



# Coherence in Neutrino Oscillations

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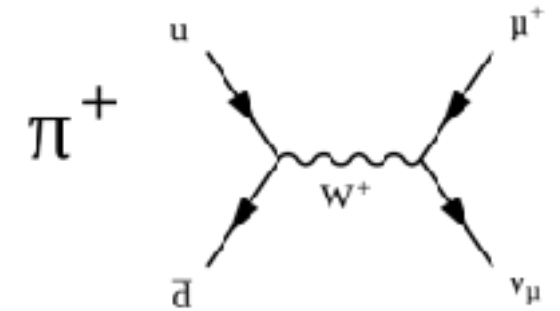
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# Pion Decay At Rest



Center of Mom.:  $p_\mu = p_\pi - p_\nu$

Square both sides:  $p_\mu^2 = p_\pi^2 + p_\nu^2 - 2p_\pi \cdot p_\nu$

$\mathbf{p}_\pi = \mathbf{0}$ :  $m_\mu^2 = m_\pi^2 + m_\nu^2 - 2m_\pi E_\nu$

$$E_\nu = \frac{m_\pi^2 + m_\nu^2 - m_\mu^2}{2m_\pi}$$

So there is an *intrinsic uncertainty* on neutrino energy corresponding to the *difference in mass eigenstates*

...but then how are these ever coherent?

$$\left\langle e^{-i(E_1 - E_2)t} \right\rangle_t = 0 \quad \text{unless } E_2 = E_1$$

“Only neutrino mass eigenstates with a common energy  $E$  are coherent” —Stodolsky

# Intrinsic Momentum Uncertainty

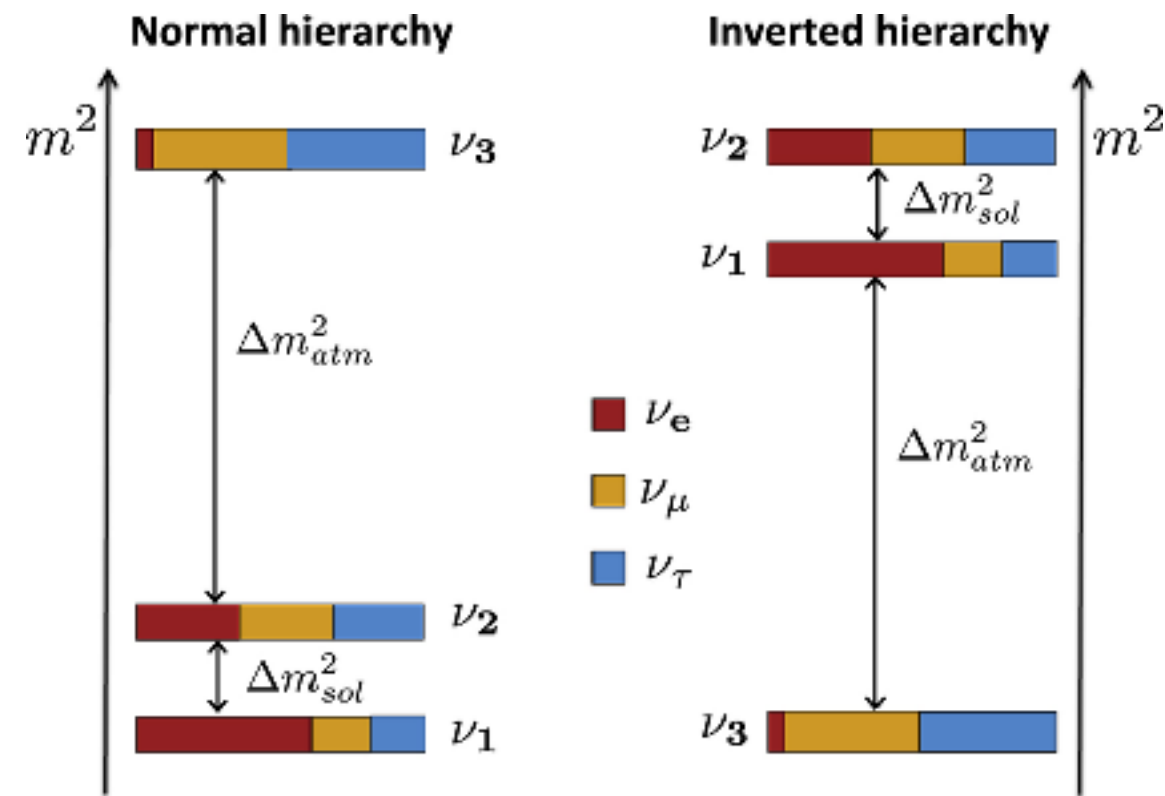
*of a superposition of mass-split eigenstates*

Momentum of a neutrino given its energy  $E$

$$p = (E^2 - m^2)^{\frac{1}{2}} \approx E - \frac{m^2}{2E}$$

Taylor Expand around small  $m^2/E^2$

Heisenberg:  $\Delta x \Delta p \geq \frac{1}{2}$



- Plug in Taylor-approximated momentum
- Atmospheric mass splitting dominates

$$\Delta p = \frac{m_{\max}^2 - m_{\min}^2}{2E} \approx \frac{\Delta m_{\text{atm}}^2}{2E}$$

$$\Delta x \geq \frac{E}{\Delta m_{\text{atm}}^2}$$

# Decoherence

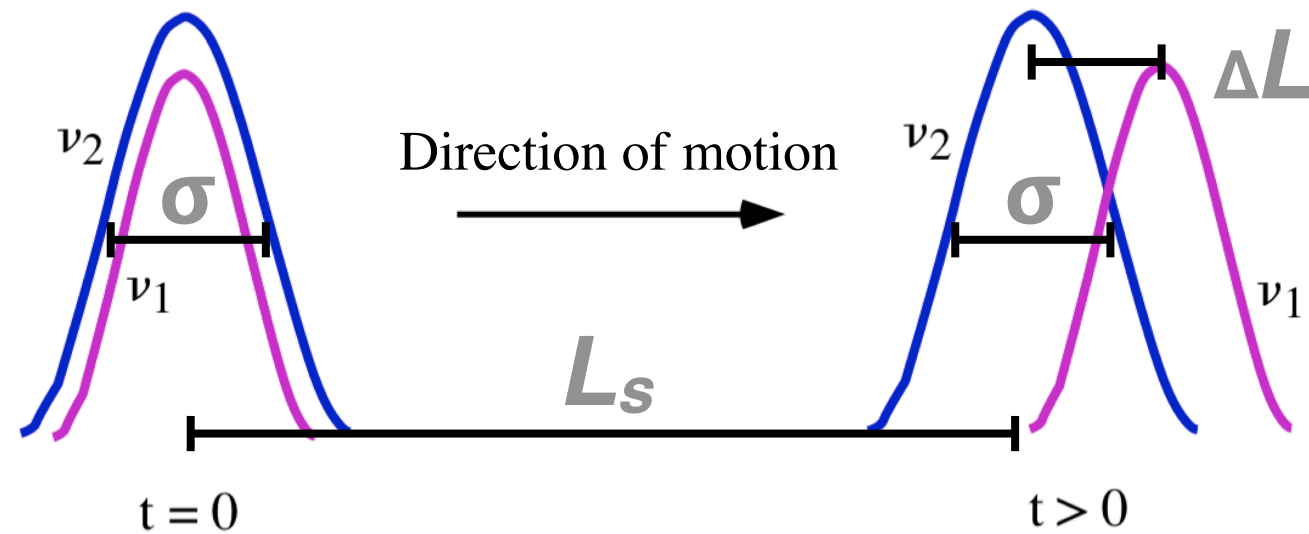
$$\Delta L = t\Delta v = \frac{L_s}{v_{\min}} \Delta v$$

$$v = \beta = \frac{p}{E} = \left(1 - \frac{m^2}{E^2}\right)^{\frac{1}{2}} \approx 1 - \frac{m^2}{2E^2}$$

$$\Delta L = L_s \left( \frac{2E^2}{2E^2 - m_j^2} \right) \left( \frac{m_j^2 - m_i^2}{2E^2} \right)$$

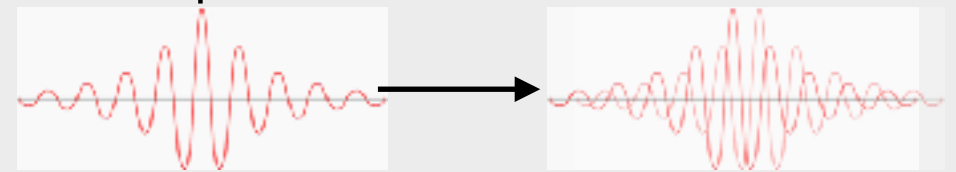
$$L_s = \Delta L \left( \frac{2E^2 - m_j^2}{\Delta m_{ij}^2} \right)$$

Estimating the width  $\sigma$  of the wave packet depends on the source. How precisely *can* you know the position the neutrino was created ( $\Delta x$ ). *Not how well do you know*



**It's much more subtle...**

Wave packets are more like:



What happens when  $\nu_3$  is incoherent while  $\nu_1$  and  $\nu_2$  are coherent? 1 and 2 still interfere, so they still oscillate. Writing that down is a matter of wave packet formalism.

<https://arxiv.org/abs/hep-ph/0109119>

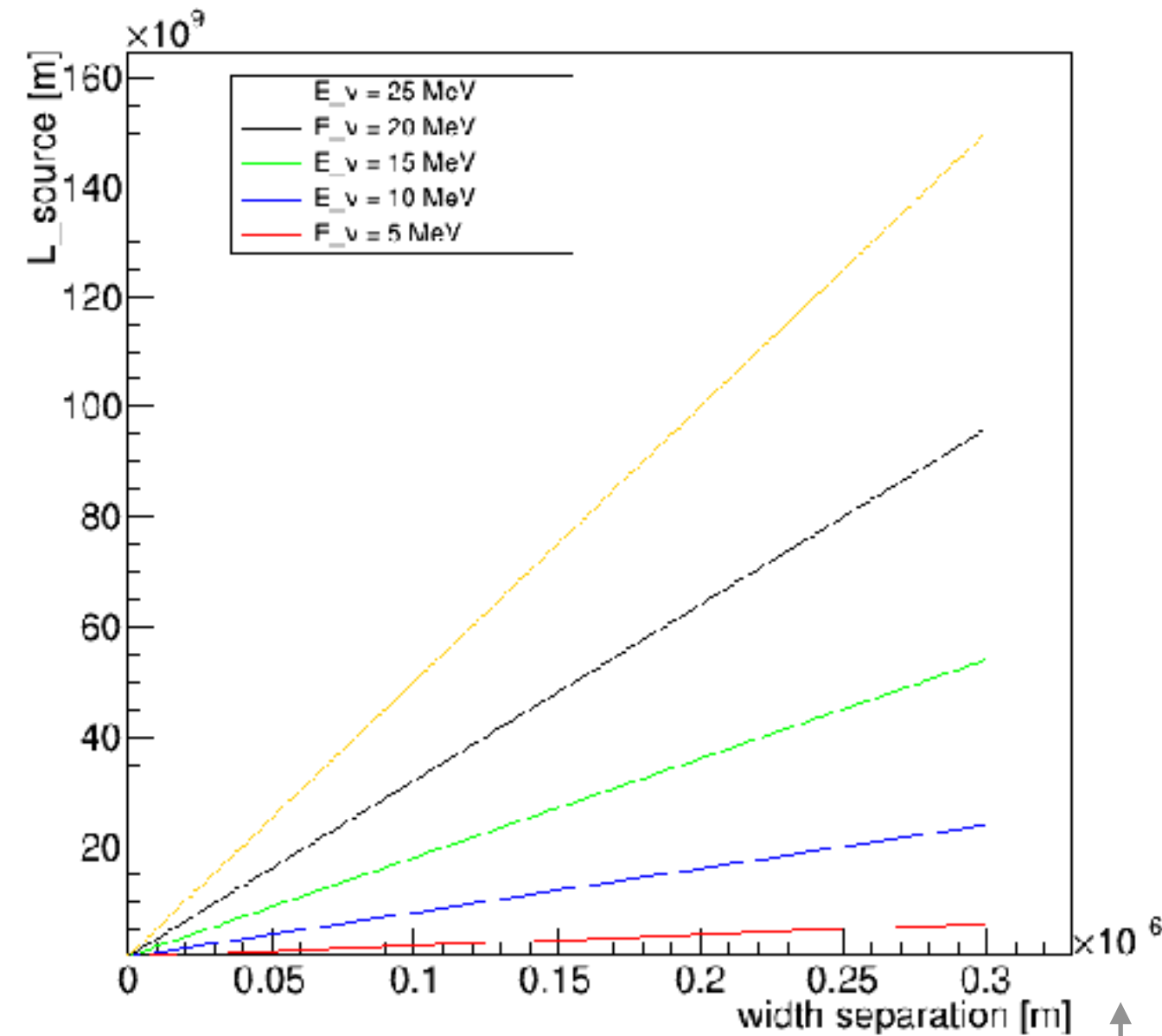
<https://arxiv.org/abs/1001.4815>

Define width  $\sigma$

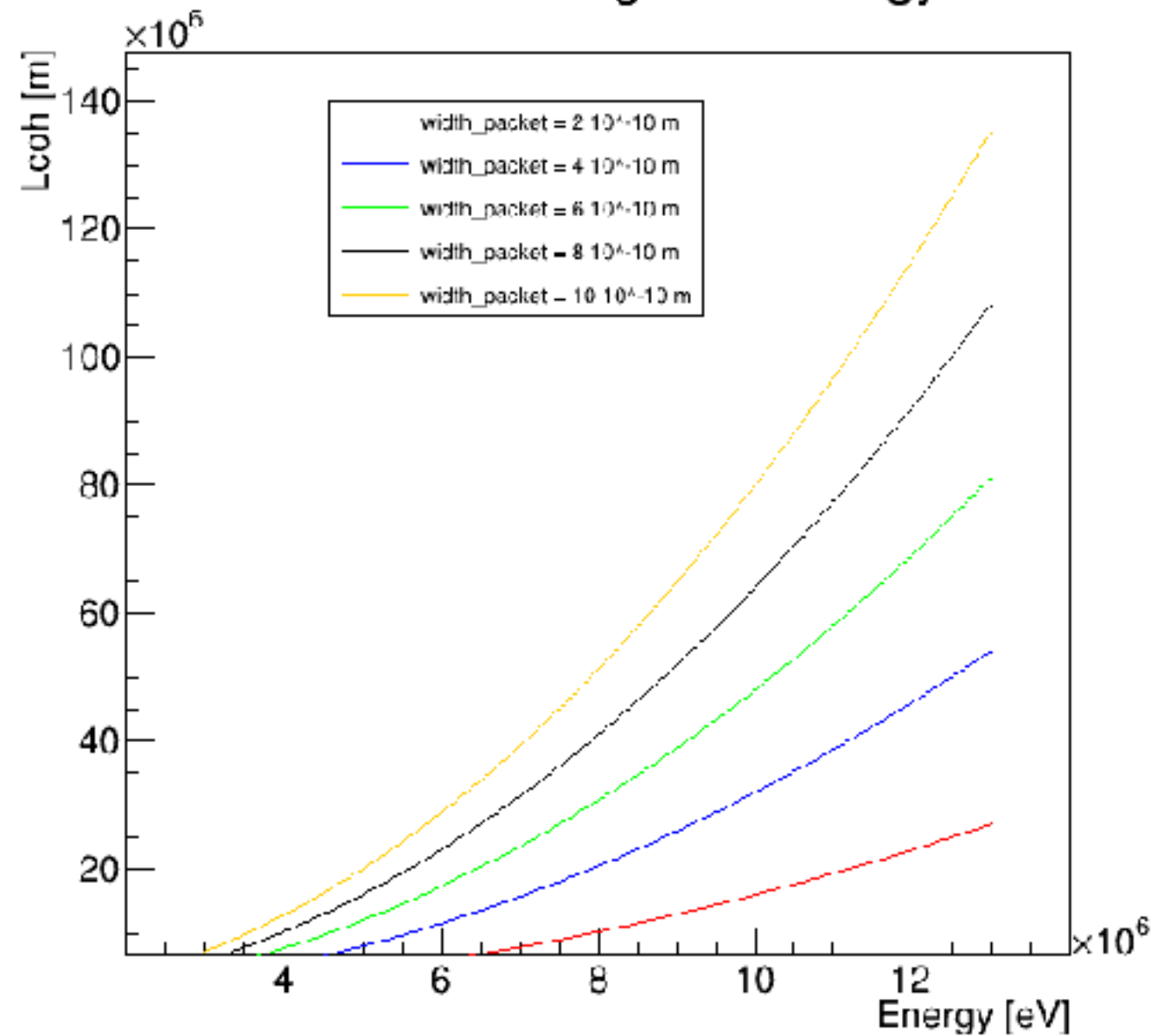
$$L_{\text{coh}} = \sigma \left( \frac{2E^2}{\Delta m_{ij}^2} \right)$$

# Decoherence

Length from source vs length separation of wave packet



Coherence length vs energy



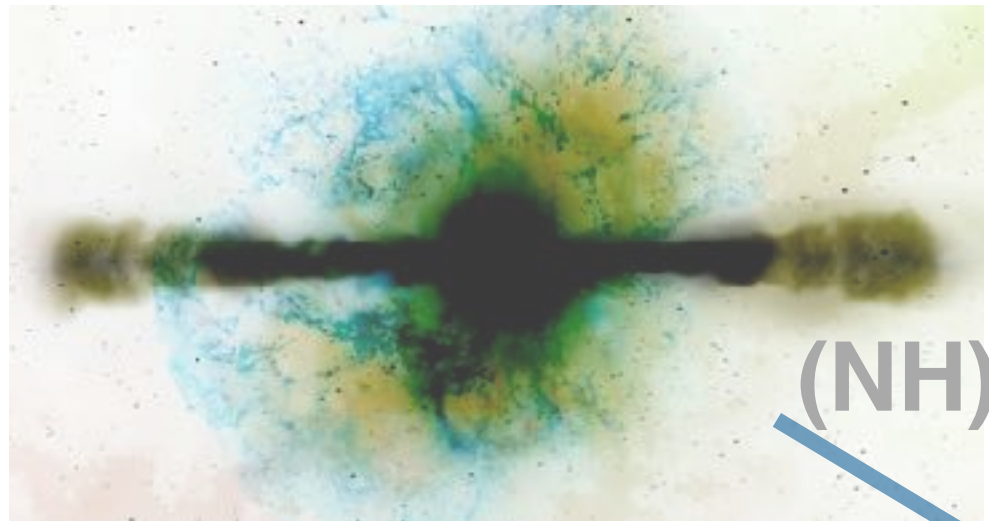
$$L_s = \Delta L \left( \frac{2E^2 - m_j^2}{\Delta m_{ij}^2} \right)$$

Nucleon density  
Atomic widths  
Length of decay pipe  
but can still pinpoint nu  
creation much more closely

$$L_{\text{coh}} = \sigma \left( \frac{2E^2}{\Delta m_{ij}^2} \right)$$

# Supernova Neutrinos

*arrive at Earth with incoherent mass states*



(NH)

$$|\langle \nu_e | \nu_2 \rangle|^2 = |U_{e2}|^2 = \frac{1}{3}$$

In *vacuum*, it's easiest to think  
of the time evolution  
in the *mass* eigenstate basis

$$\mathcal{A}_{\alpha \rightarrow \beta} = \langle \nu_\beta(0) | e^{-i(\hat{H}_{\text{matter}}t - \vec{p} \cdot \vec{x})} | \nu_\alpha(0) \rangle$$

$$\mathcal{A}_{\alpha \rightarrow \beta} = \langle \nu_\beta(0) | e^{-i(\hat{H}_{\text{vac}}t - \vec{p} \cdot \vec{x})} | \nu_\alpha(0) \rangle$$

In *matter*, it's easiest to think of  
the time evolution of the *mass* eigenstates  
in the *matter* eigenstate basis

incoherent mass eigenstates  
are then superpositions  
of the matter eigenstates

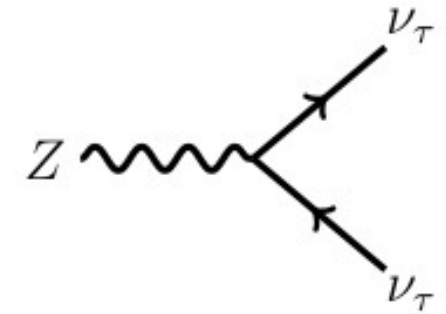
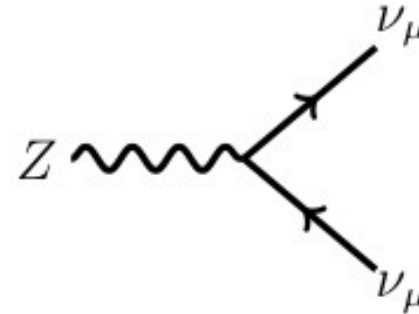
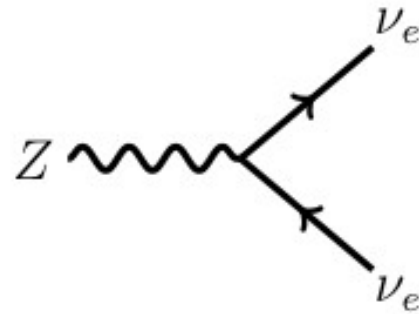
$\nu_2$

Begin coupling  
to a new matter  
Hamiltonian

$\nu_2$

# Do neutrinos from $Z^0$ decays oscillate ?

Flavor blind  $Z^0$  decay :



Production flux from  $Z^0$  decay :

$$\phi_e = \phi_\mu = \phi_\tau = \phi^0$$

General misconception : the observed neutrinos are decoherent

Observed flux for a given flavor  $\beta$  :

$$\phi_\beta(r) = \sum_{\alpha} \phi_\alpha P_{\alpha\beta}(r)$$

$$\Rightarrow \phi_\beta(r) = \phi^0 \sum_{\alpha} P_{\alpha\beta}(r) = \phi^0$$

Flux for one given flavor independent of distance  $\rightarrow$  no oscillation process observed



# What would it take to observe these oscillation ?

Coherence  $\Rightarrow$  We write the neutrino state as a combination of the mass eigenstates

$$\Rightarrow |\nu_z(r, \bar{r})\rangle = \frac{1}{\sqrt{3}} \sum_{1,2,3} |\nu_i\rangle |\bar{\nu}_i\rangle \boxed{e^{i\phi_i(r, \bar{r})}}$$

↓  
Phase factor

Probability to observe a  $\nu_\alpha$  at distance  $r$  and a  $\bar{\nu}_\beta$  at distance  $\bar{r}$  :

$$W_{\alpha\beta} = \frac{1}{3} \sum_i |U_{i\alpha}|^2 |U_{i\beta}|^2 + \frac{2}{3} \sum_{i>j} |U_{i\alpha}^* U_{i\beta} U_{j\alpha} U_{j\beta}^*| \cos(\phi_i - \phi_j - \xi_{\alpha\beta ij})$$

Oscillation factor

Probability to observe a  $\nu_\alpha$  :

$$W_\alpha = \sum_\beta W_{\alpha\beta} = \frac{1}{3}$$

We need to look both the neutrino and antineutrino to actually observe the oscillations



# Thank You



# Safe Travels!