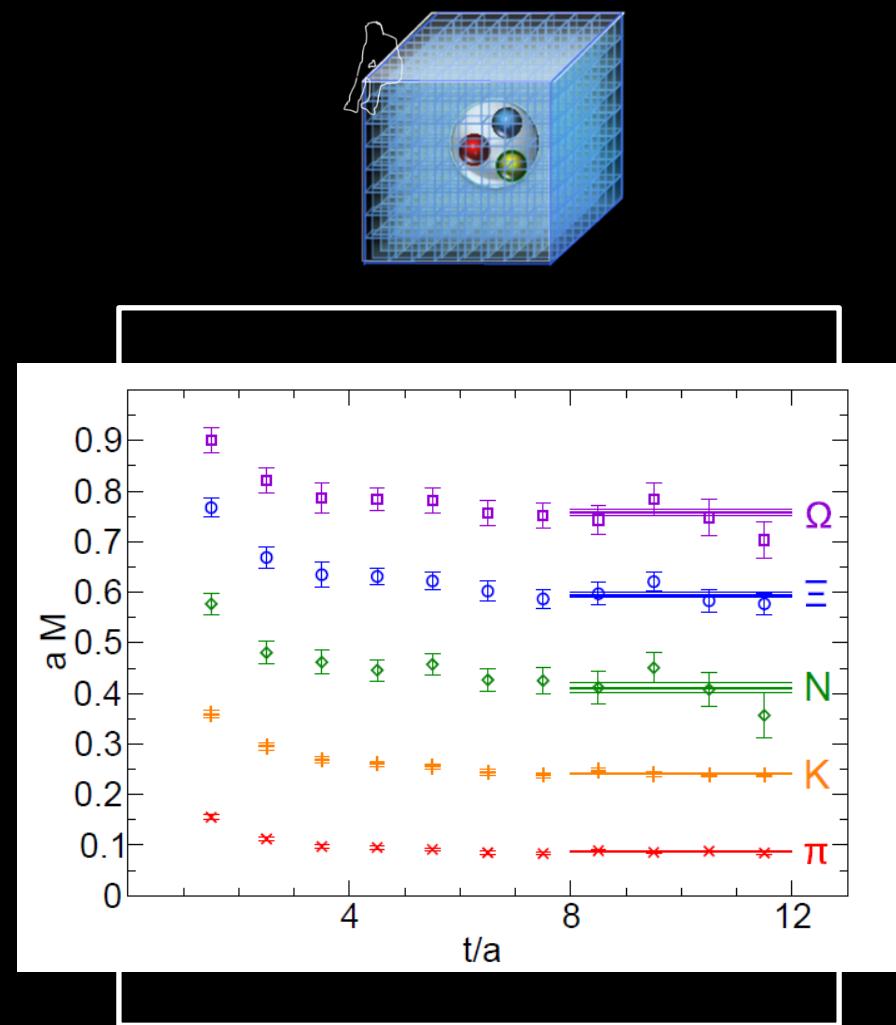
$$\sim 10^{-2} \times \sqrt{N}$$

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Difference between single-baryon and multi-baryon in LQCD



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Space-Time Hadronic Correlation to overcome the difficulty

$$\left\{ \frac{1}{4M_B} \frac{\partial^2}{\partial \tau^2} - \frac{\partial}{\partial \tau} - H_0 \right\} \mathcal{R}(\mathbf{r}, \tau) = \int d^3 r' U(\mathbf{r}, \mathbf{r}') \mathcal{R}(\mathbf{r}', \tau)$$

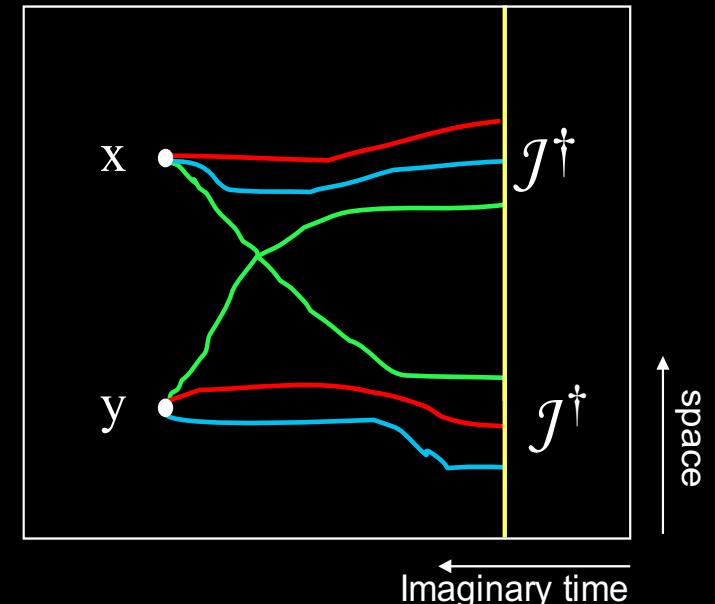
$\mathcal{R}(\mathbf{r}, \tau) \rightarrow U(\mathbf{r}, \mathbf{r}')$ → phase shift, binding energy

$$U(\mathbf{r}, \mathbf{r}') = V(\mathbf{r}, \mathbf{v}) \delta(\mathbf{r} - \mathbf{r}'),$$

$$V(\mathbf{r}, \mathbf{v}) = \underbrace{V_C(r)}_{\text{LO}} + \underbrace{V_T(r) S_{12}}_{\text{NLO}} + \underbrace{V_{LS}(r) \mathbf{L} \cdot \mathbf{S}}_{\text{N}^2\text{LO}} + \mathcal{O}(\mathbf{v}^2) + \dots$$

Ishii, Aoki & Hatsuda, PRL 99 (2007) 022001

Ishii et al. [HAL QCD Coll.], PLB 712 (2012) 437



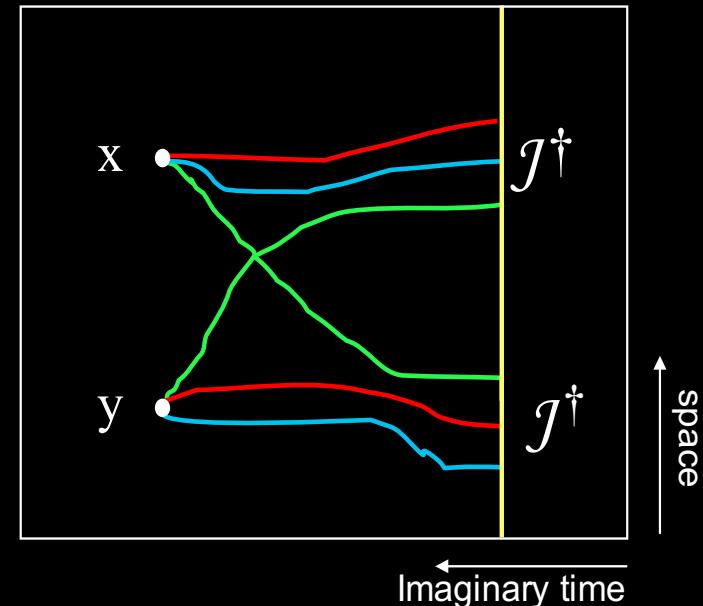
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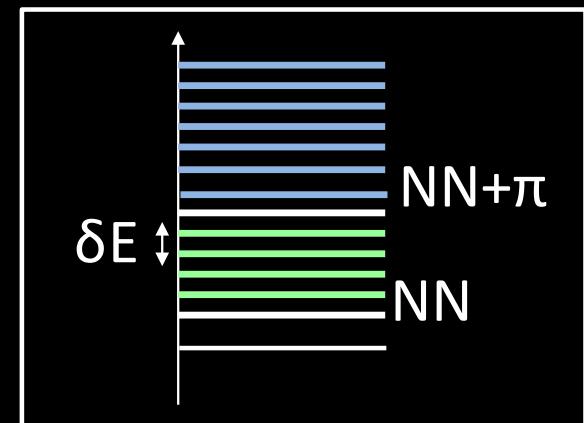
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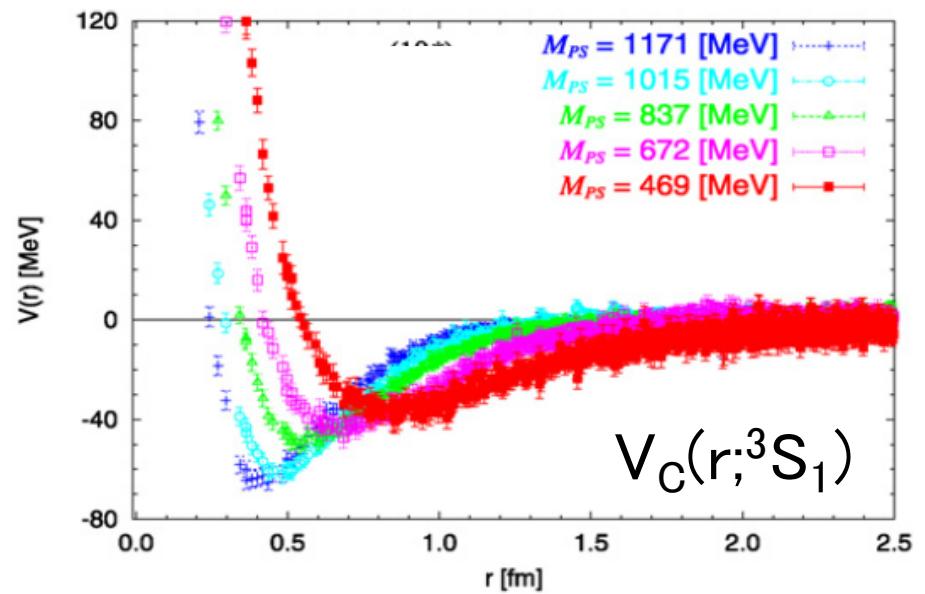
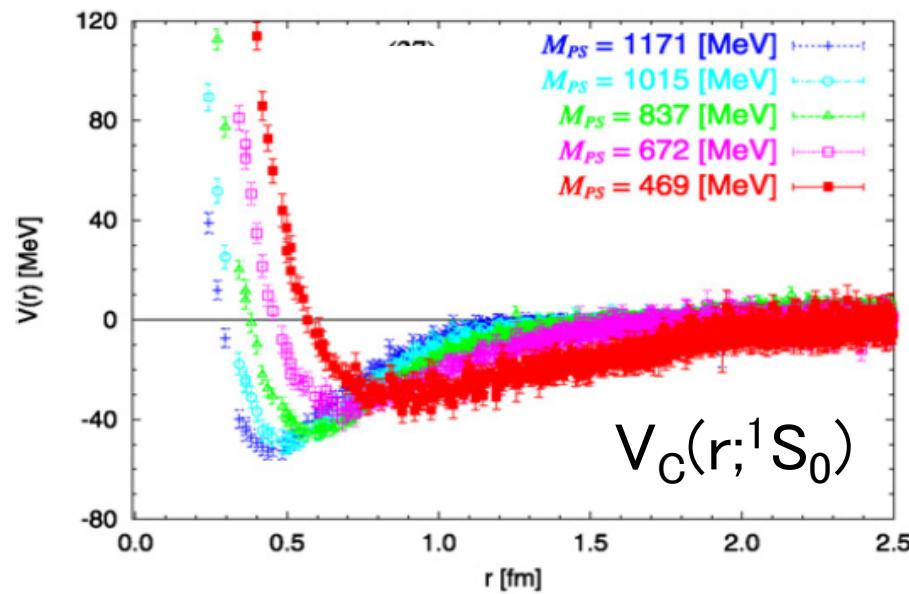
Ishii et al. [HAL QCD Coll.], PLB 712 (2012) 437

	time correlation	Space-time correlation
inelastic	Noise	Noise
elastic	Noise	Signal
ground	Signal	Signal
necessary τ	$\tau > 10 \text{ fm}$	$\tau \sim 1 \text{ fm}$



IV. Some examples

Pilot study ($L=3.9$ fm) in LQCD with $N_F=3$



$$U(\mathbf{r}, \mathbf{r}') = V(\mathbf{r}, \mathbf{v})\delta(\mathbf{r} - \mathbf{r}'),$$

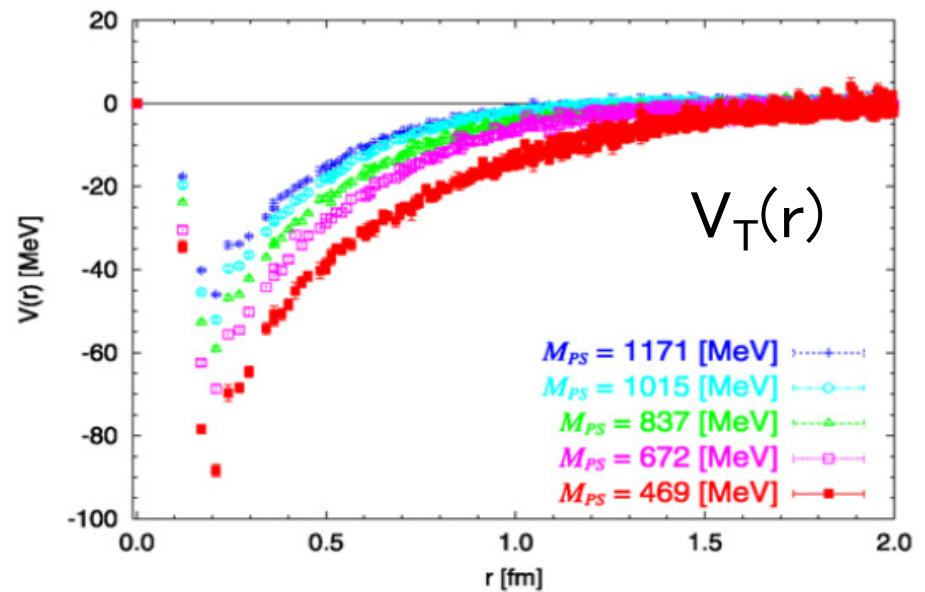
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HAL QCD Coll.

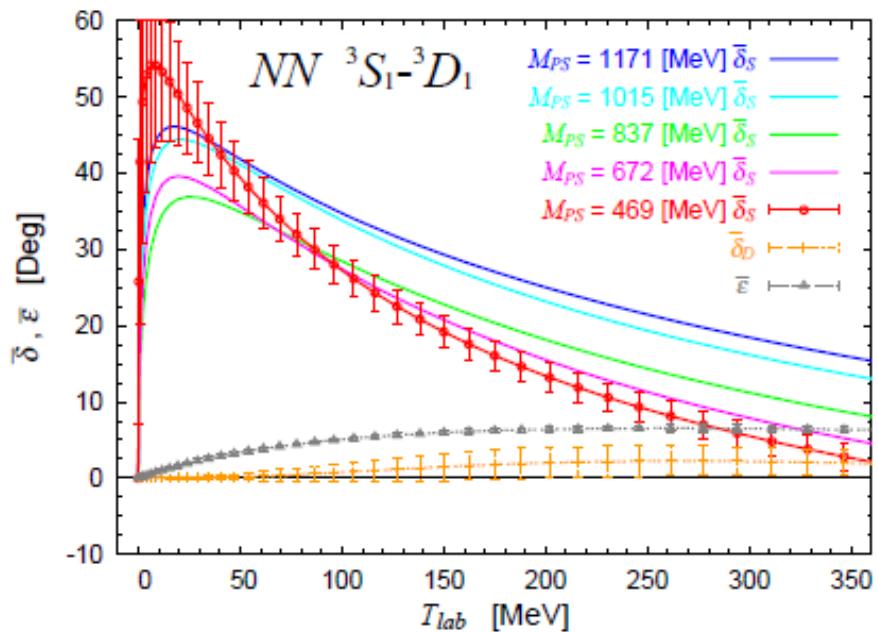
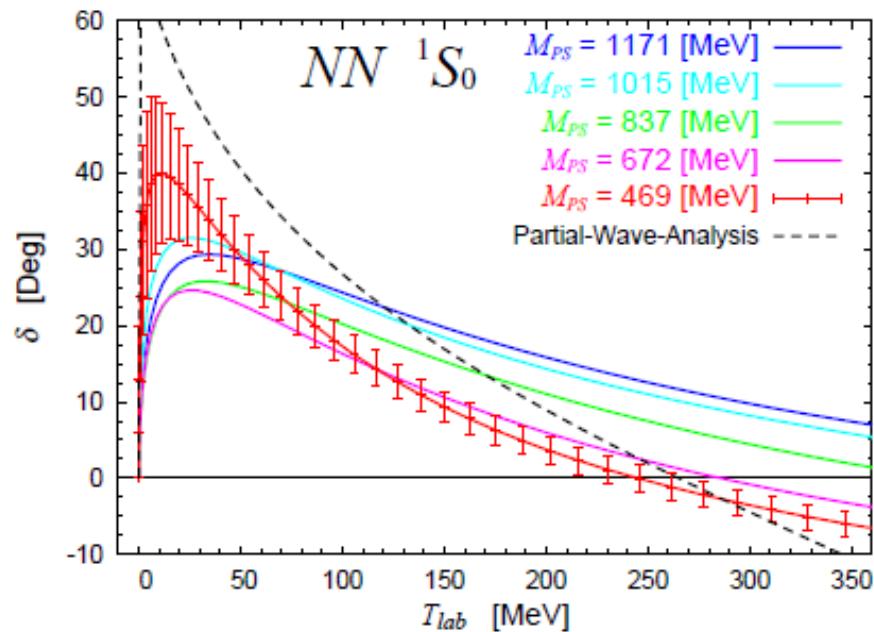
Phys. Rev. Lett. 106 (2011) 162002

Nucl. Phys. A881 (2012) 28

Phys. Rev. Lett. 111 (2013) 112503



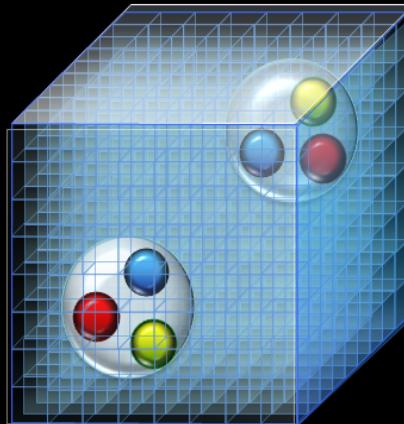
Pilot study ($L=3.9$ fm) in LQCD with $N_F=3$: NN phase shifts



- Stronger attraction in the deuteron channel
- Physical point is close to the unitary region

HAL QCD Coll.,
 Phys. Rev. Lett. 106 (2011) 162002
 Nucl. Phys. A881 (2012) 28
 Phys. Rev. Lett. 111 (2013) 112503

Large scale LQCD simulations with $N_F=2+1$



$a = 0.085 \text{ fm}$

$L = 8.1 \text{ fm}$

$m_\pi = 146 \text{ MeV}$
 $M_K = 525 \text{ MeV}$



$S=0$ $S=-1$ $S=-2$ $S=-3$ $S=-4$ $S=-5$ $S=-6$

NN

N Λ , N Σ

$\Lambda\Lambda, \Lambda\Sigma, \Sigma\Sigma, \mathbf{N}\Xi$

$\Lambda\Xi, \Sigma\Xi$

$\Xi\Xi$

$\Xi\Omega$

$\Omega\Omega$



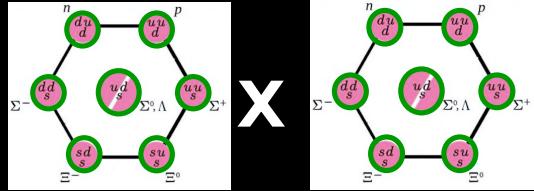
EXP

rich data

LQCD

better S/N

Flavor SU(3) Classification : Two-baryon



$$8 \times 8 = 27 + 8_s + 1 + 10^* + 10 + 8_a$$

$H_{\Lambda\Lambda-N\Xi-\Lambda\Sigma}(J=0)$

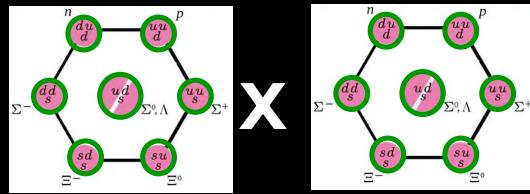
Jaffe (1977)

$D(J=1)$

Rarita-Schwinger (1941)

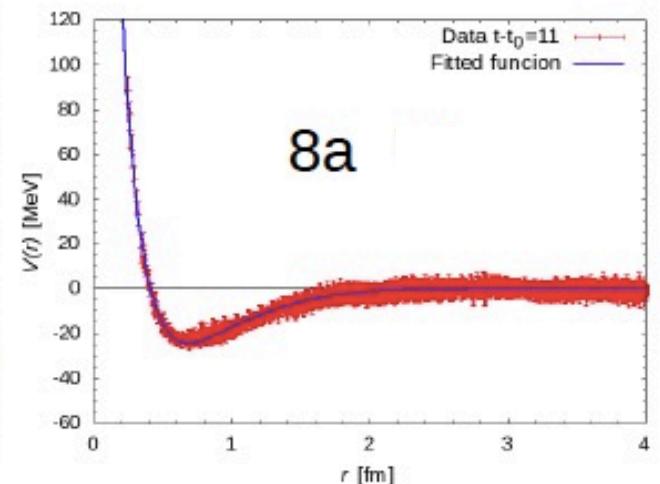
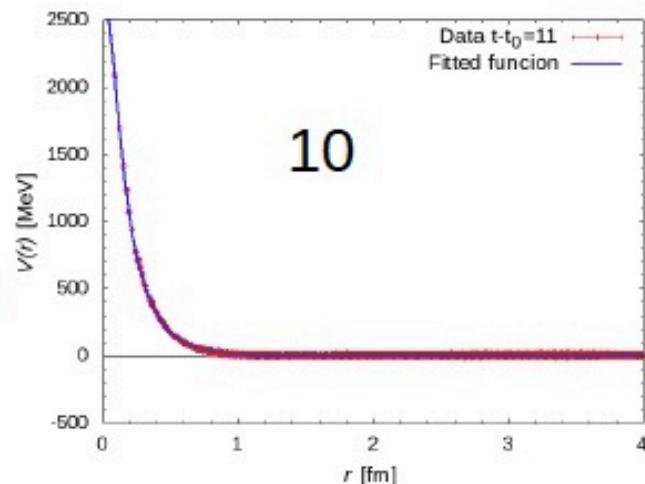
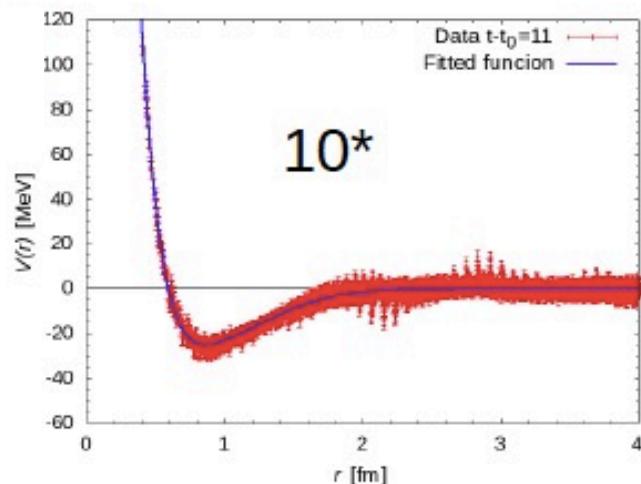
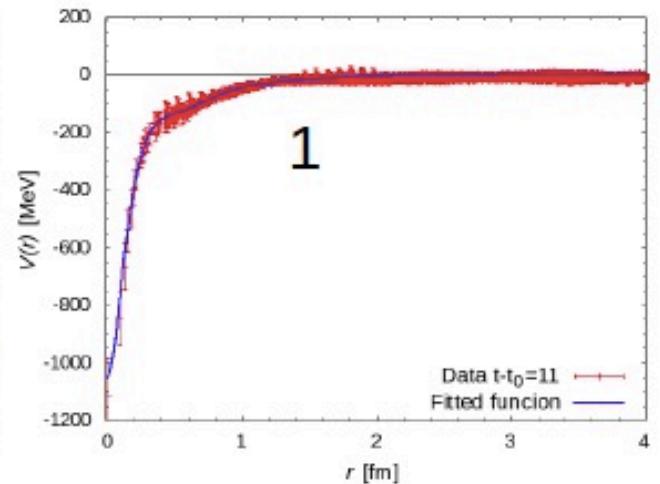
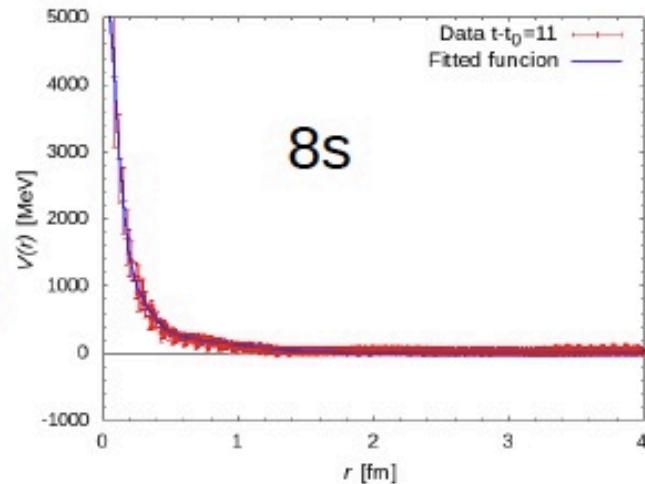
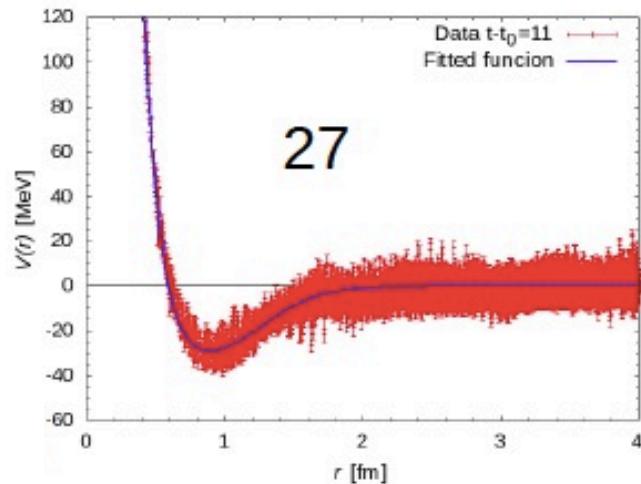
BB interactions in the flavor basis: $V_C(r)$

preliminary

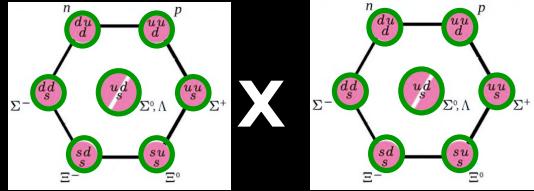


X

$$8 \times 8 = \frac{27 + 8s + 1}{^1S_0} + \frac{10^* + 10 + 8a}{^3S_1, ^3D_1}$$



Flavor SU(3) Classification : Two-baryon



X

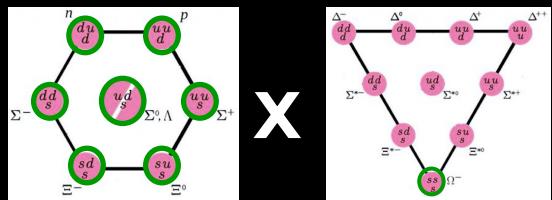
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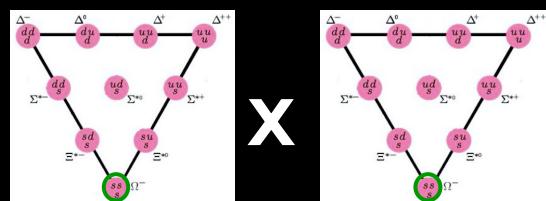


X

$$8 \times 10 = 35 + 8 + 10 + 27$$

$N\Omega(J=2)$

Goldman et al (1987)



X

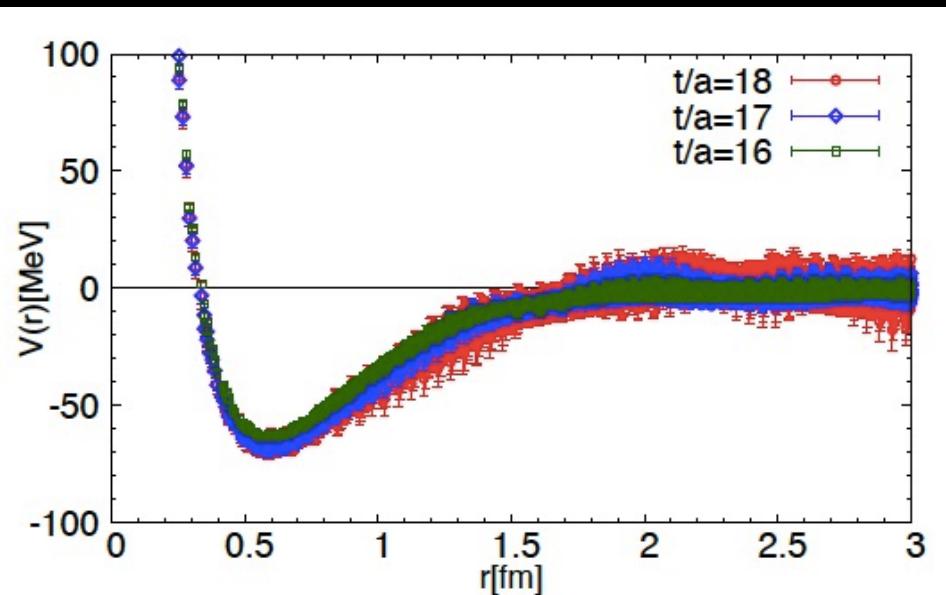
$$10 \times 10 = 28 + 27 + 35 + 10^*$$

$\Omega\Omega(J=0)$

Zhang et al (1997)

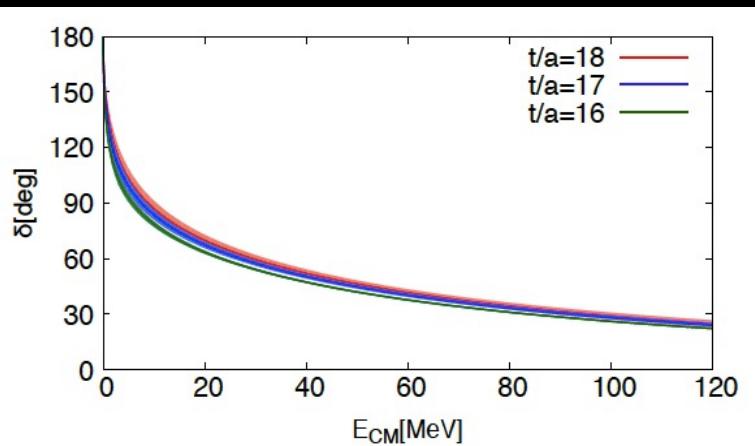
Most strange dibaryon $\Omega\Omega$

HAL QCD Coll.,
arXiv:1709.00654 [hep-lat]

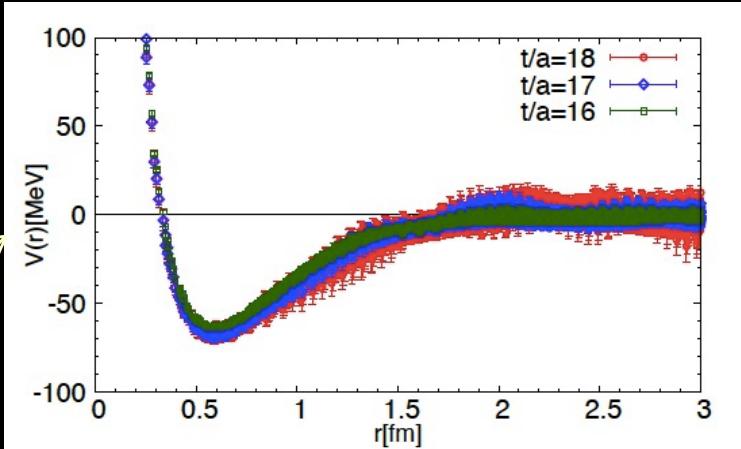


Most strange dibaryon $\Omega\Omega$

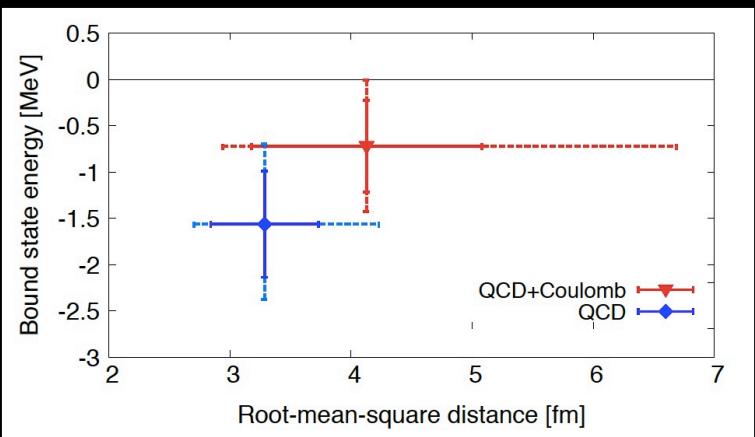
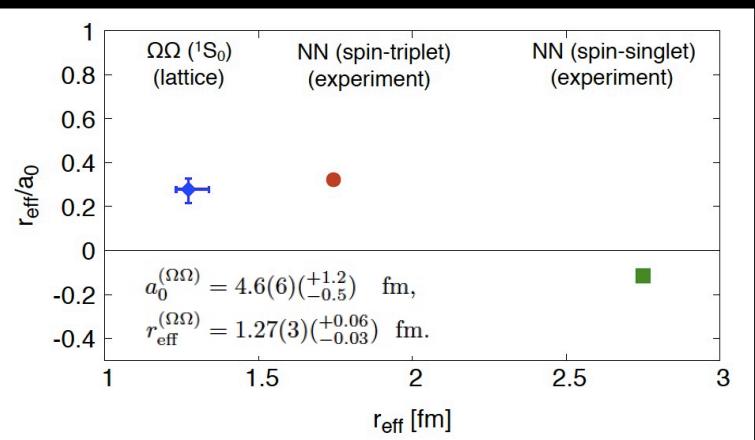
HAL QCD Coll.,
arXiv:1709.00654 [hep-lat]



Phase shift



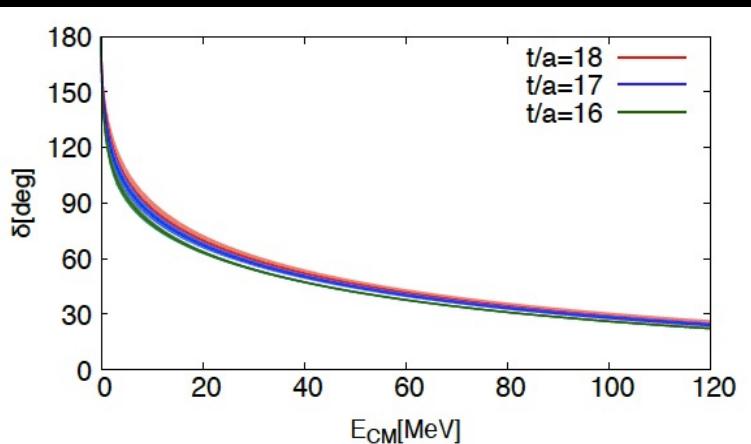
Scatt.length
Effective range



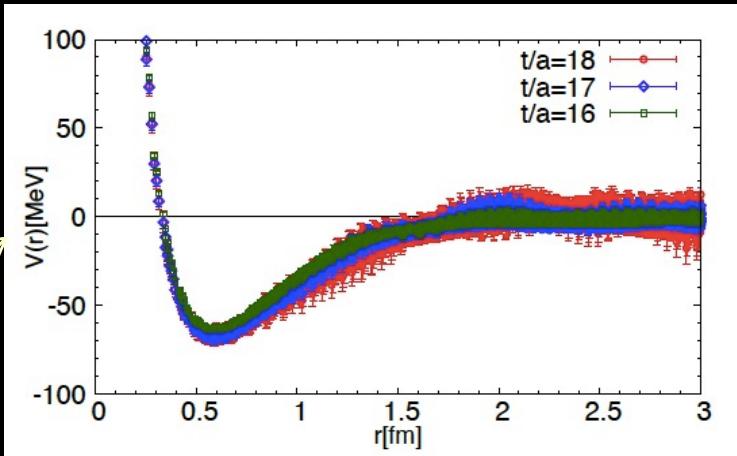
Binding energy
Matter size

Most strange dibaryon $\Omega\Omega$

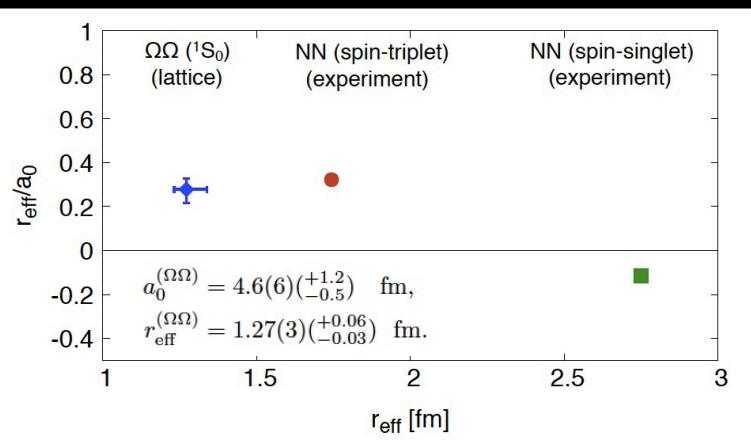
HAL QCD Coll.,
arXiv:1709.00654 [hep-lat]



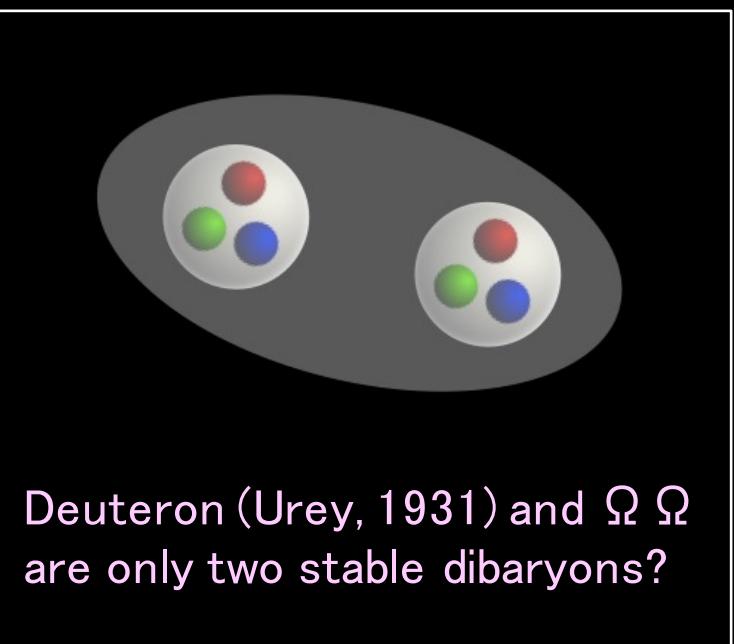
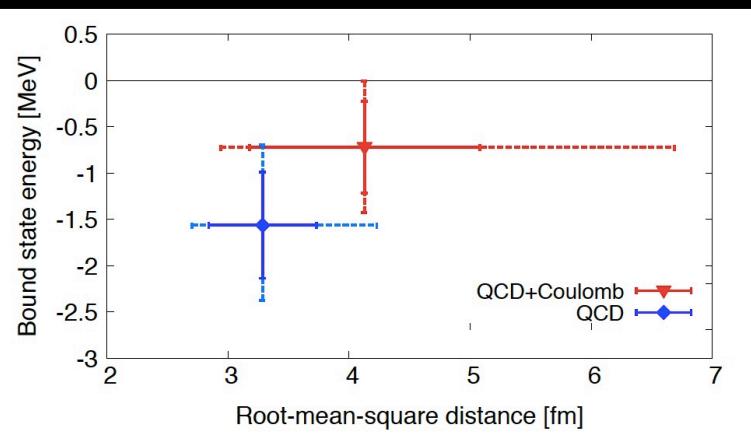
Phase shift



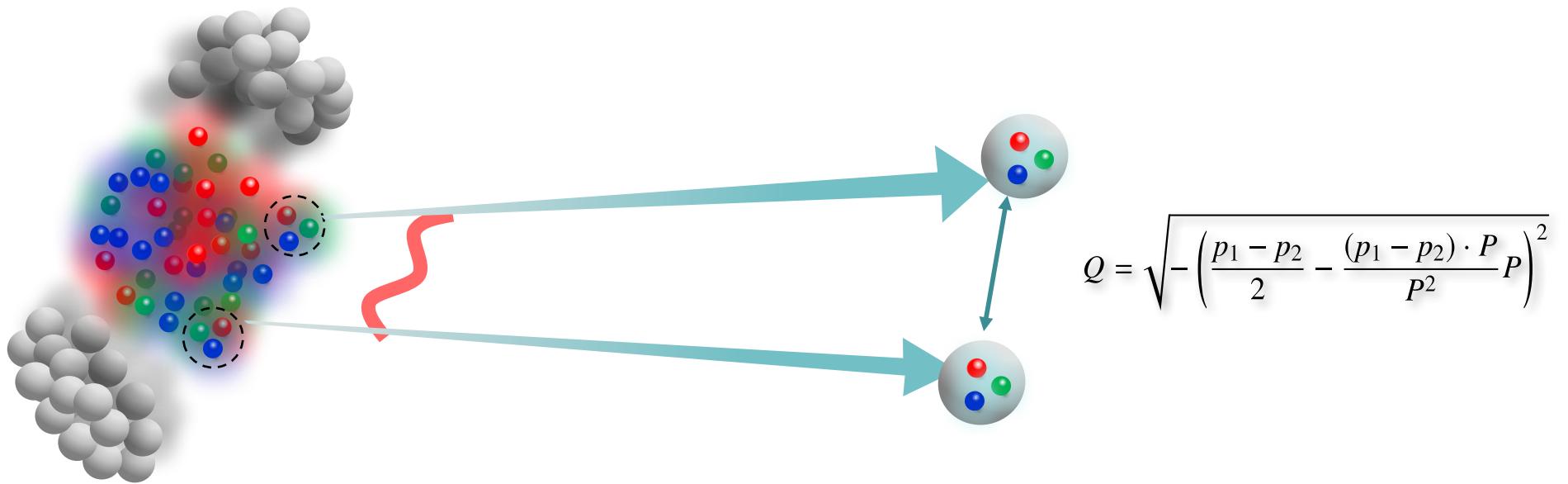
Scatt.length
Effective range



Binding energy
Matter size



How HIC can tell us about interaction?



Measuring Pair Correlation → Constraint on Pairwise Interaction

$$C_{AB}(Q) = \frac{N_{AB}^{\text{pair}}(Q)}{N_A N_B(Q)} = \begin{cases} 1 & \textbf{No Correlation} \\ \text{others} & \textbf{Interaction, Quantum Interference etc} \end{cases}$$

Summary: From QCD to Nuclear Physics

Lattice QCD

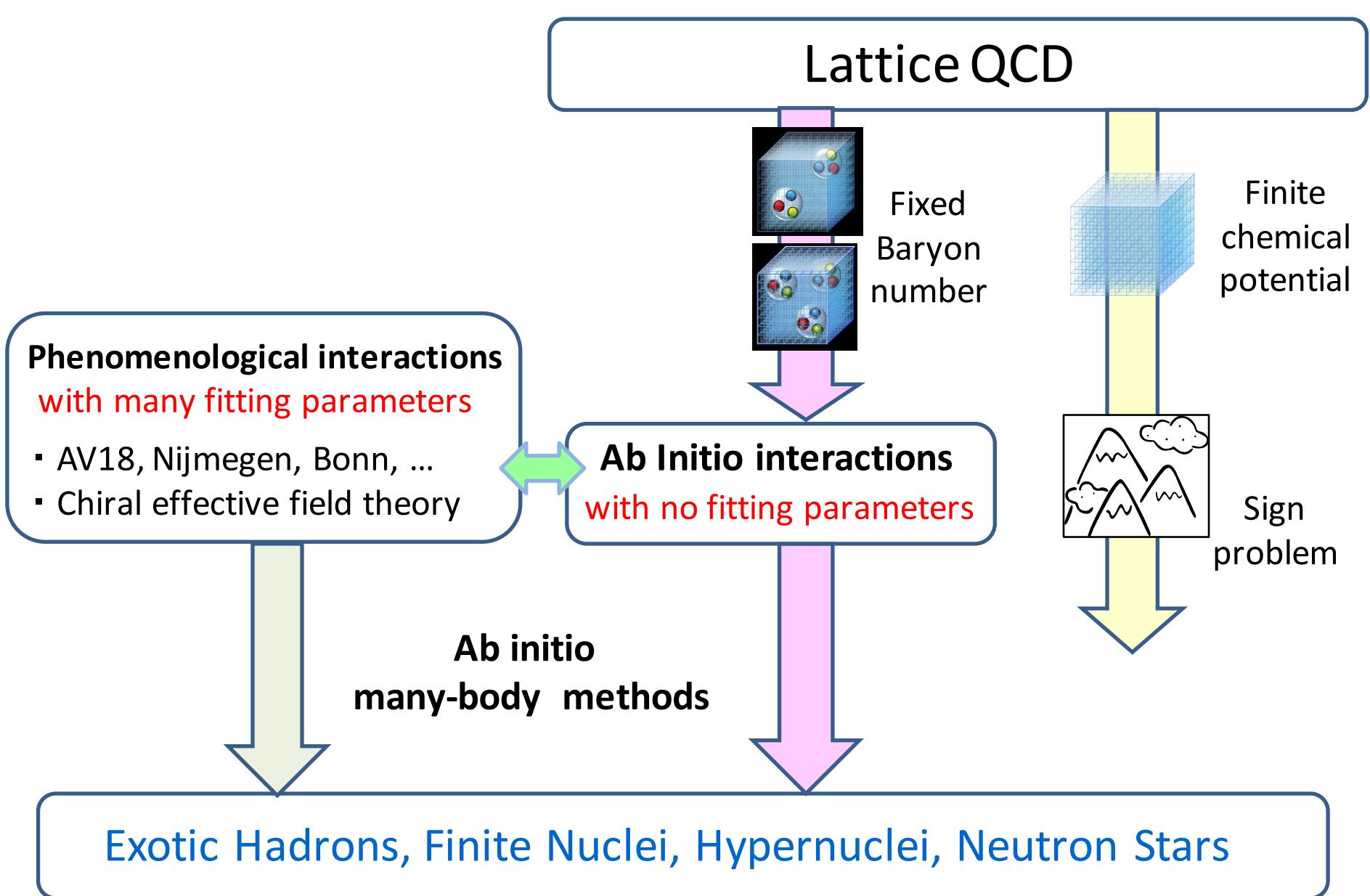
**Phenomenological interactions
with many fitting parameters**

- AV18, Nijmegen, Bonn, ...
- Chiral effective field theory

**Ab initio
many-body methods**

Exotic Hadrons, Finite Nuclei, Hypernuclei, Neutron Stars

Summary: From QCD to Nuclear Physics

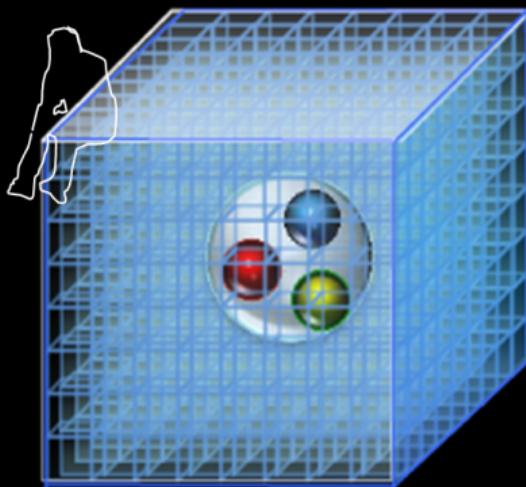


END

Precision LQCD

-- neutron-proton mass difference --

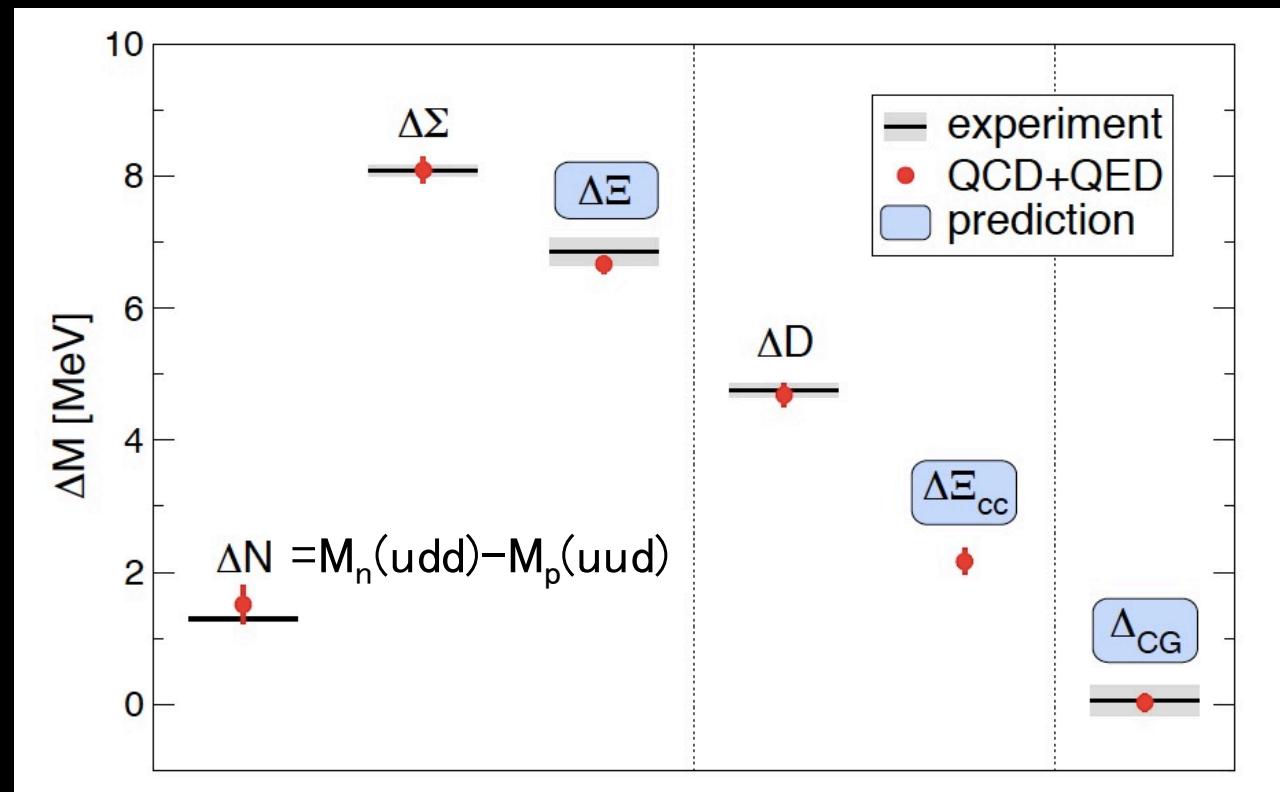
BMW Coll., Science 347 (2015) 1452



$$a_{\min} = 0.054 \text{ fm}$$

$$L_{\max} = 8 \text{ fm}$$

$$m_{\pi,\min} = 190 \text{ MeV}$$



$$(M_n - M_p)_{\text{lat}} = 1.51(16)(23) \text{ MeV}$$

$$(M_n - M_p)_{\text{exp}} = 1.29 \text{ MeV}$$

Problem of Signal to Noise Ratio

Parisi, Lepage (1989)

$$G(r, t) = \langle 0 | \mathcal{O}(r, t) \bar{\mathcal{O}}(0) | 0 \rangle = \sum_n \alpha_n \psi_n(r) e^{-E_n t} \xrightarrow[t \rightarrow \infty]{} \alpha_0 \psi_0(r) e^{-E_0 t}$$

Single pion $\frac{\langle \pi(t)\pi(0) \rangle}{\sqrt{\langle \pi\pi(t)\pi\pi(0) \rangle}} \sim \frac{\exp(-m_\pi t)}{\sqrt{\exp(-2m_\pi t)}} \sim \boxed{\text{const.}}$ Signal/Noise $\sim \sqrt{N_{\text{conf}}}$

Multi pion Signal/Noise $\sim \sqrt{N_{\text{conf}}}$

Single nucleon $\frac{\langle N(t)\bar{N}(0) \rangle}{\sqrt{\langle |N(t)\bar{N}(0)|^2 \rangle}} \sim \frac{\exp(-m_N t)}{\sqrt{\exp(-3m_\pi t)}} \sim \exp[-(m_N - 3/2m_\pi)t]$

Signal/Noise $\sim \exp(-m_N t) \times \sqrt{N_{\text{conf}}}$

Multi nucleon $\frac{\langle N^A(t)\bar{N}^A(0) \rangle}{\sqrt{\langle |N^A(t)\bar{N}^A(0)|^2 \rangle}} \sim \frac{\exp(-\mathbf{A}m_N t)}{\sqrt{\exp(-3\mathbf{A}m_\pi t)}} \sim \exp[-\mathbf{A}(m_N - 3/2m_\pi)t]$

Signal/Noise $\sim \sqrt{\exp(-\mathbf{A}m_N t)} \times \sqrt{N_{\text{conf}}}$

Fake plateaus in temporal correlation for two baryons

"Mirage in temporal correlation functions for baryon-baryon interactions in lattice QCD",
JHEP 10 (2016) 101 by HAL QCD Coll.

"Are two nucleons bound in lattice QCD for heavy quark masses ?
– Sanity check with Luscher's finite volume formula –"

Phys. Rev. D96 (2017) 034521 by HAL QCD Coll.

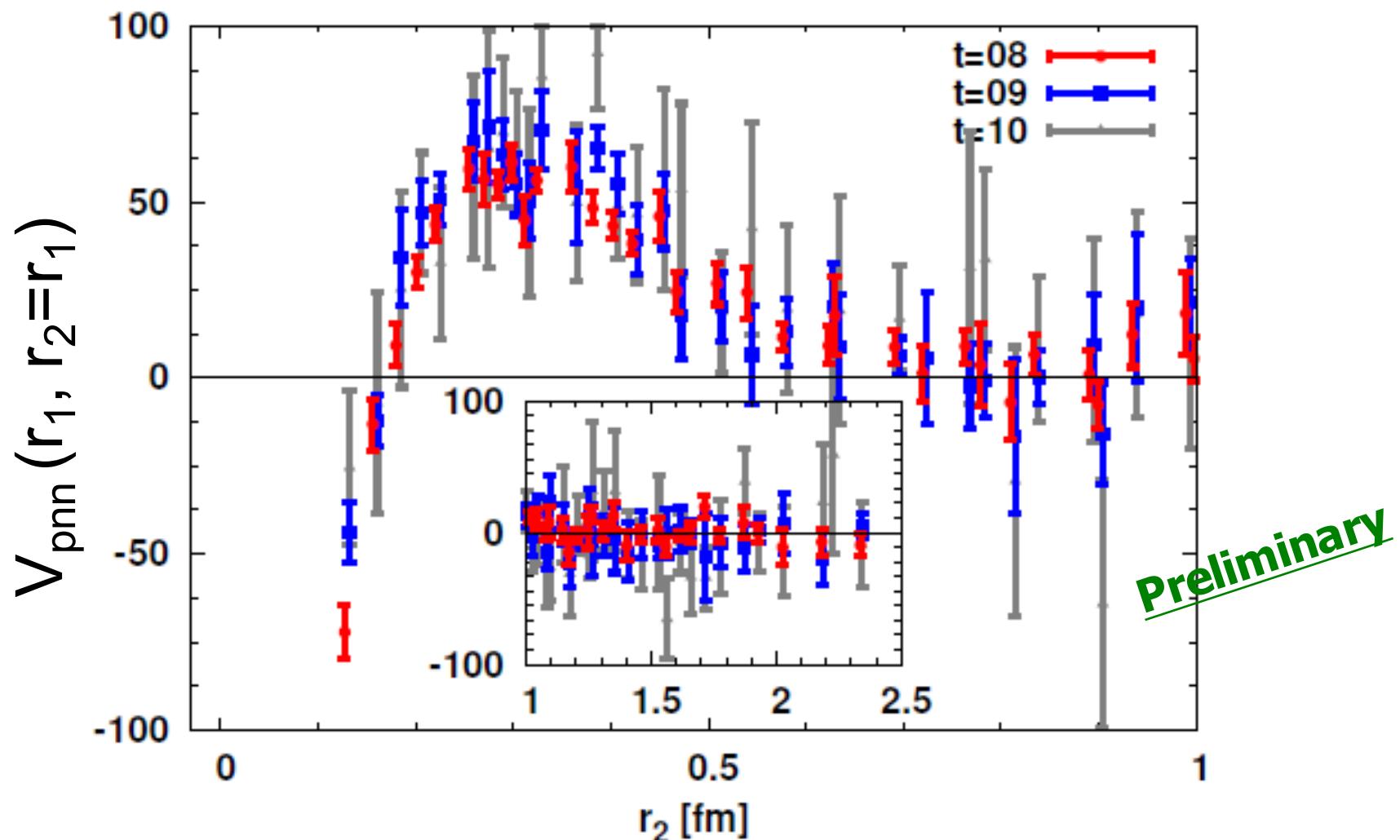
"Sanity check for NN bound states in lattice QCD with Luscher's finite volume
formula -- Exposing Symptoms of Fake Plateaux -- "

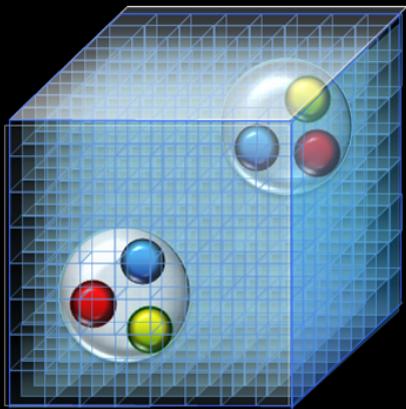
arXiv:1707.08800 [hep-lat] by Aoki, Doi, Iritani

Data	Source independence	NN(1S_0)			NN(3S_1)		
		Consistency check			Consistency check		
		(i)	(ii)	(iii)	(i)	(ii)	(iii)
YKU2011 [24]	†	No	No	*	†	No	No
YIKU2012 [25]	No	†	No	*	No	†	No
YIKU2015 [26]	†	†	No	*	†	†	No
NPL2012 [27]	†	†	No	*	†	†	*
NPL2013 [28,29]	No	*	*	No	No	*	?
NPL2015 [30]	†	No	*	No	†	No	*
CalLat2017 [31]	No	?	*	No	No	?	No

Pilot Study ($L=2.9\text{-}5.8$ fm) in (2+1)-flavor QCD :
3N force in Triton channel

$m_\pi = 0.51 \text{ GeV}$

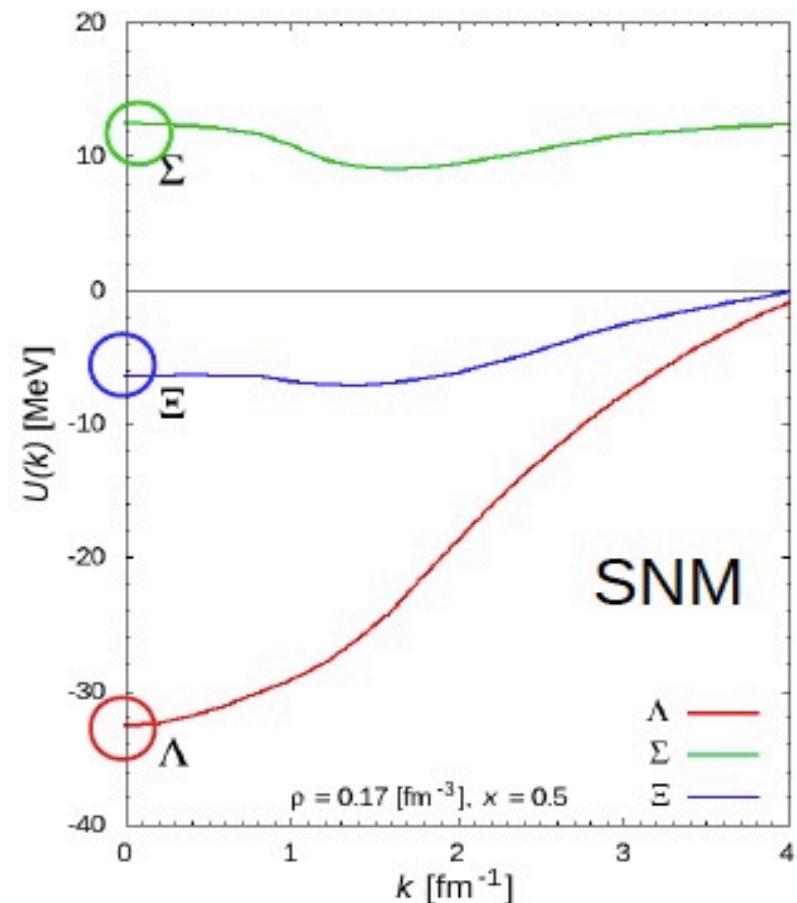
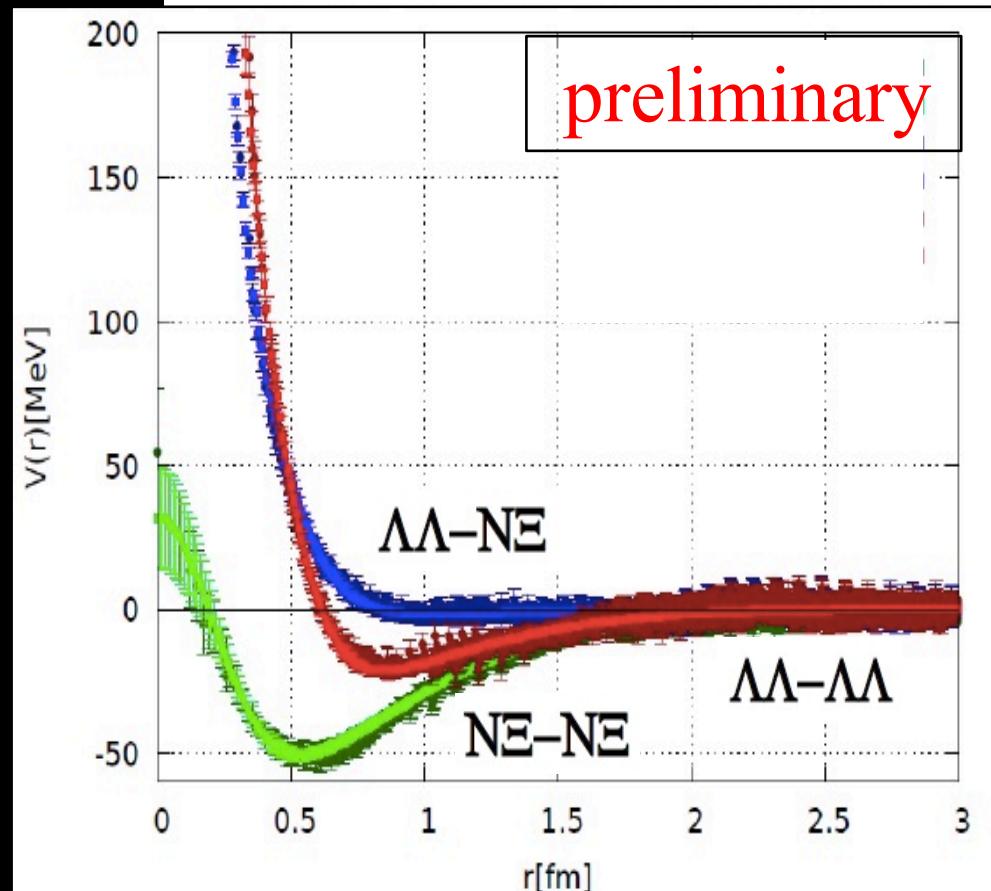




$L = 8.1 \text{ fm}$
 $M_\pi \sim 146 \text{ MeV}$

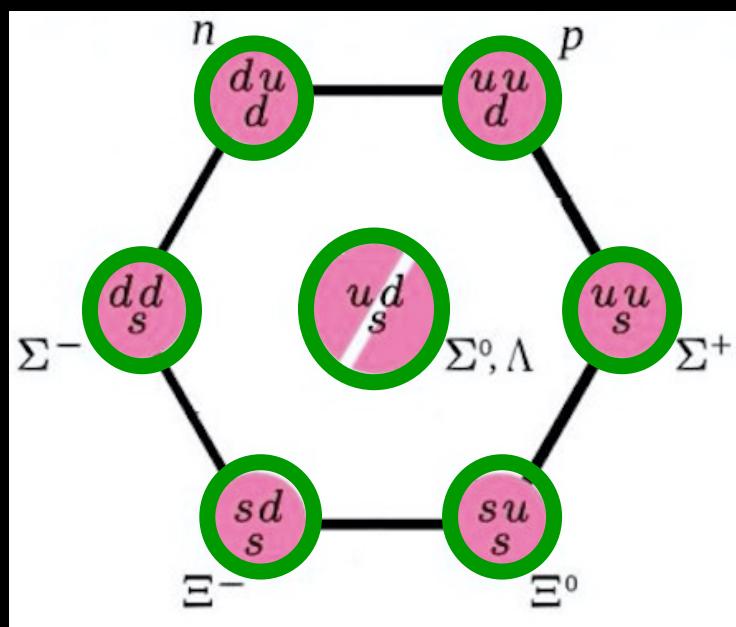
Example: Hyperon Forces from LQCD Hyperons in Nuclear Matter (LQCD+BHF)

HAL QCD Coll. (2016)

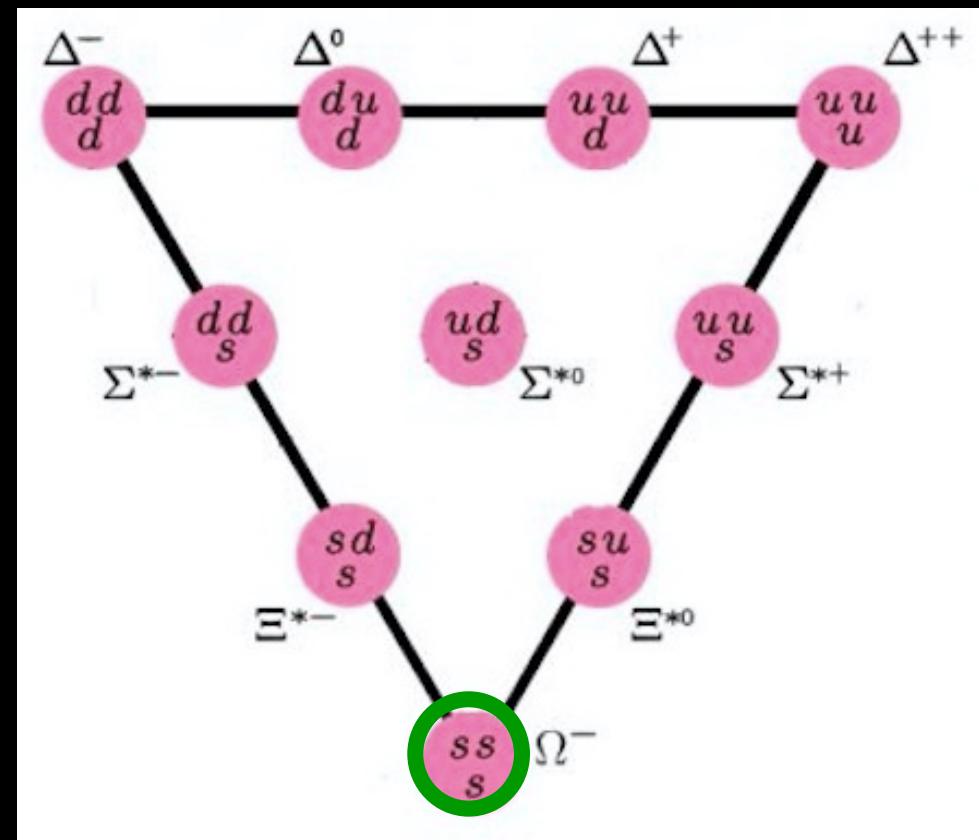


Flavor SU(3) Classification : single-baryon

8 (Octet)

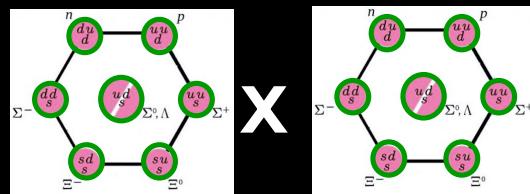


10 (Decuplet)



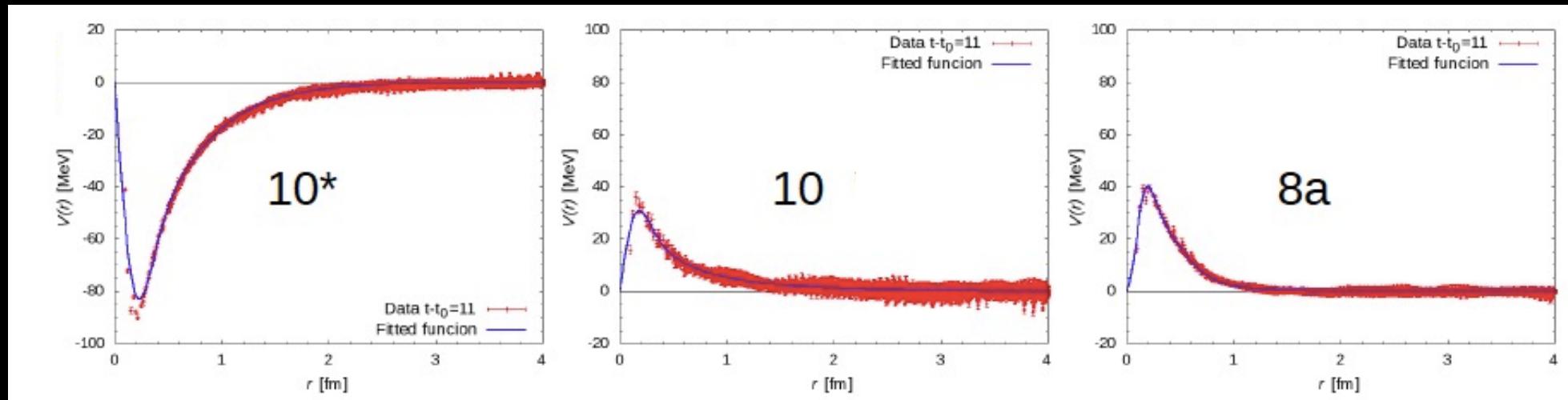
BB interactions in the flavor basis: $V_T(r)$

preliminary



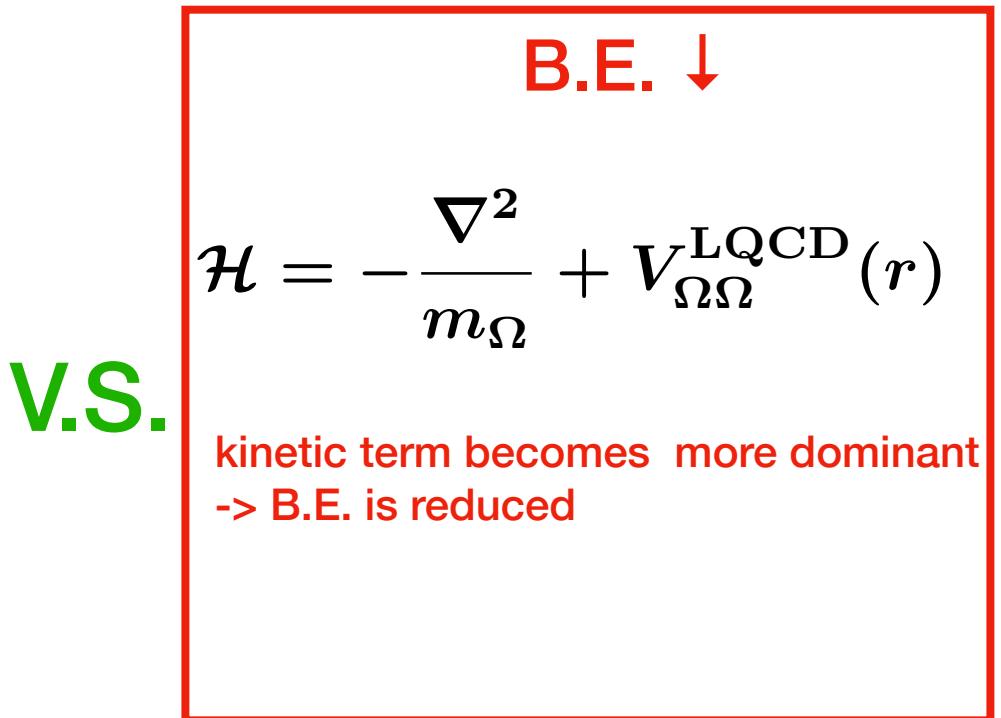
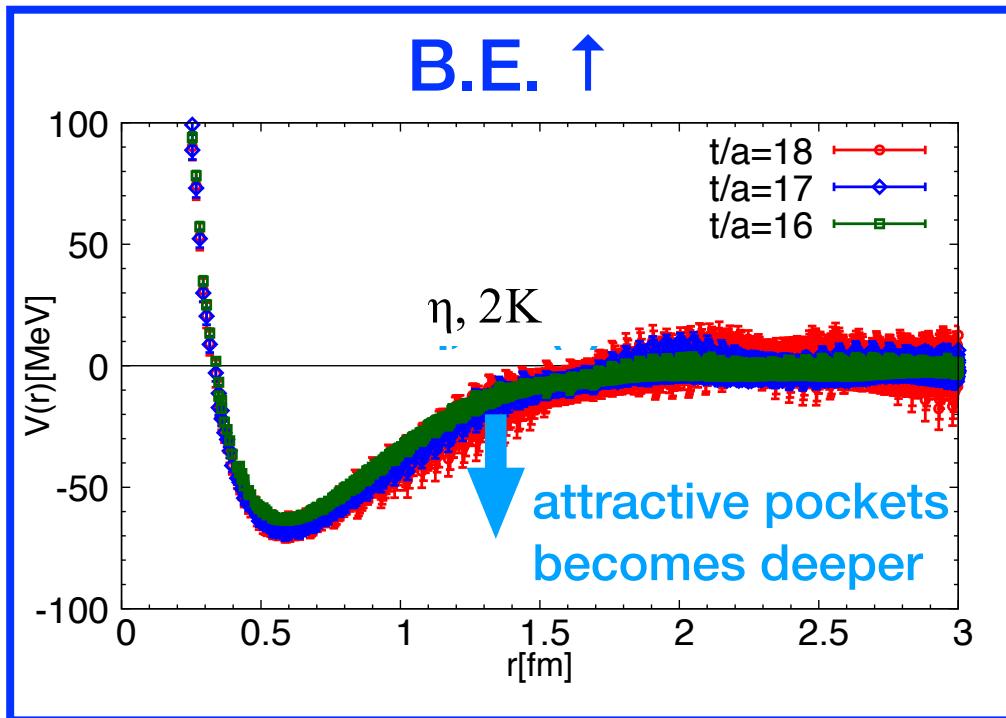
X

$$8 \times 8 = \frac{27 + 8s + 1}{^1S_0} + \frac{10^* + 10 + 8a}{^3S_1, ^3D_1}$$



Conservative estimate at exact phys. pt.

$m_{\pi}=146 \text{ MeV} \rightarrow 135 \text{ MeV}$, $m_{\Omega}=1712 \text{ MeV} \rightarrow 1672 \text{ MeV}$



conservative estimate:
only change the mass of schroedinger eq.

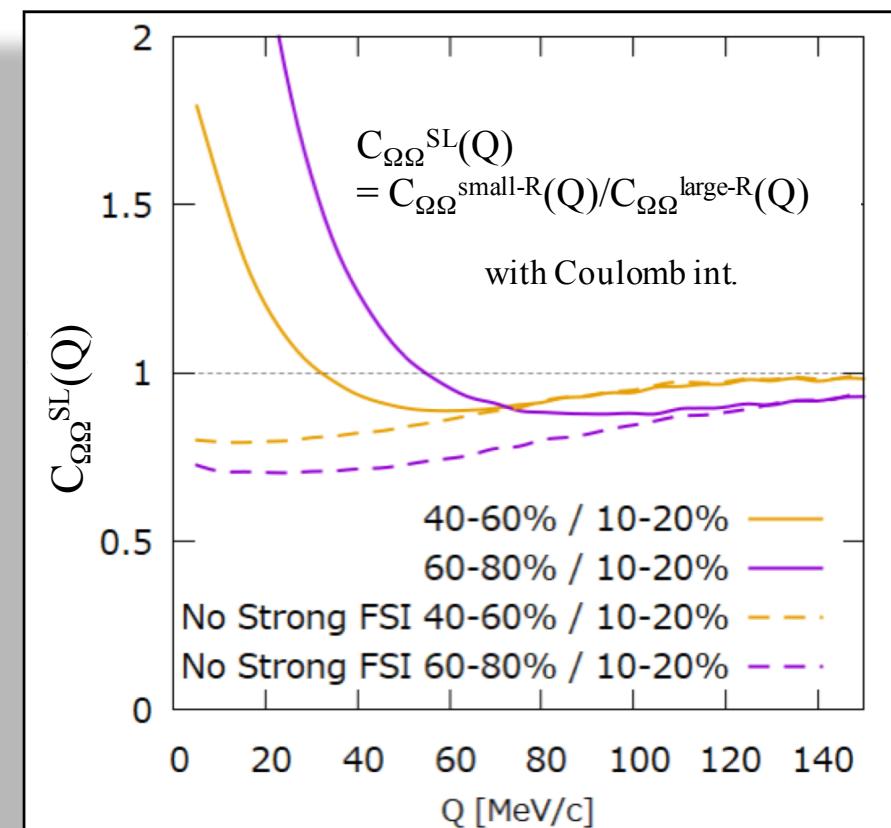
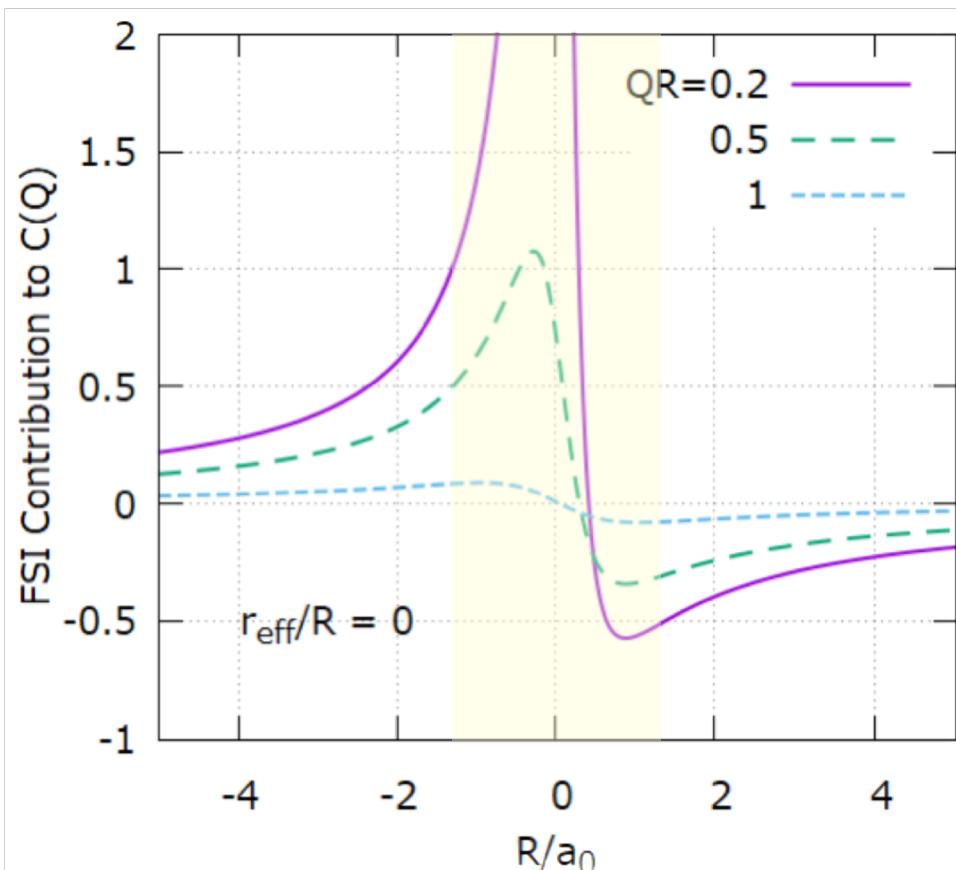
$$(B_{\Omega\Omega}^{(\text{QCD})}, B_{\Omega\Omega}^{(\text{QCD+Coulomb})}) = (1.6(6)\text{MeV}, 0.7(5)\text{MeV})$$
$$\rightarrow (1.3(5)\text{MeV}, 0.5(5)\text{MeV})$$

These changes are well within errors

Correlation from FSI

$$C_{AB}(Q) - 1 = \frac{4\pi}{(2\pi R^2)^3} \int dr r^2 S^{\text{rel}}(r) [|\chi_Q(r)|^2 - |j_0(Qr)|^2]$$

Lednicky+
1982



Morita+ (in preparation)

$\Upsilon=2$ States in $Su(6)$ Theory

Freeman J. Dyson and Nguyen-Huu Xuong

Phys. Rev. Lett. **13**, 815 – Published 28 December 1964; Erratum Phys. Rev. Lett. **14**, 339 (1965)

[From Dyson, Sep. 13, 2017]

Thank you very much for sending me your paper on the Omega-Omega calculation. This is a beautiful piece of work, and it will be a big step forward if the experimenters are able to confirm it.

Thank you also for referring to our 1964 paper. I am amazed that you remember that paper after 53 years. The predictions that we made in that paper turned out to be wrong, and the SU6 theory was soon abandoned. Luckily you did not assume SU6 symmetry when you made your prediction.

Now I wish you the joy of seeing it confirmed.

Yours sincerely,

Freeman Dyson.

