A brief update on Nambu-Goto cosmic string models and the stochastic GW background from cosmic string loops

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long strings loose energy in the form of smaller loops + GWs





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 \Longrightarrow stochastic GW background: **GW frequencies are multiples of** $\,\omega=4\pi/l$

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Constraints on Cosmic Strings from the LIGO-Virgo Gravitational-Wave Detectors.



- loops assumed to be formed with tiny size (fraction of horizon size), decay in a Hubble time

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power law shape of loop distribution (through NG simulations)

Ringeval, Sakellariadou, Bouchet (2007) Lorentz, Ringeval, Sakellariadou (2010) Blanco-Pillado, Olum, Shlaer (2014)

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○ GW emission

 $P_{
m GW} = \Gamma G \mu^2$ $\Gamma \sim \mathcal{O}(50)$

Assumes about 1 kink/cusp per loop

GW evaporation dominates for loops of size

 $l < \Gamma G \mu t$

max scale to trust simulations



GW emission length scale in units of t

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- GW emission
- $_{
 m \circ}$ GW back-reaction $~l_{
 m c} < l_{
 m d}$

GW back-reaction scale in units of t

$$\gamma_{\rm c} = \Upsilon(G\mu)^{1+2\chi}$$
 where $\Upsilon \sim 10$ and $\chi = 1 - P/2$

$$P = 1.41^{+0.08}_{-0.07} |_{\text{mat}}, P = 1.60^{+0.21}_{-0.15} |_{\text{rad}}$$

Polchinski, Rocha (2006) Dubath, Polchinsk, Rocha (2018) Ringeval, Sakellariadou, Bouchet (2007) Lorentz, Ringeval, Sakellariadou (2010)

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- GW emission
- GW back-reaction
- GW signal depends on number of cusps and kinks

Not known from simulations

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Damour, Vilenkin (2001)

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Blanco-Pillado, Olum, Shlaer (2015) Wachter, Olum (2017) Wachter, Olum (2017)

Ringeval, Suyama 1709.03845

Lorentz, Ringeval, Sakellariadou (2010) Ringeval, Sakellariadou, Bouchet (2007)

$$\frac{\rho_{\rm gw}}{\rho_{\rm crit}} \equiv \int_0^{+\infty} \frac{\mathrm{d}\omega}{\omega} \Omega_{\rm gw}(\omega)$$

Distinction between Individually separable events (bursts) and the stochastic background is done (as usually) through the value of $|h_{\mu\nu}(\omega)|^2$

$$\hat{\Omega}_{\rm sgw}(\omega) \equiv \frac{\Omega_{\rm sgw}(\omega)}{c_{\alpha}^2}$$

 Γ may change...

Numerical constant that contains all theoretical uncertainties associated with type of GW source (cusps, kinks, kink-kink collisions) amplitude decays as $[\omega^{-2}]$ but #events per loop

$$c_{3} \equiv \frac{\eta_{\rm s}^2 \sqrt{2}}{(2\pi\beta)^{2/3}}, \qquad c_{3/2} \equiv \frac{2\eta_{\rm s} \sqrt{2v_{\pm}^2}}{(2\pi\beta)^{1/3}}, \qquad c_{1} \equiv 4\sqrt{2v_{\pm}^2 v_{-}^2} \qquad \stackrel{\text{oscillations goes as square of $\#$ kinks}}{\eta_{\rm s} \simeq 4.1}, \qquad \beta = \mathcal{O}(1)$$

Ringeval, Suyama 1709.03845



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 $G\mu = 10^{-7}$ 10⁻⁵ peak 10⁻⁶ knee Produced by loops thermal history 1 cusp 10⁻⁷ *in the matter era* 1 kink 1 collision 10⁻⁸ 10⁻⁹ 10⁻¹⁰ 10⁻¹¹ High-frequency plateau produced $^{\rm Sgw}$ by loops in the radiation era 10⁻¹² 10⁻¹³ -adiation era contribution 10⁻¹⁴ matter era contribution 10⁻¹⁵ 10⁻¹⁶ 10⁻¹⁷ 10⁻¹⁸ 10⁻¹⁹ 100 1..... $10^{-16} 10^{-14} 10^{-12} 10^{-10} 10^{-8} 10^{-6} 10^{-4} 10^{-2} 10^{0} 10^{2} 10^{4} 10^{6} 10^{8} 10^{10} 10^{12} 10^{14} 10^{16}$ $f\left(Hz\right)$

Ringeval, Suyama 1709.03845

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The overall peak is significantly higher than the knee while its amplitude decreases with Gµ much slower than the amplitude at the knee frequency.



Ringeval, Suyama 1709.03845





All studies find a plateau at high frequencies, a maximum and a fast decay at low frequencies.

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VS.

Ringeval, Suyama 1709.03845

Lorentz, Ringeval, Sakellariadou (2010) Ringeval, Sakellariadou, Bouchet (2007)

 $\gamma_{\rm c} \neq \gamma_{\rm d}$

Blanco-Pillado, Olum, Shlaer (2014)

$$\gamma_{\rm c} = \gamma_{\rm d} = \Gamma G \mu$$

No thermal history effects Cusp events only

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Blanco-Pillado, Olum, Shlaer (2014)

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The overall peak is significantly higher than the knee while its amplitude decreases with $G\mu$ much slower than the amplitude at the knee frequency.

Also the effect of thermal history lowers the plateau by 3 orders or magnitude around LIGO frequencies.

Effects of microstructure



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Model	LIGO	EPTA	LIGO + EPTA
$2\mathrm{C}$	$G\mu \leq 1.1 \times 10^{-10}$	$\overline{G\mu} \leq 3.4 \times 10^{-11}$	$G\mu \leq 1.0 \times 10^{-11}$
LNK	—	$G\mu \leq 6.8 \times 10^{-11}$	$G\mu \leq 7.2 \times 10^{-11}$
HNK	$G\mu \leq 8.8 \times 10^{-14}$	$G\mu \leq 6.4 \times 10^{-12}$	$G\mu \leq 6.7 \times 10^{-14}$

Blanco-Pillado, Olum, Siemens 1709.02434

Blanco-Pillado, Olum (2017) Blanco-Pillado, Olum, Shlaer (2014)

Blanco-Pillado, Olum, Siemens 1709.02434

Blanco-Pillado, Olum (2017) Blanco-Pillado, Olum, Shlaer (2014)

• Thermal history is taken into account

$$H(z) = H_0 \sqrt{\Omega_{\Lambda} + (1+z)^3 \Omega_m + G(z)(1+z)^4 \Omega_r}$$
$$G(z) = \frac{T(z)^4 g_*(z)}{T_0^4 (1+z)^4 g_{*,0}}$$

- Gravitational back-reaction is considered through a toy model of smoothing
 -- kinks are not considered (smoothed out by `convolution')
- Assumed flat background
- Find cusps and compute radiation power spectrum

Blanco-Pillado, Olum (2017)



Pulsar Timing Array:

$$G\mu < 1.5 \times 10^{-11}$$

