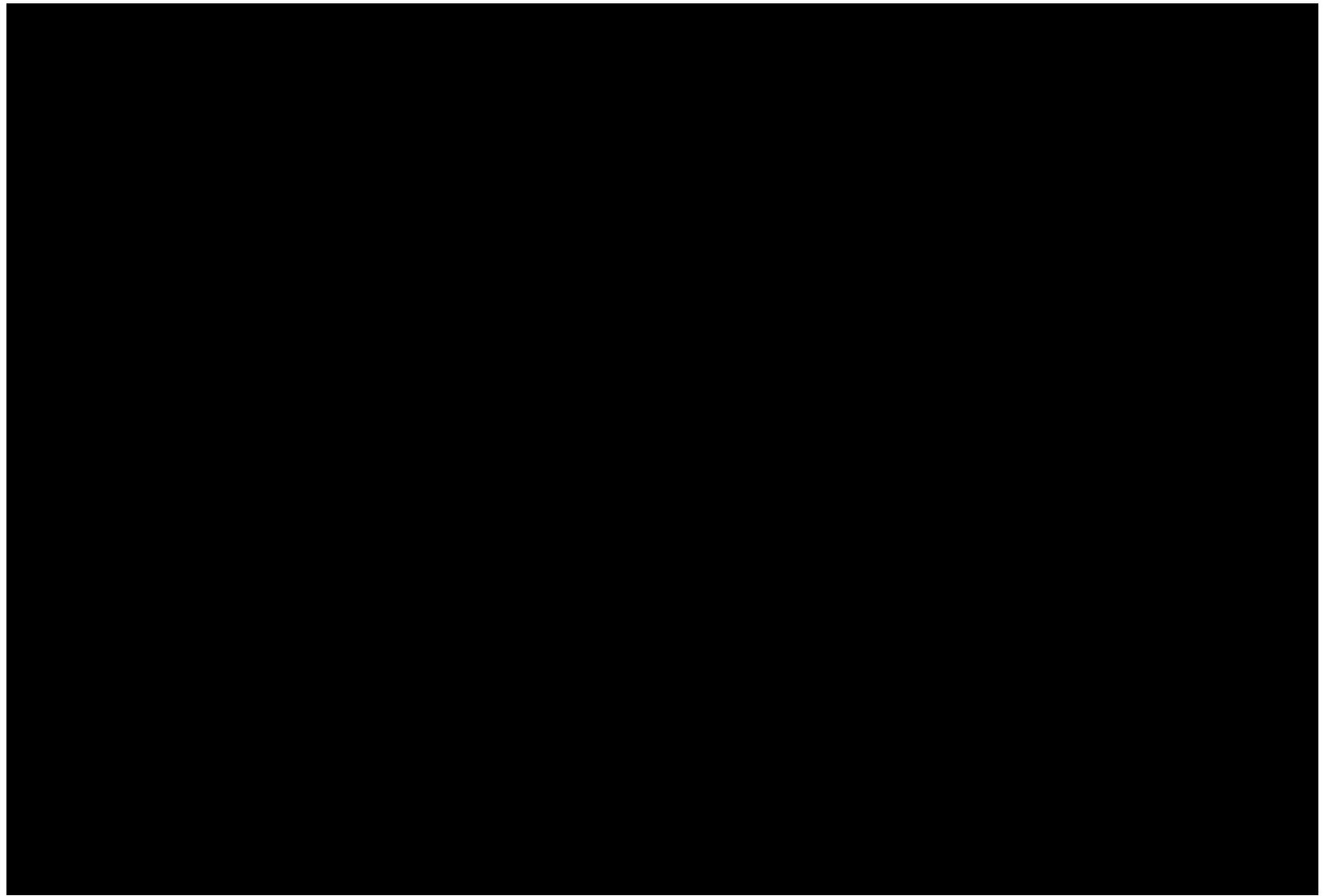


Effective Field Theory Analysis of CRESST-II Data

56th International Winter Meeting on Nuclear Physics

January 25, 2018, Vincent Schipperges





CRESST

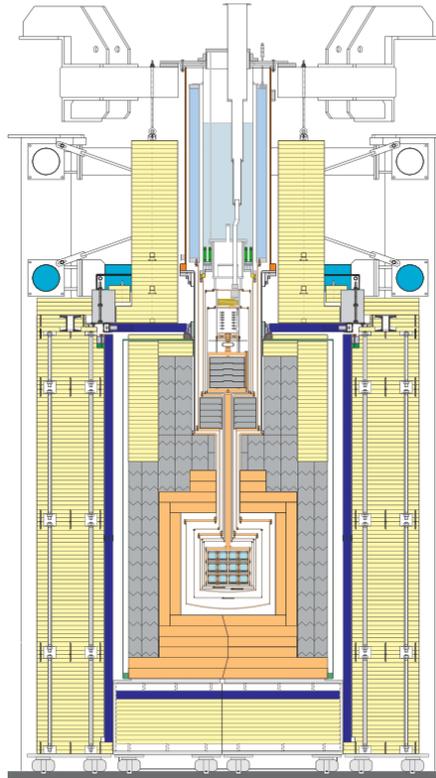
- direct Dark Matter detection experiment (model independent):
nuclear recoils
- Cryogenic Rare Event Search with Superconducting
Thermometers
- located at the LNGS underground lab in central Italy



CRESST

- additional shielding

- temperature ~ 15 mK





CRESST-II detectors

- scintillating target (CaWO_4)
- separate phonon detector and light detector (for scintillation light detection)

Example: detector *Lise*

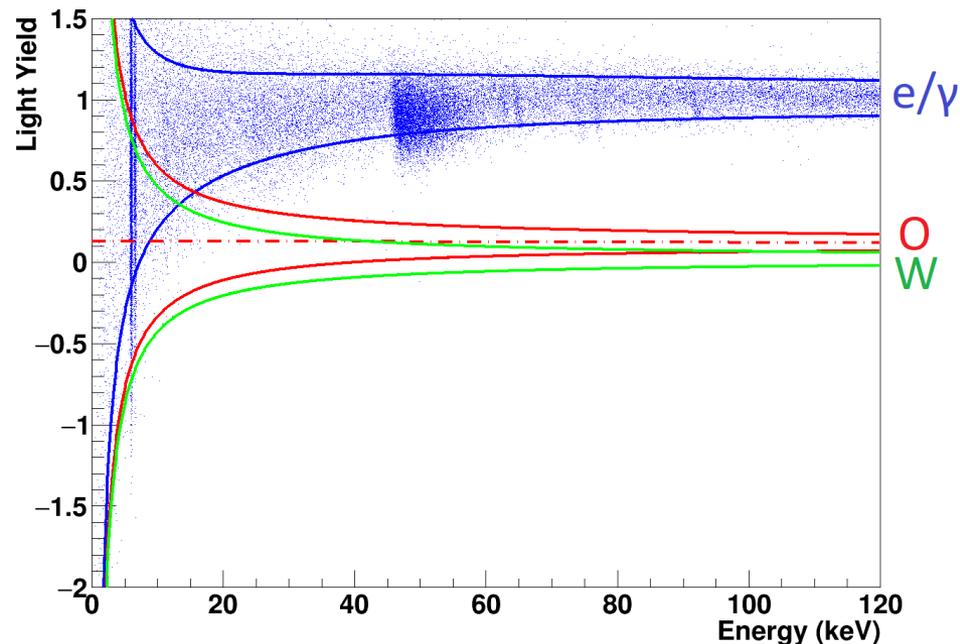
- $E_{th} = 307 \pm 4$ eV for nuclear recoils





CRESST Data Analysis

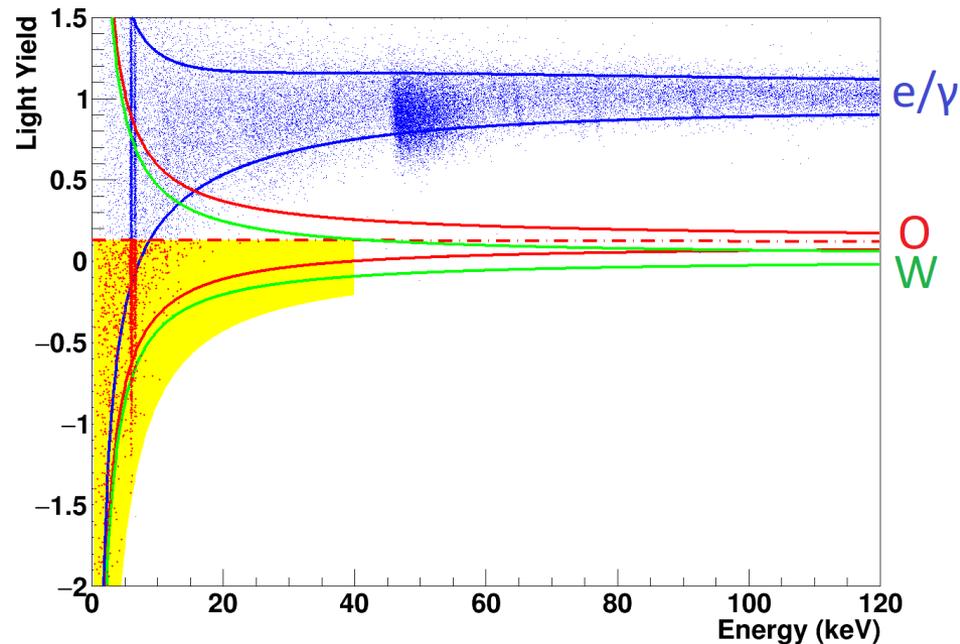
- signal: nuclear recoils, most background: electron recoils
- particle discrimination by fraction of light and heat (phonon)
signal: light yield





CRESST Data Analysis

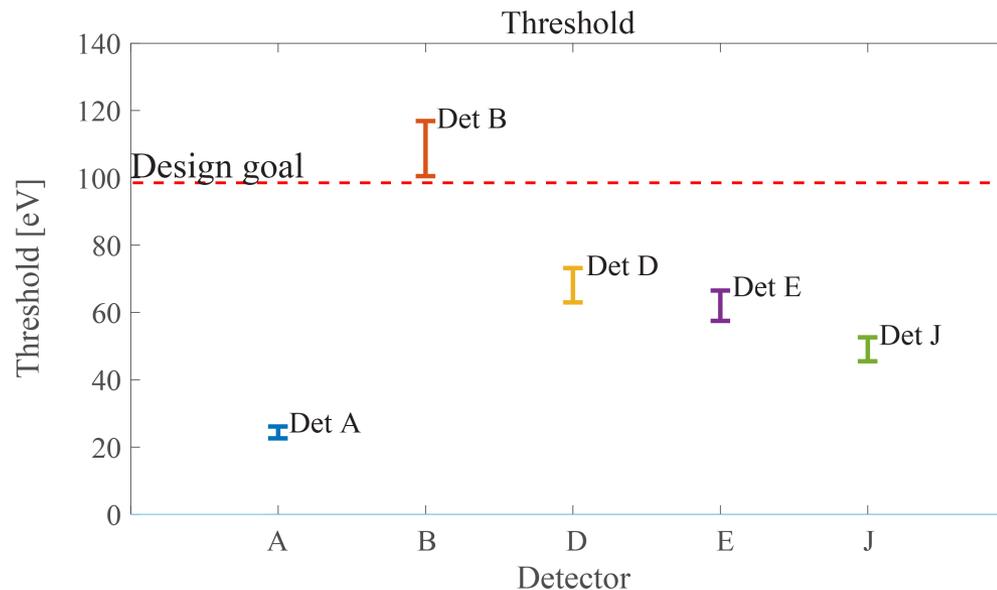
- signal: nuclear recoils, most background: electron recoils
- particle discrimination by fraction of heat (phonon) and light signal: light yield
- assuming all events in ROI as signal
- using Yellin optimum interval method for limit calculation





CRESST-III

- upgrade to CRESST-III, data taking since 2016 with 10 new detectors
- focus on low mass dark matter, detector design improvements
- design goal: energy threshold below 100 eV





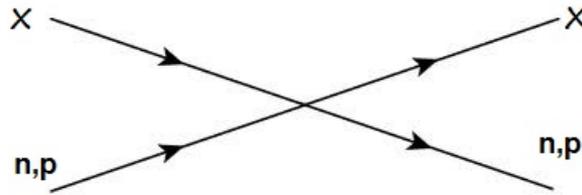
Dark Matter Search

- model-independent? in general assuming scalar (or spin-dependent) DM-nucleon interaction to be dominant
 - dark matter nature unknown
 - structure might be more complicated than assumed
- set limits in a more general way



Effective Field Theory of Dark Matter Detection

low momentum transfer (keV)



non-relativistic



Effective Field Theory of Dark Matter Detection

possible relevant Galilean invariant quantities:

$$\vec{S}_\chi, \quad \vec{S}_N, \quad \vec{q}, \quad \vec{v}$$

Hermitian and orthogonal:

$$\vec{S}_\chi, \quad \vec{S}_N, \quad i\vec{q}, \quad \vec{v}^\perp = \vec{v} + \vec{q}/(2\mu_N)$$



Dark-Matter-Nucleon Interaction

possible combinations:

$$\mathcal{O}_1 = \mathbb{I}$$

$$\mathcal{O}_3 = i\vec{S}_N \cdot (\vec{q} \times \vec{v}^\perp)$$

$$\mathcal{O}_4 = \vec{S}_\chi \cdot \vec{S}_N$$

$$\mathcal{O}_5 = i\vec{S}_\chi \cdot (\vec{q} \times \vec{v}^\perp)$$

$$\mathcal{O}_6 = (\vec{S}_\chi \cdot \vec{q})(\vec{S}_N \cdot \vec{q})$$

$$\mathcal{O}_7 = \vec{S}_N \cdot \vec{v}^\perp$$

$$\mathcal{O}_8 = \vec{S}_\chi \cdot \vec{v}^\perp$$

$$\mathcal{O}_9 = i\vec{S}_\chi \cdot (\vec{S}_N \times \vec{q})$$

$$\mathcal{O}_{10} = i\vec{S}_N \cdot \vec{q}$$

$$\mathcal{O}_{11} = i\vec{S}_\chi \cdot \vec{q}$$

$$\mathcal{O}_{12} = \vec{S}_\chi \cdot (\vec{S}_N \times \vec{v}^\perp)$$

$$\mathcal{O}_{13} = i(\vec{S}_\chi \cdot \vec{v}^\perp)(\vec{S}_N \cdot \vec{q})$$

$$\mathcal{O}_{14} = i(\vec{S}_\chi \cdot \vec{q})(\vec{S}_N \cdot \vec{v}^\perp)$$

$$\mathcal{O}_{15} = -(\vec{S}_\chi \cdot \vec{q})((\vec{S}_N \times \vec{v}^\perp) \cdot \vec{q})$$

interaction with nucleon described by $\sum_{\tau=0,1} \sum_{k=1}^{15} c_k^\tau \mathcal{O}_k t^\tau$



Dark-Matter-Nucleus Interaction

$$\hat{\mathcal{L}}(\vec{r}) = \sum_{i=1}^A \sum_{\tau=0,1} \sum_{k=1}^{15} c_k^\tau \mathcal{O}_k^{(i)} t_{(i)}^\tau$$

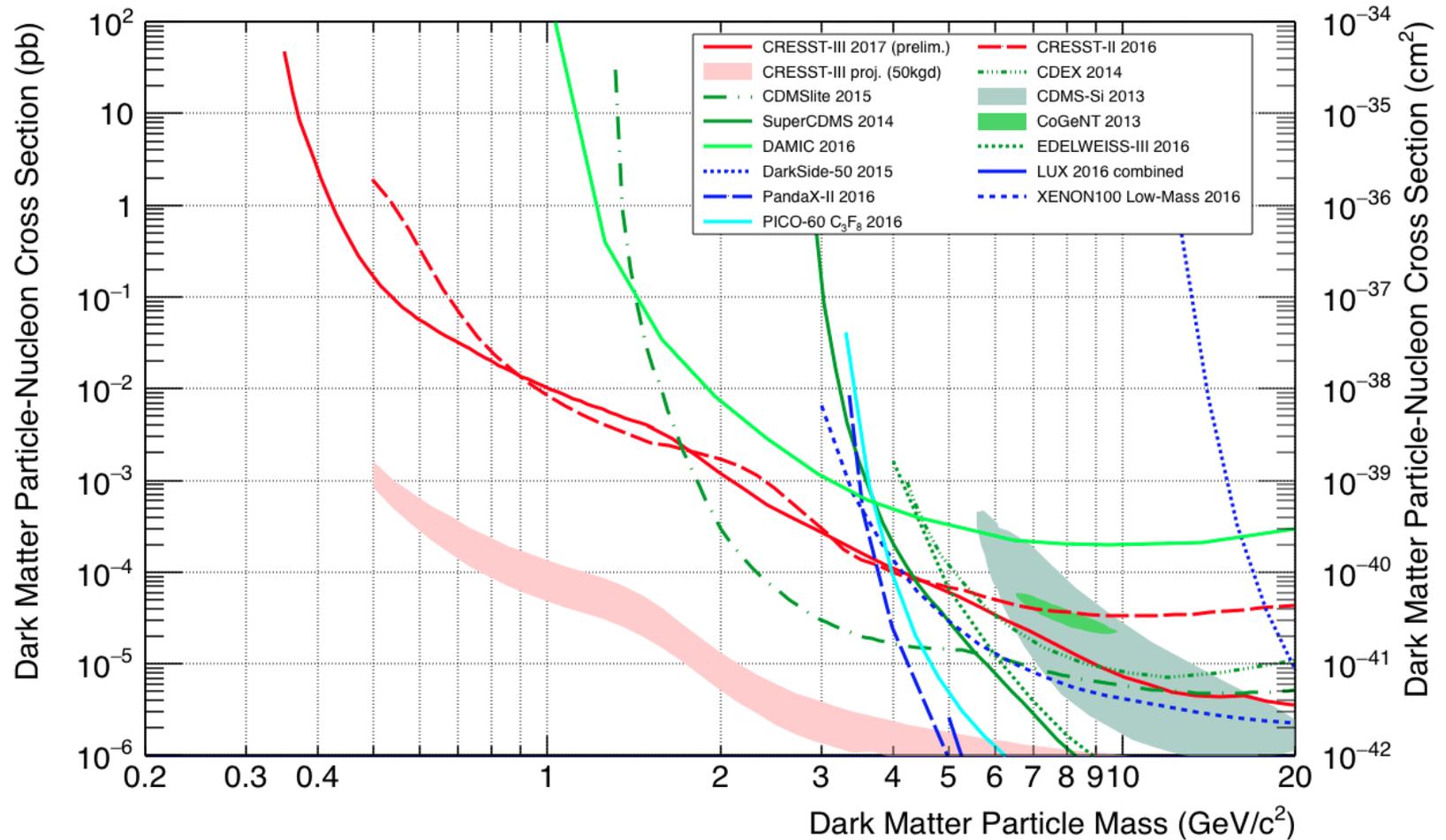
goal: set limits on c , for each one separately (setting all the other $c = 0$)

for movement inside nucleus: input from nuclear physics (shell model)

⇒ set limits (isoscalar operators not depending on nuclear spin)

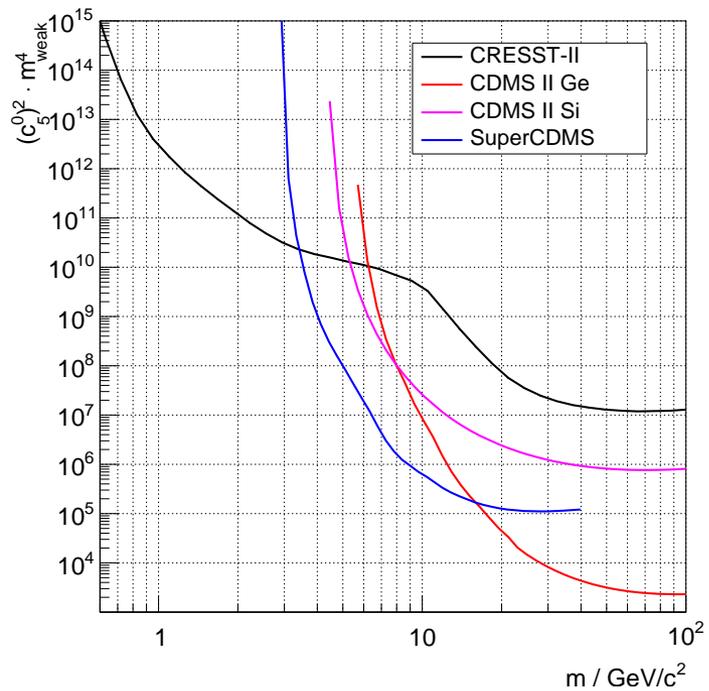
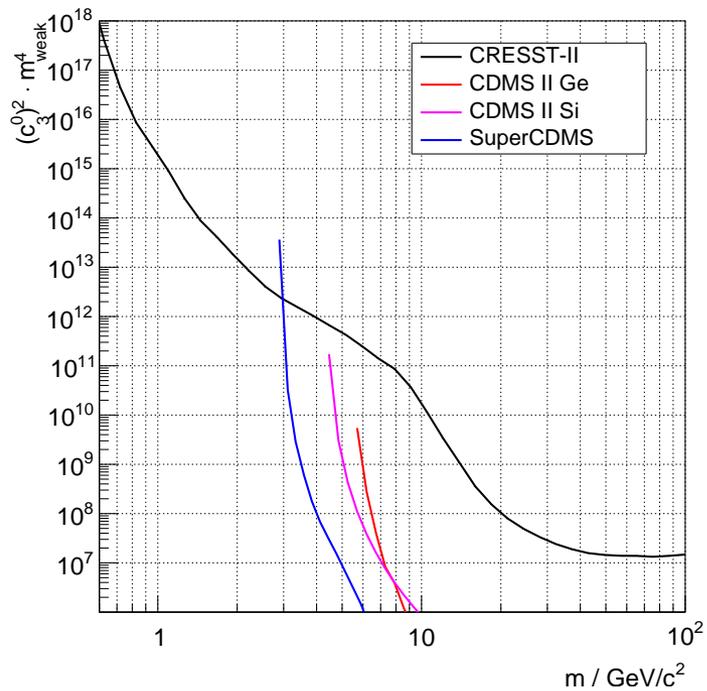


Results





Limits for Wilson coefficients c_3^0 and c_5^0





Summary

- set limits for more general Effective Field Theory
- CRESST-II: best DM limit for SI interaction below $1.7 \text{ GeV}/c^2$
- leading limits for masses below $\sim 3\text{-}4 \text{ GeV}/c^2$ for coefficients $c_1^0, c_3^0, c_5^0, c_8^0, c_{11}^0, c_{12}^0$ and c_{15}^0
- upgrade to CRESST-III: improved threshold, precision, design and radiopurity