

etc.

Axions ~~and defects~~ in the early Universe

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MAX-PLANCK-GESELLSCHAFT

MITP Scientific Program

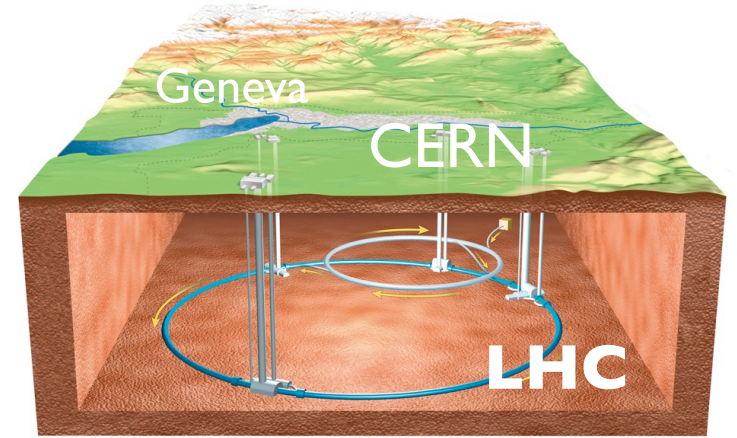
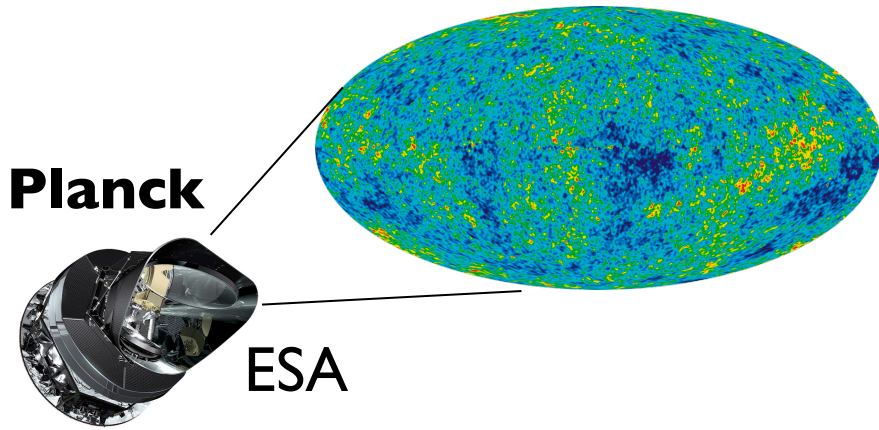
Jets, particle production and transport properties
in collider and cosmological environments

Mainz, July 28, 2014

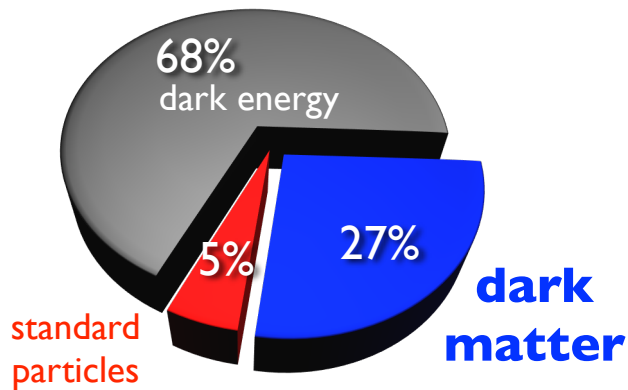
Astroparticle Physics

Cosmology
largest scales

Particle Physics
smallest scales

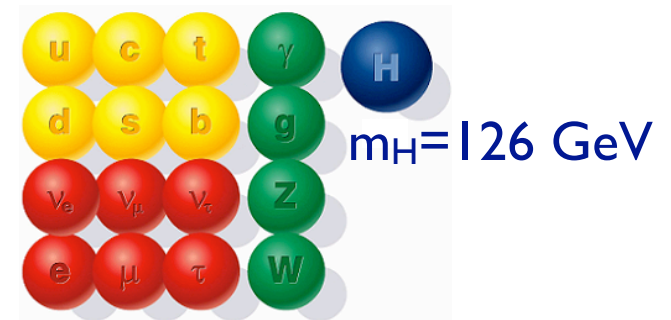


What is the energy budget of the Universe?



What is the particle identity and origin of dark matter?

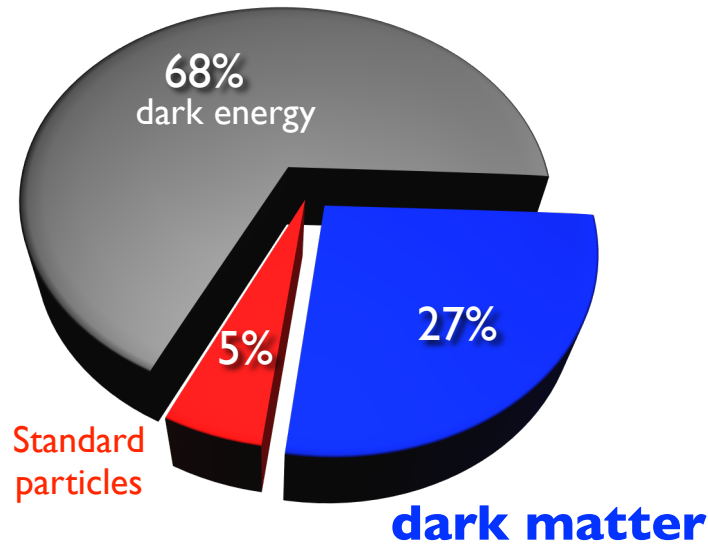
What are the fundamental constituents of matter?



Standard Model

Cosmology

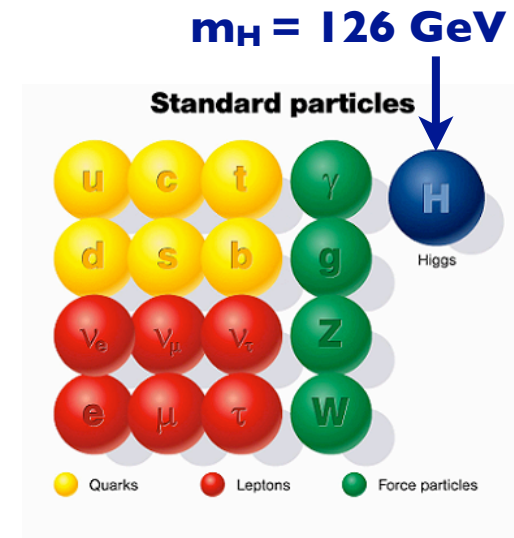
- 2013: Planck CMB sky map



- Cosmological puzzles
 - ? Matter-Antimatter Asymmetry
 - ? Particle Identity & Origin of Dark Matter
 - ? Dark Energy = Cosmological Constant

Particle Physics

- 2012: LHC Higgs-boson discovery

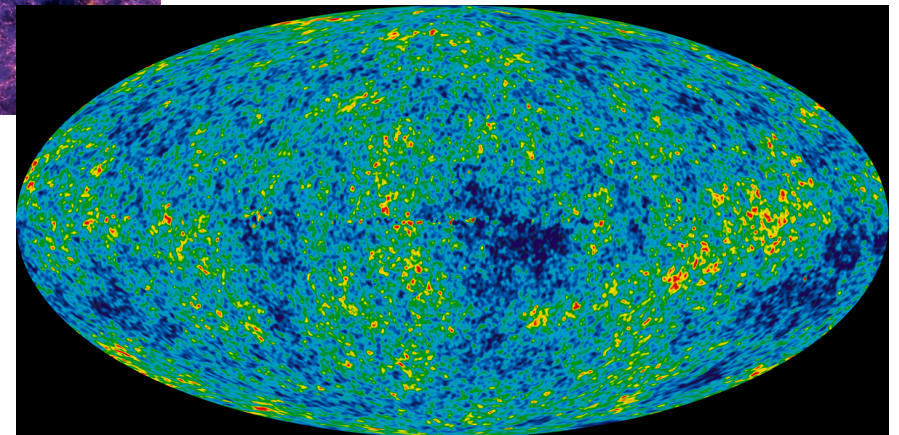
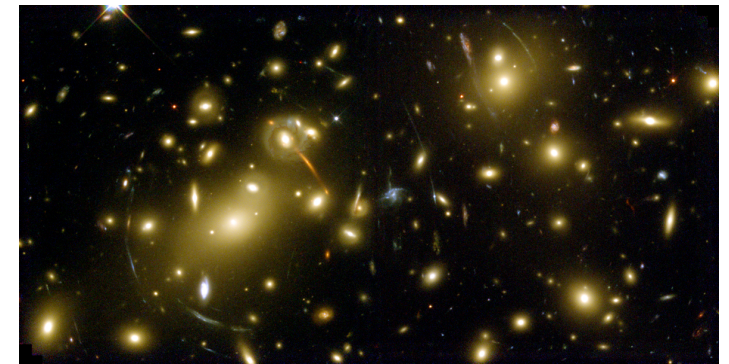
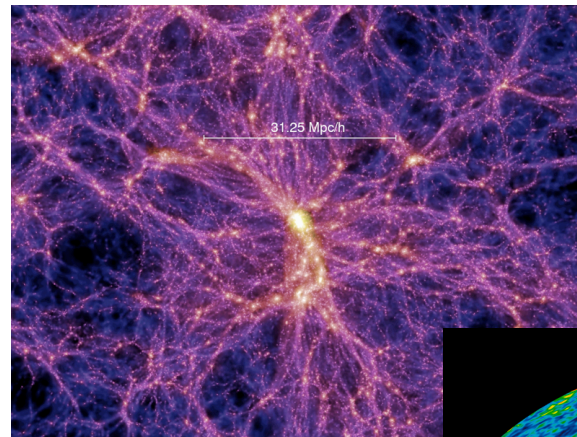
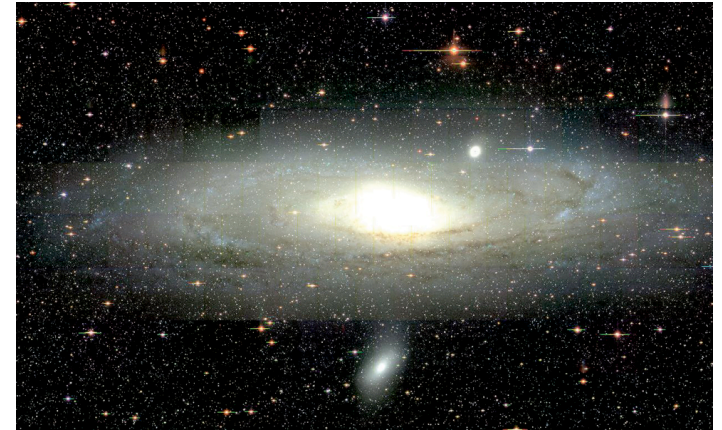


- Intrinsic fine tuning problems
 - ? Hierarchy Problem ($m_H \ll M_{\text{Planck}}$)
 - ? Strong CP Problem ($\Theta_{\text{QCD}} \ll 1$)
 - ? Small Neutrino Masses ($m_\nu \ll m_H$)

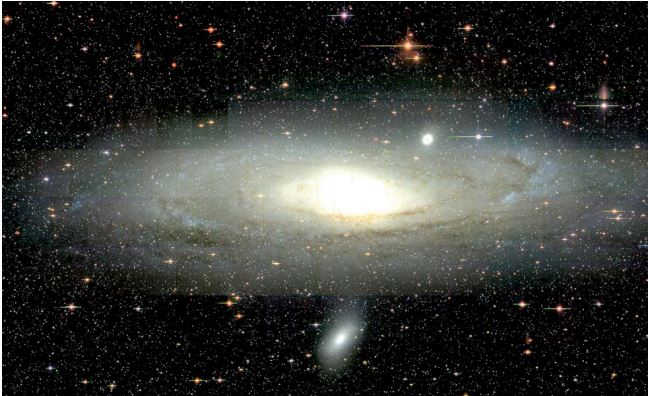
→ **Physics beyond the Standard Model**

Properties of Dark Matter

- stable or lifetime well above the age of our Universe
- electrically neutral
- clusters →
- cold / warm
- dissipationless
- color neutral



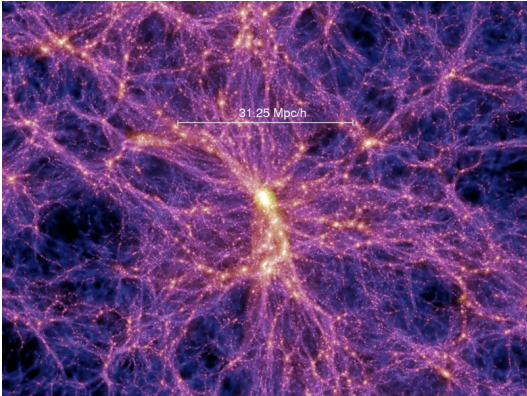
Dark matter



galaxies - rotation velocities

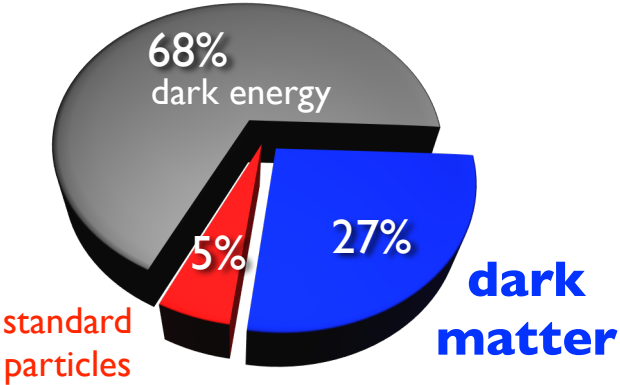


galaxy clusters - gravitational lensing



large scale structure

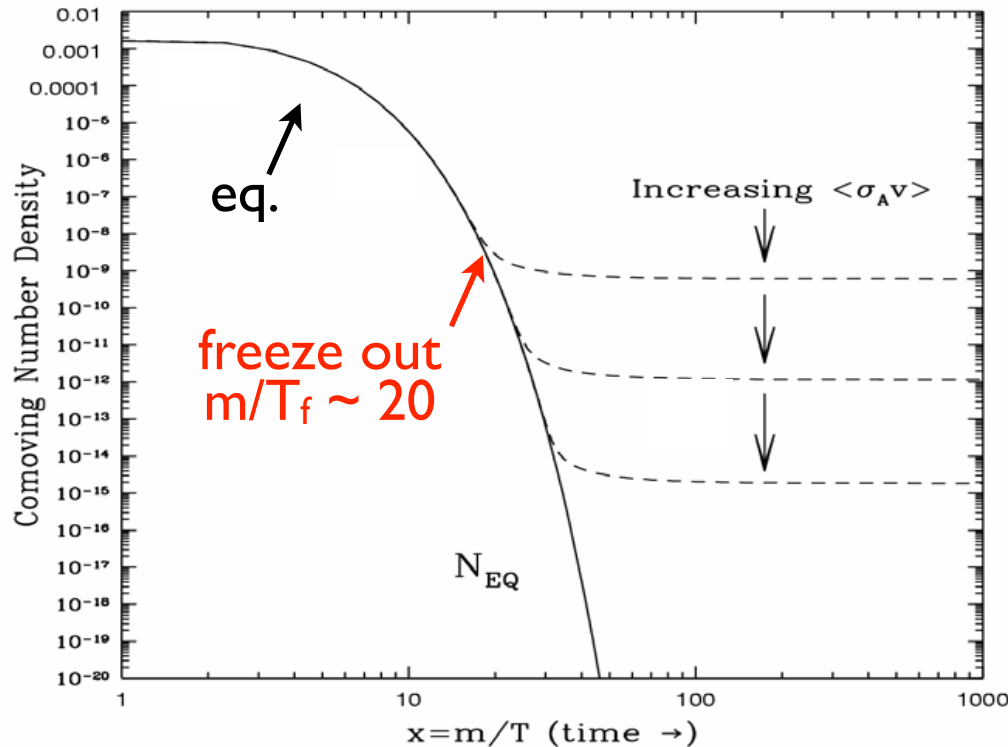
→ Particle Identity of Dark Matter



Axions etc.

Dark Matter Candidates

WIMP - Weakly Interacting Massive Particle



WIMP miracle

$$\Omega_{DM} \approx \frac{0.2 \text{ pb}}{\sigma_{\text{anni}}}$$

$$\sigma_{\text{anni}} \approx 1 \text{ pb}$$

Fermi-scale annihilation cross section

EWIP - Extremely Weakly Interacting Particle

thermally produced EWIP DM

← QCD processes
in the hot early Universe

axion condensate

Cosmic Relic Abundances

reheating temp.



decoupling temp. of X



- $T_R > T_D: 1+2 \rightleftharpoons 3+X$

$T > T_D: X$ in thermal eq. with the primordial plasma

$T \sim T_D: X$ decouples as a **thermal relic** (\rightarrow B. eq.)

- $T_R > T_D: 1+2 \rightarrow 3+X$

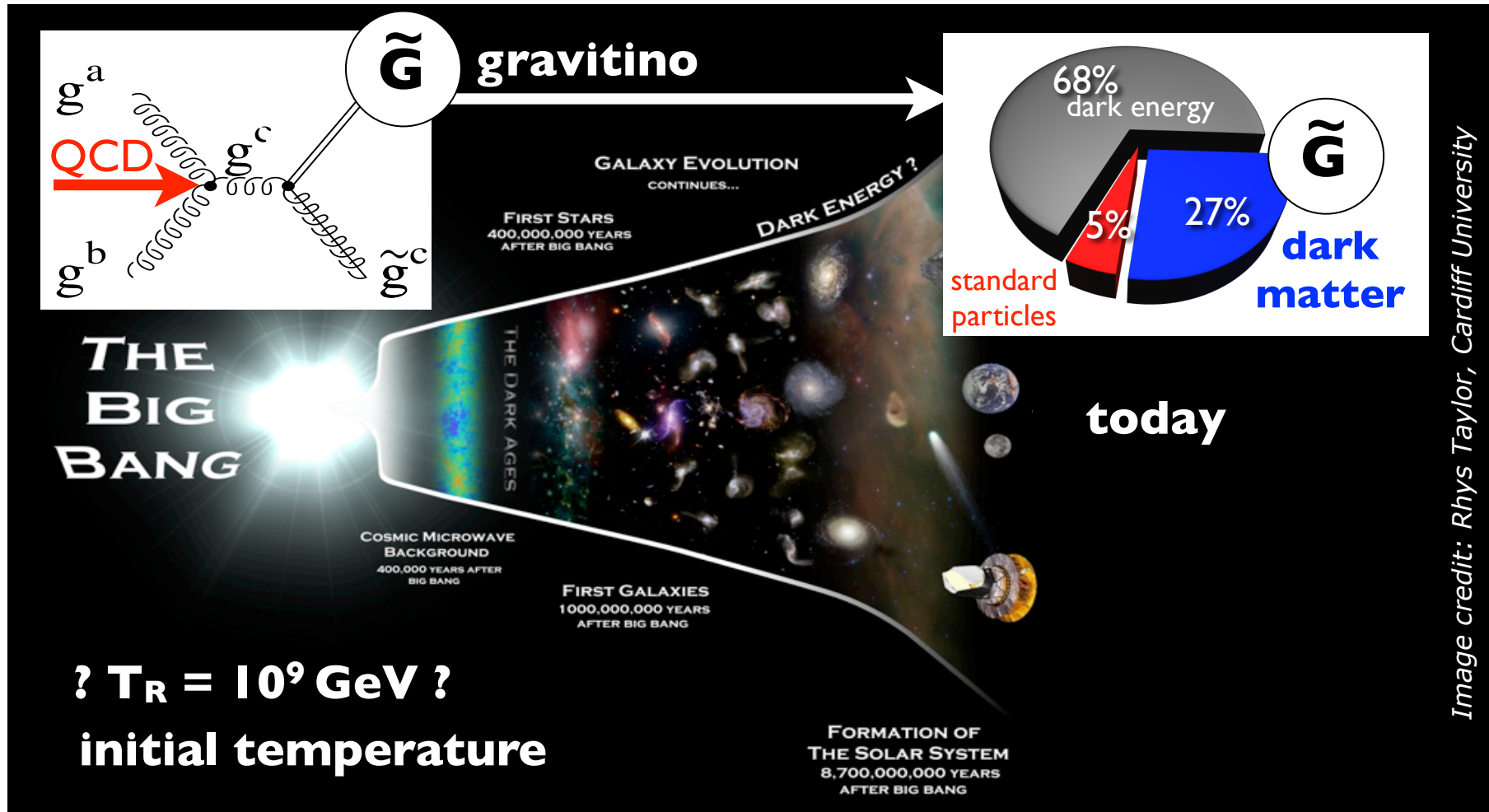
$T_D \gg T: X$ is never in th. eq. with the prim. plasma

but **thermally produced** \rightarrow Boltzmann eq.

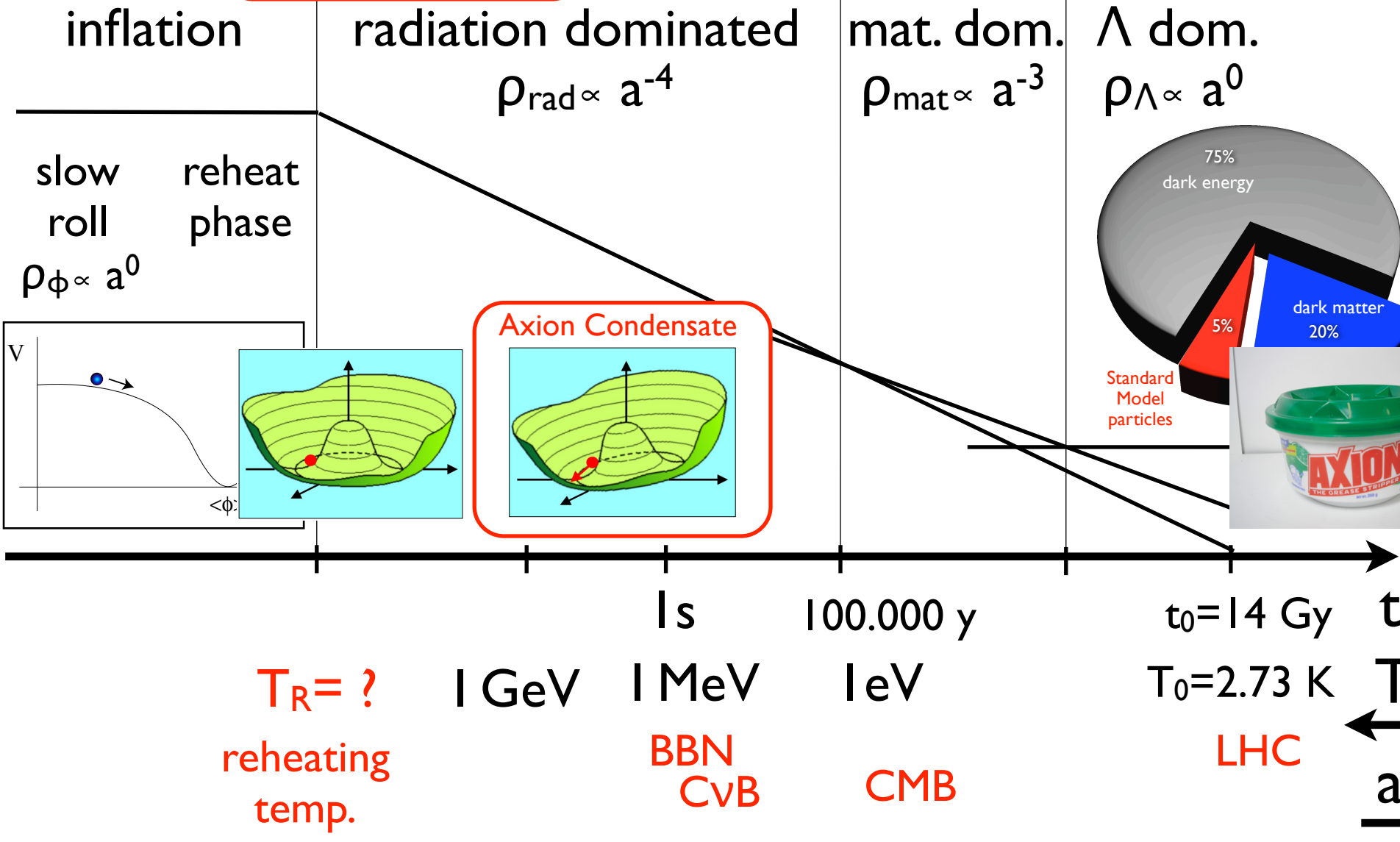
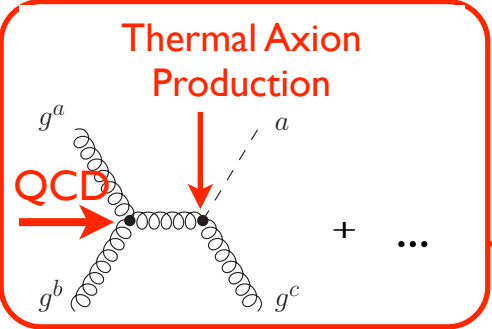


collision term

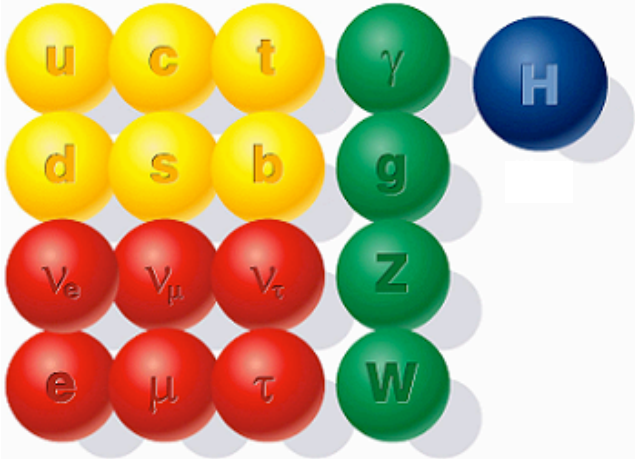
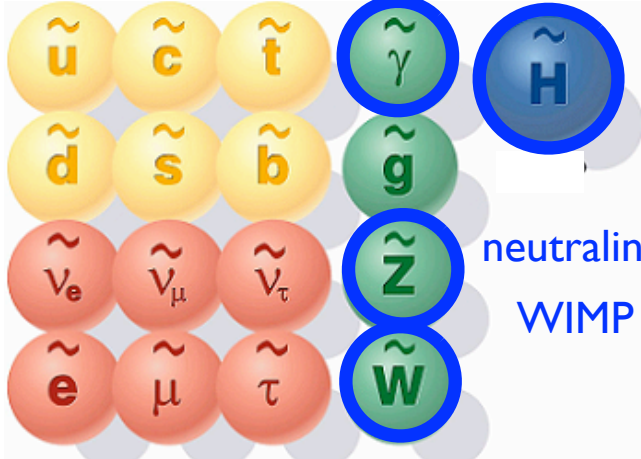
Thermal EWIP DM production in the hot early Universe



Axion Dark Matter

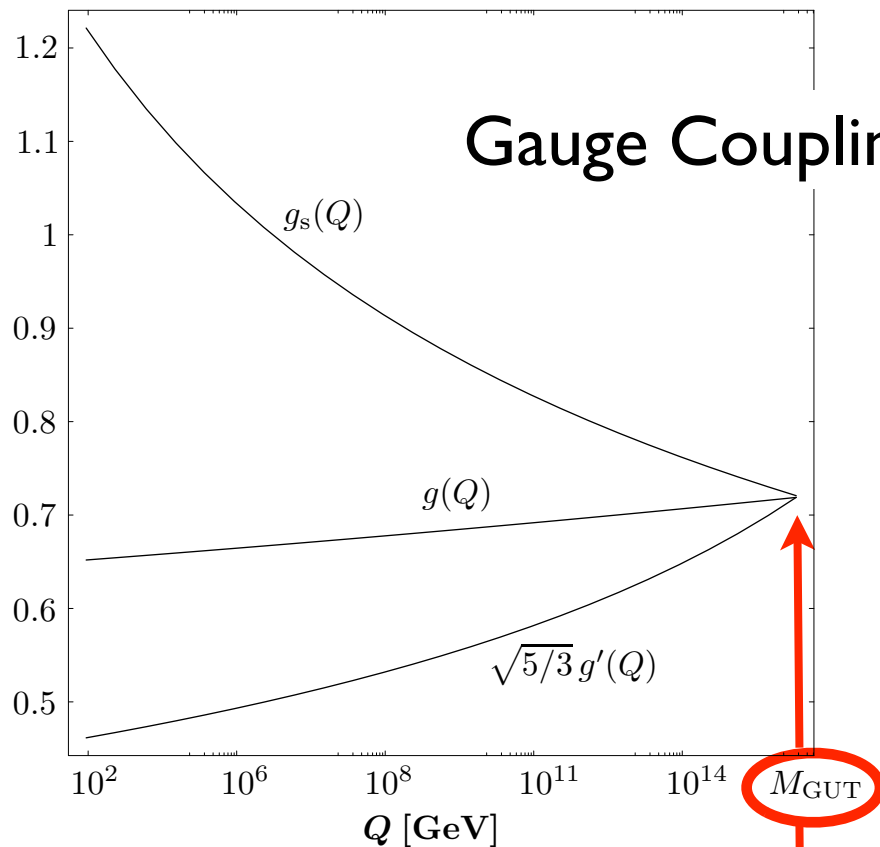


Dark Matter Candidates

interactions	standard particles	superpartners
<p>strong & electroweak</p> <p>extremely weak</p>	<p>Standard Model</p>  <p>Gravity</p> <p>G graviton</p> <p>Peccei-Quinn (PQ) Symmetry</p> <p>a axion EWIP</p>	<p>Supersymmetry</p>  <p>Supergravity</p> <p>G-tilde gravitino EWIP</p> <p>axino EWIP</p> <p>neutralino WIMP</p>

Why Supersymmetry?

Extension of Space-Time Symmetry



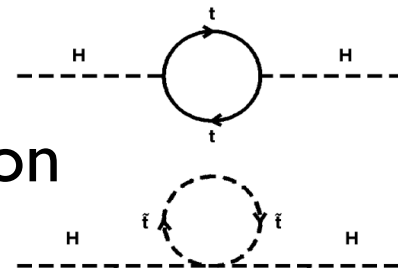
Gauge Coupling Unification

Hierarchy Stabilization

(Super-) Gravity

Consistent String Theory

Dark Matter

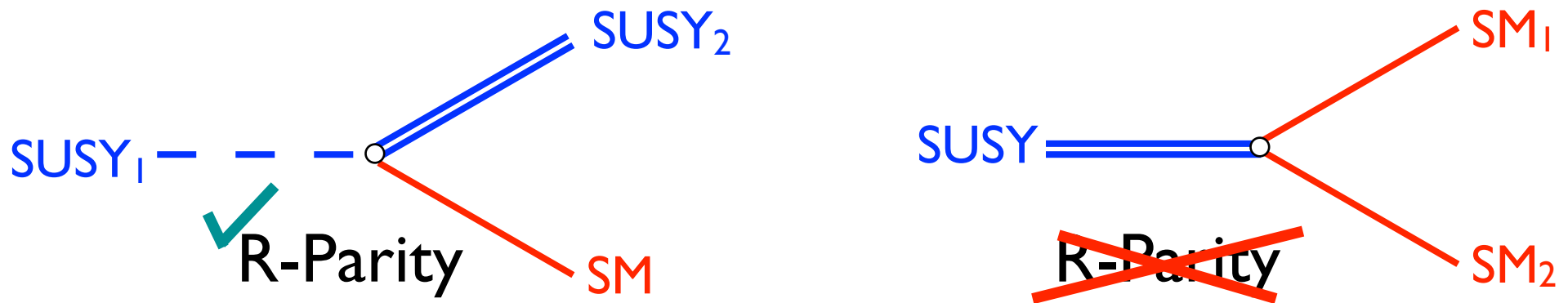


Gauge Coupling Unification at $M_{GUT} \simeq 2 \times 10^{16}$ GeV

R-Parity Conservation

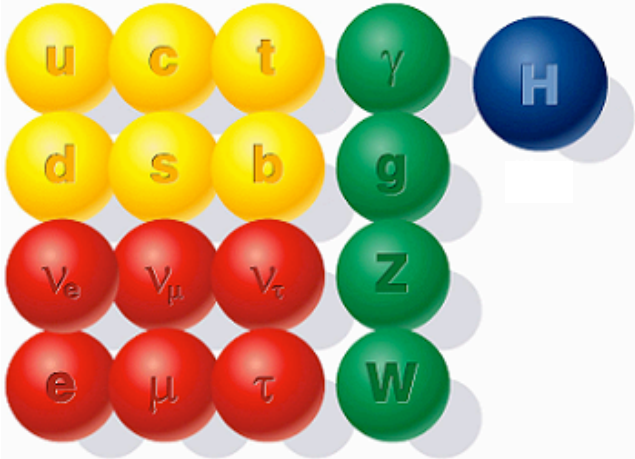
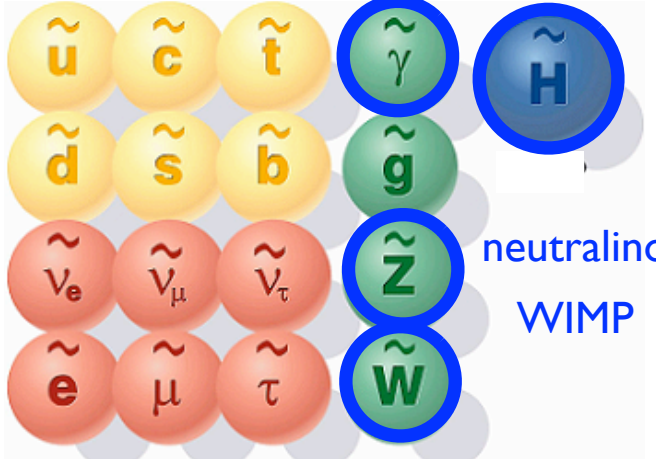
- superpotential: $W_{\text{MSSM}} \leftarrow W_{\Delta L} + W_{\Delta B}$
- non-observation of L & B violating processes (proton stability, ...)
- postulate conservation of R-Parity \leftarrow multiplicative quantum number

$$P_R = (-1)^{3(B-L)+2S} = \begin{cases} +1 & \text{for SM, } H_u, H_d \\ -1 & \text{for } \tilde{X} \leftarrow \text{superpartners} \end{cases}$$

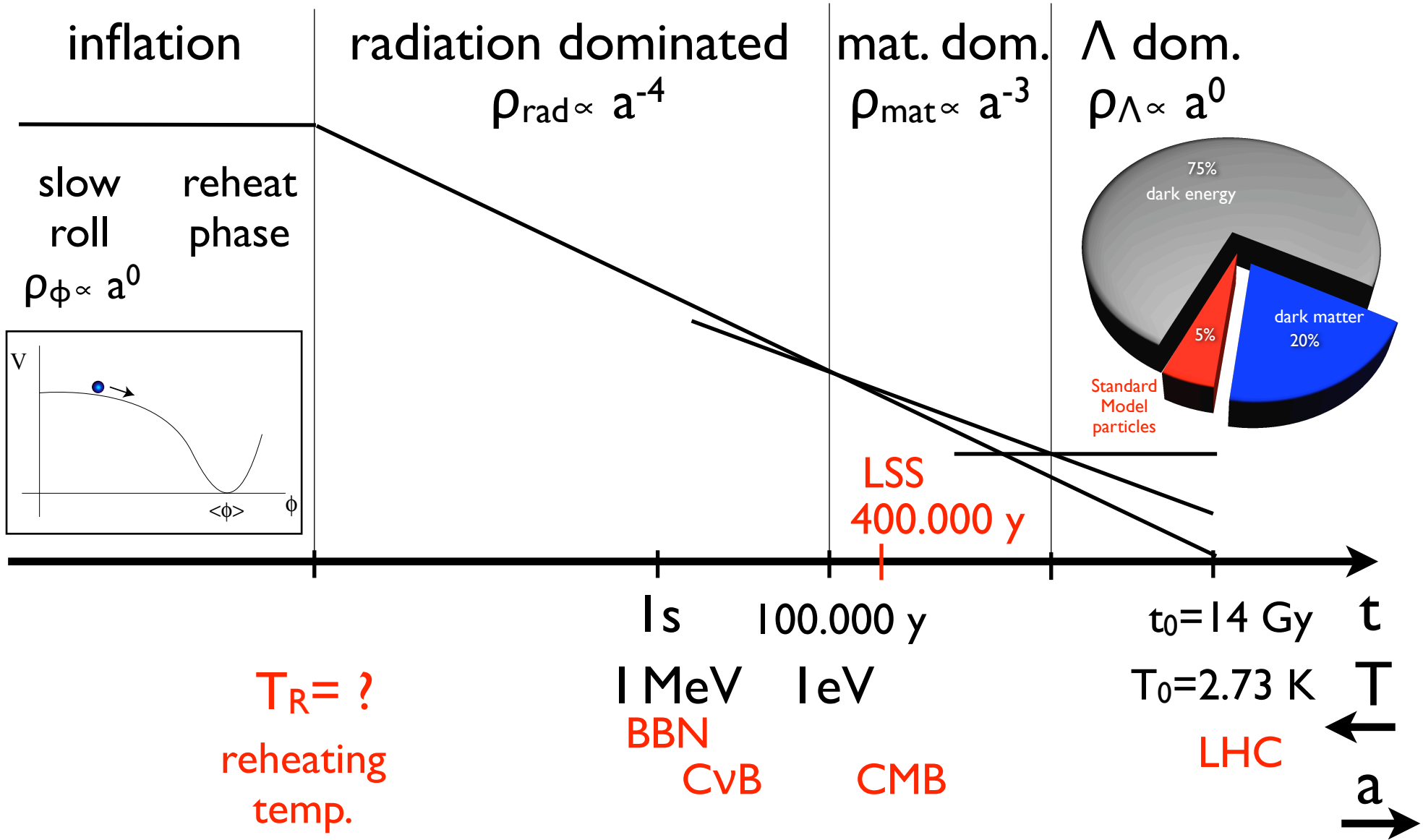


The lightest supersymmetric particle (LSP) is stable!!!

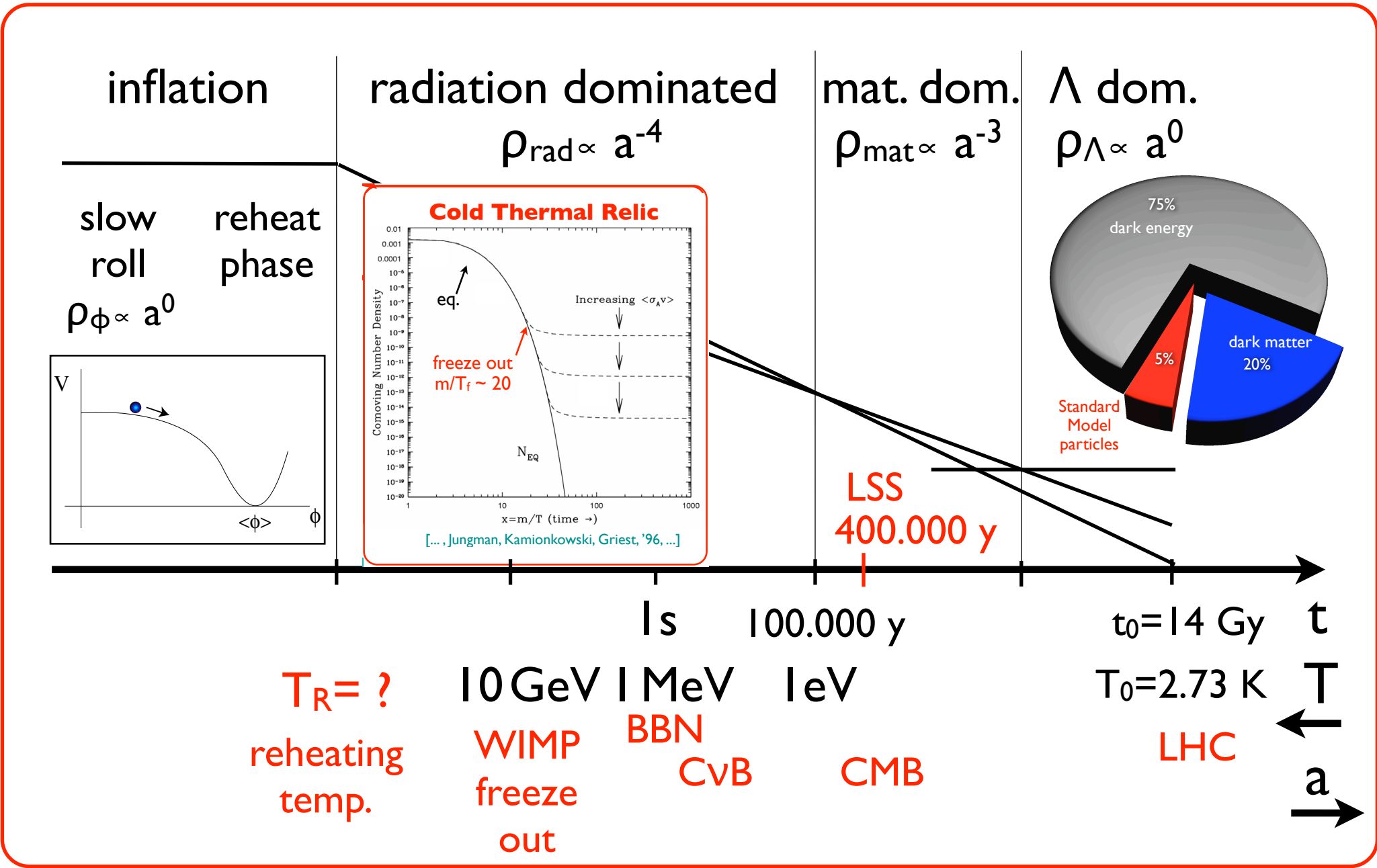
Dark Matter Candidates

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Standard Thermal History of the Universe

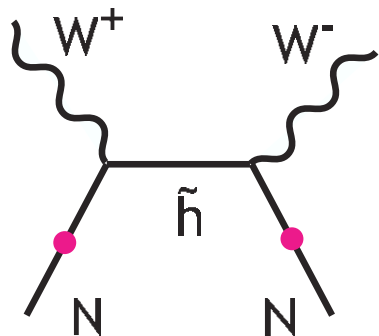


Standard Thermal History of the Universe

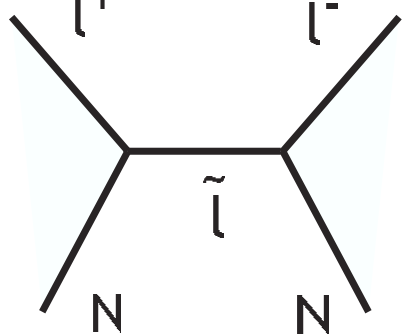


Neutralino LSP Case

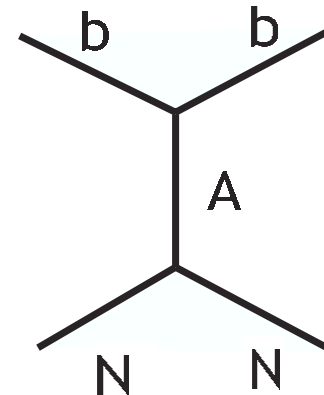
focus point region



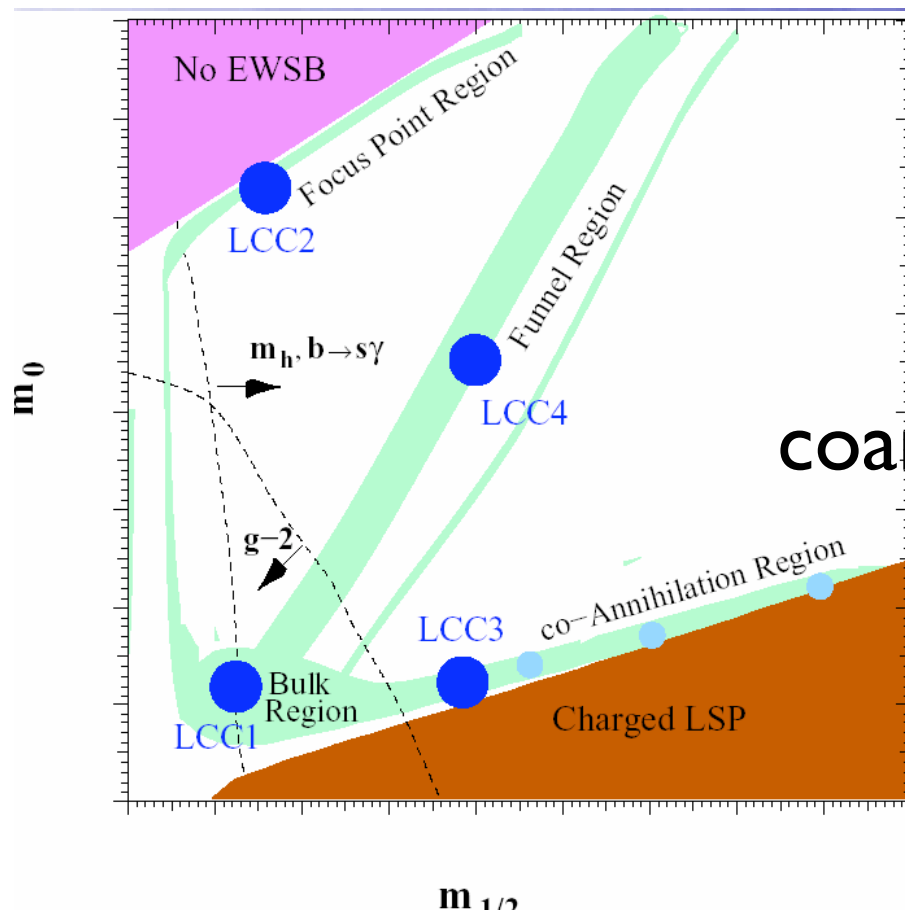
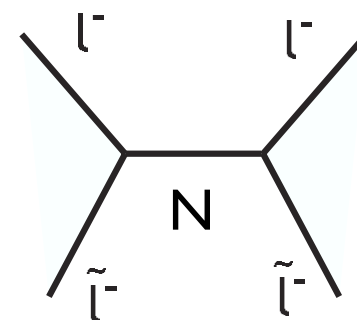
bulk region



annihilation funnel



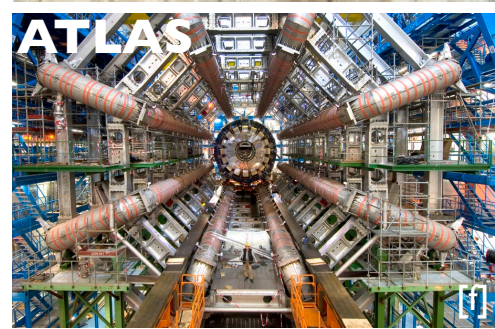
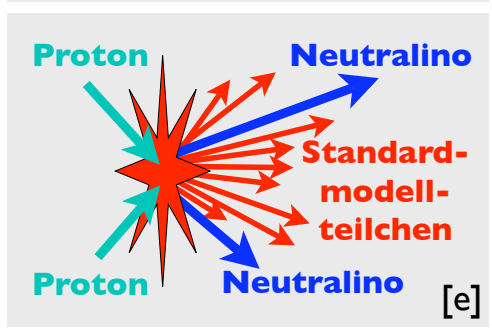
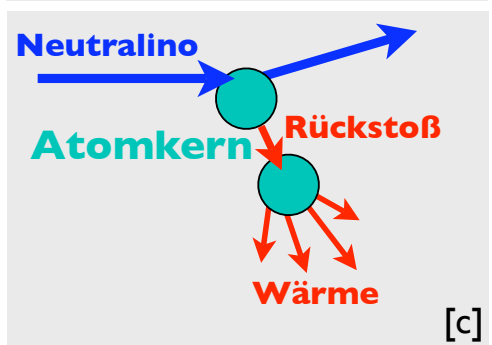
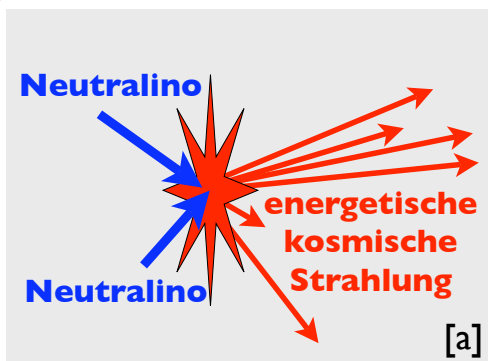
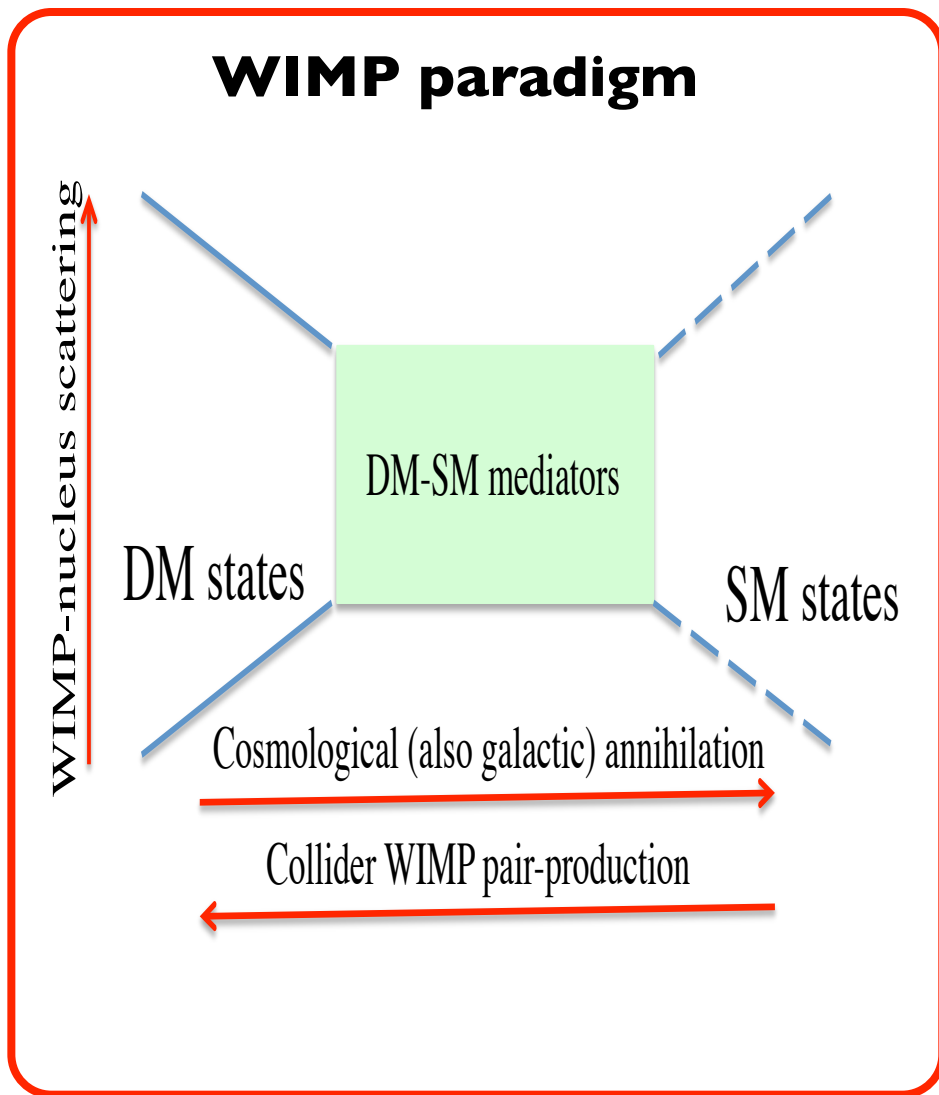
coannihilation region



[see Baltz, Battaglia, Peskin, Wizansky, '06]

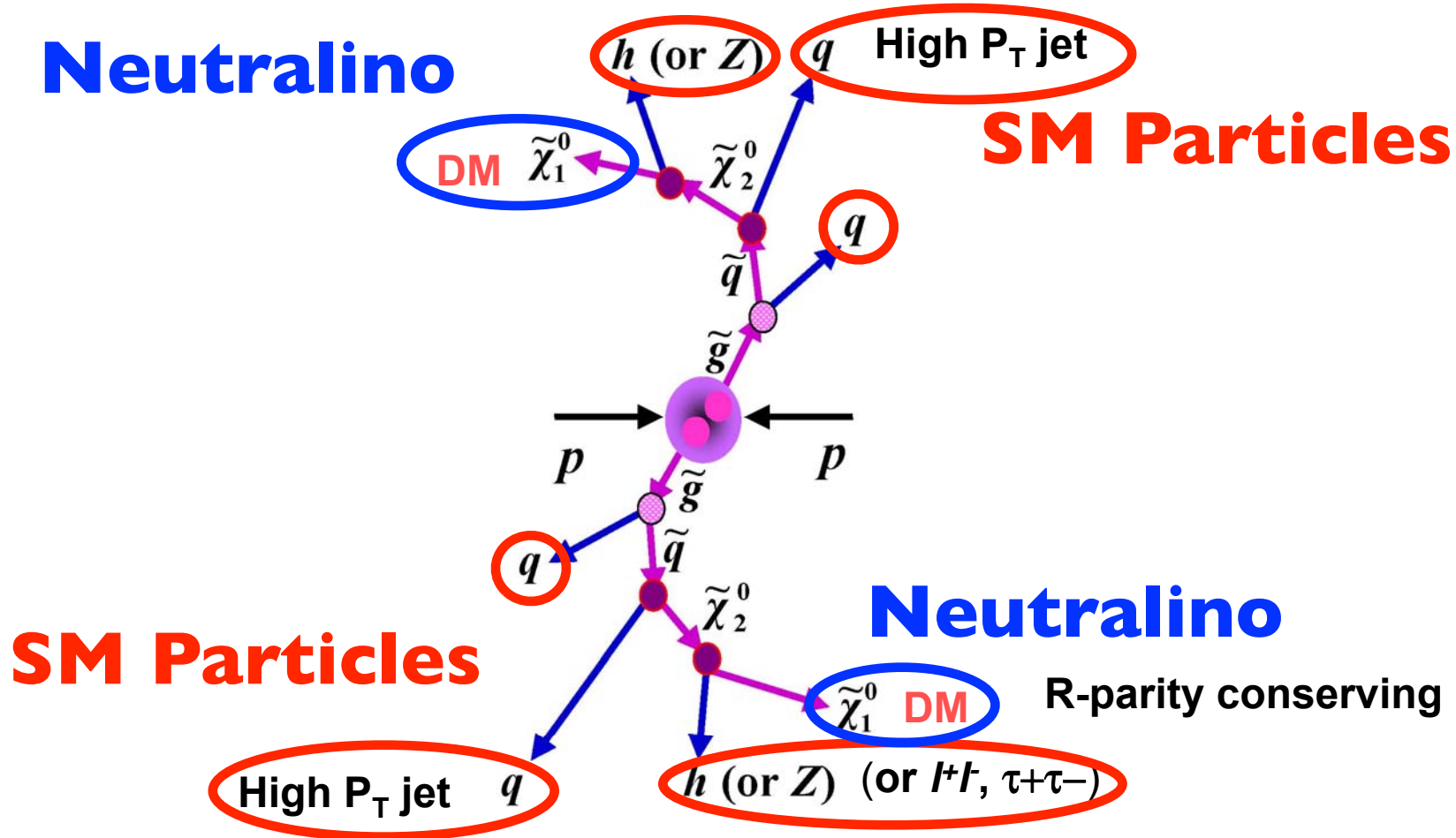
no colored sparticles involved

WIMP paradigm & prospects



Neutralino DM Production at the LHC

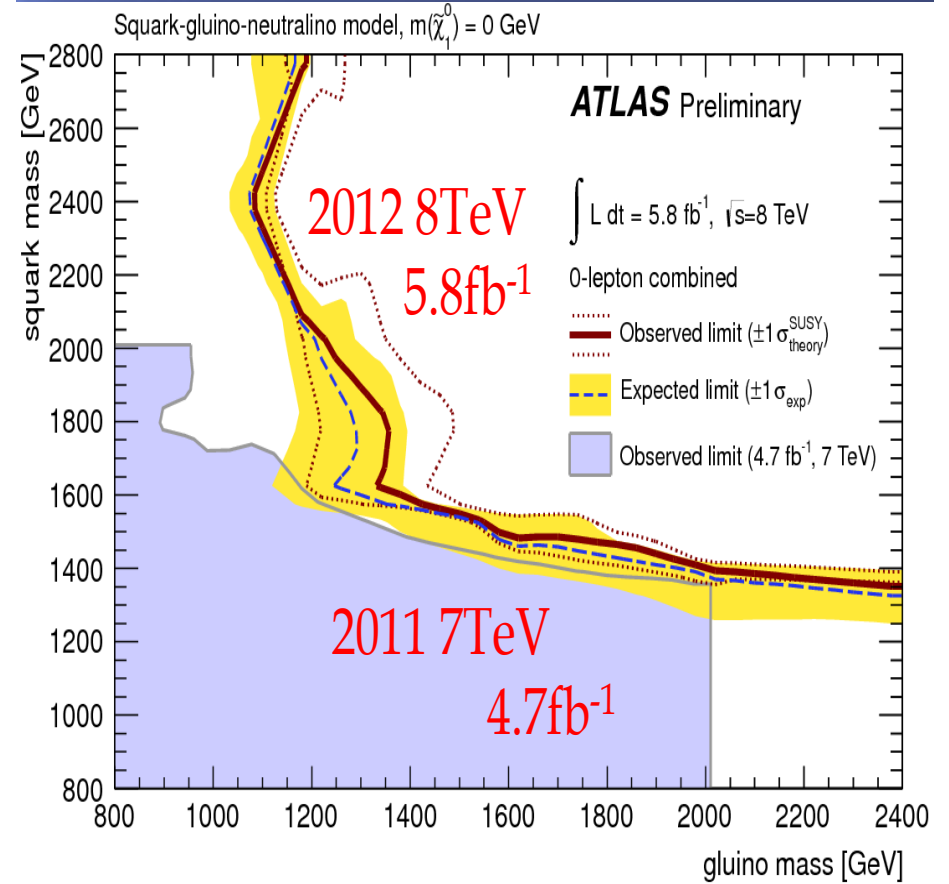
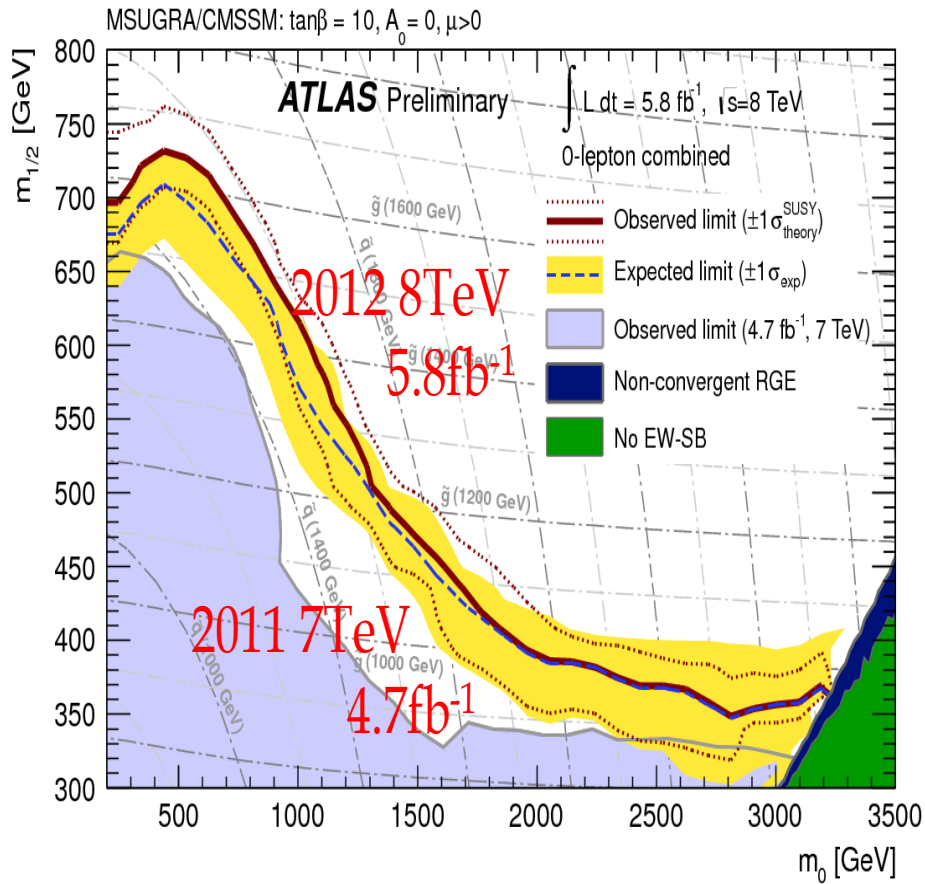
Neutralino



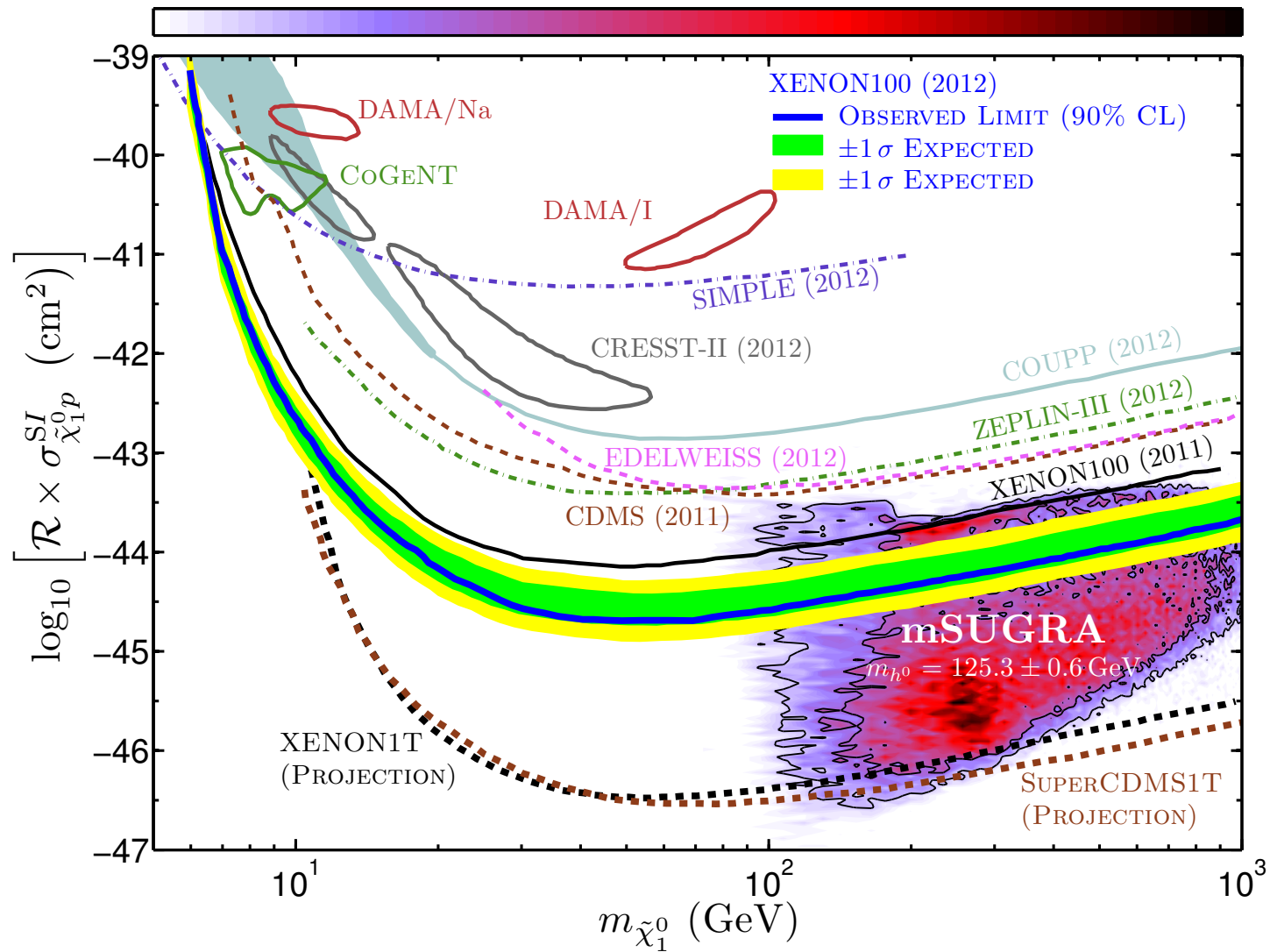
The signal : jets + leptons + missing E_T

[from B. Dutta's Talk, SUSY 2007]

Collider Dark Matter Searches: Limits Only

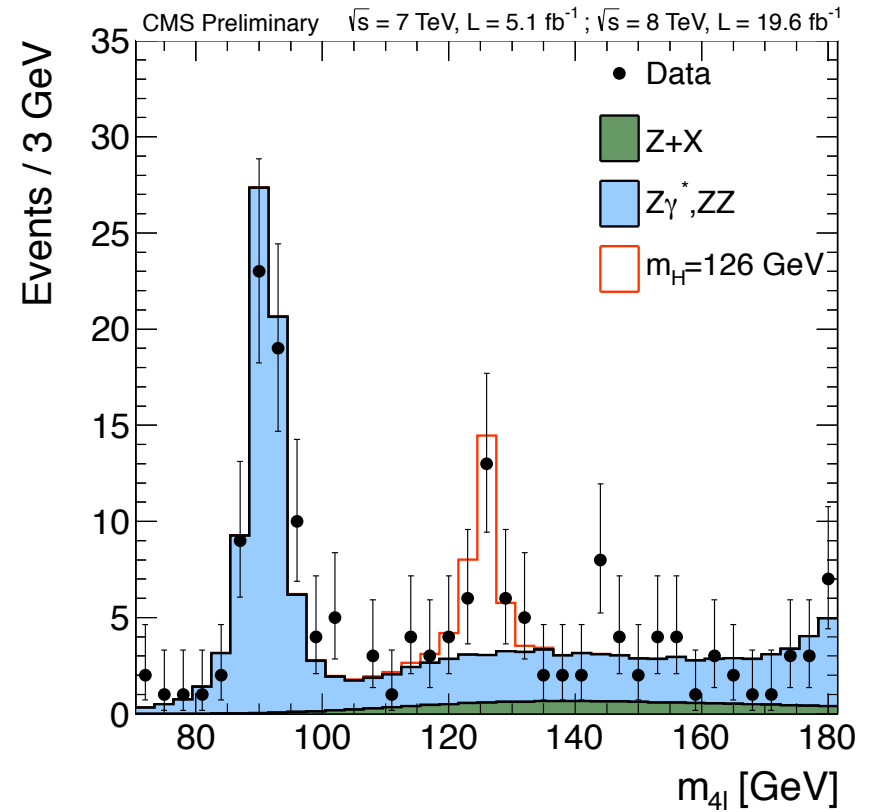
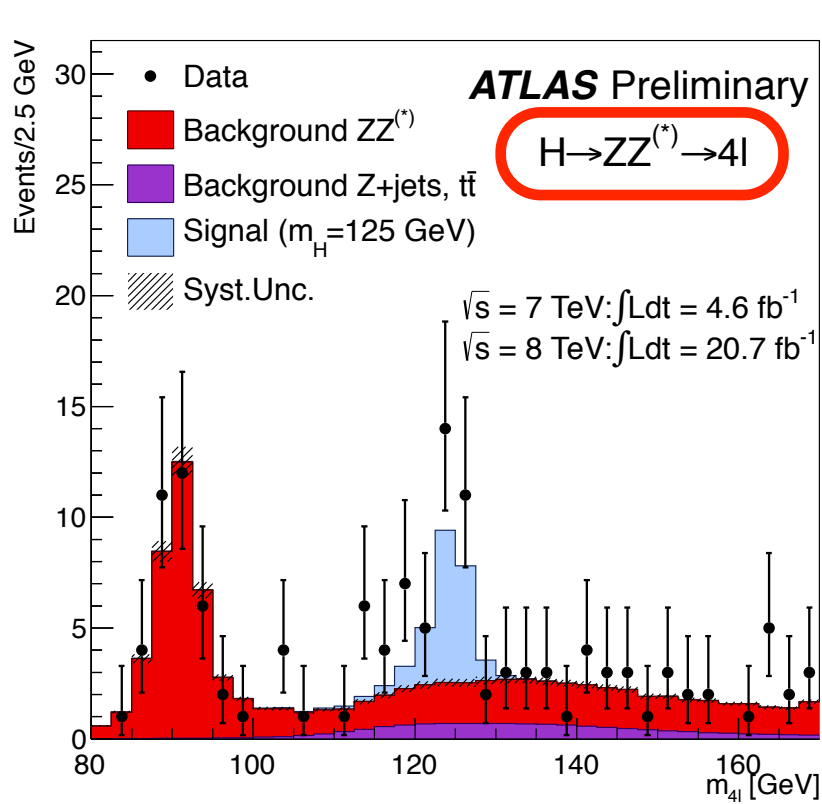


Direct neutralino WIMP dark matter searches



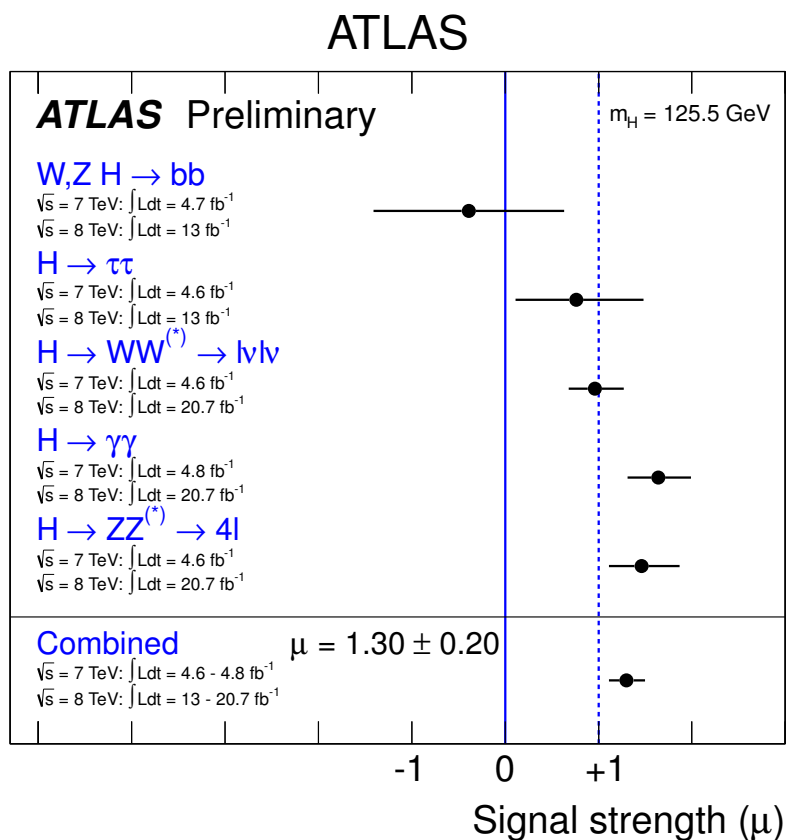
[Akula, Nath, arXiv:1210.0520]

Higgs discovery



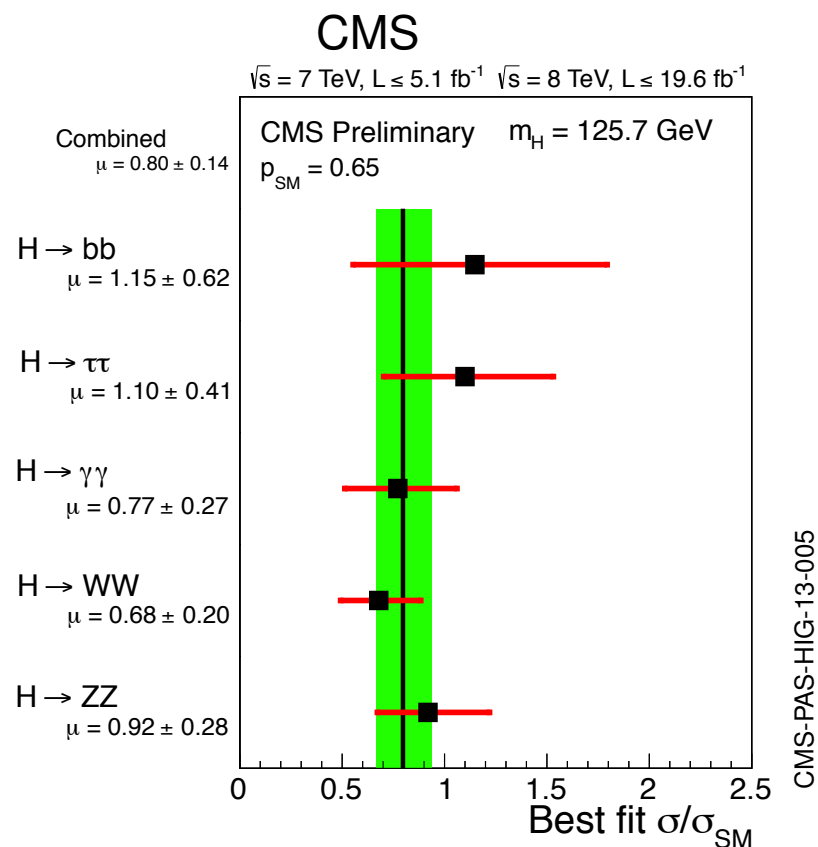
! very impressive !

Higgs discovery



ATLAS-CONF-2013-034

note: bb and $\tau\tau$ have been updated to full 2012 dataset (recently)

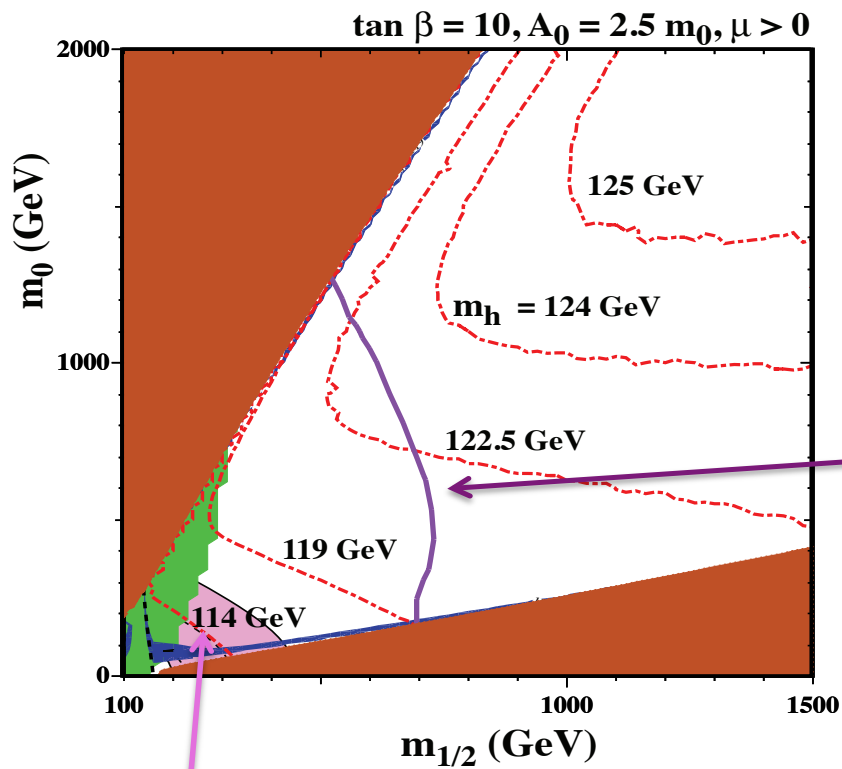
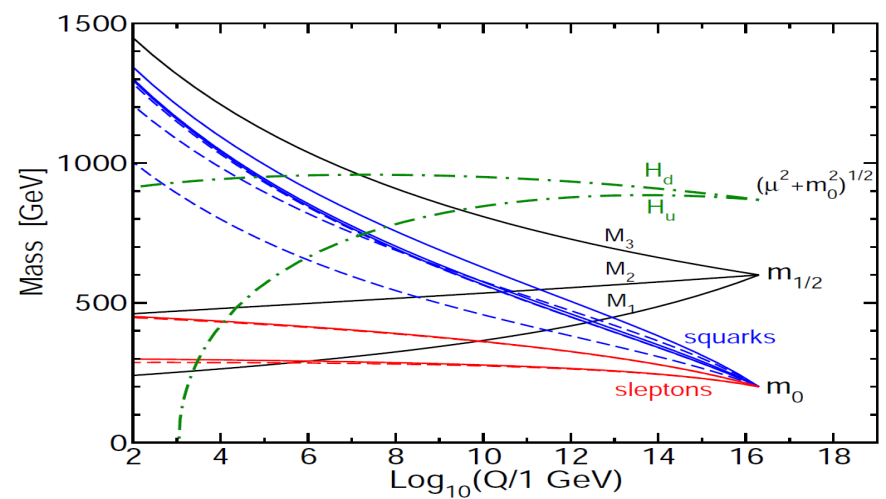


CMS-PAS-HIG-13-005

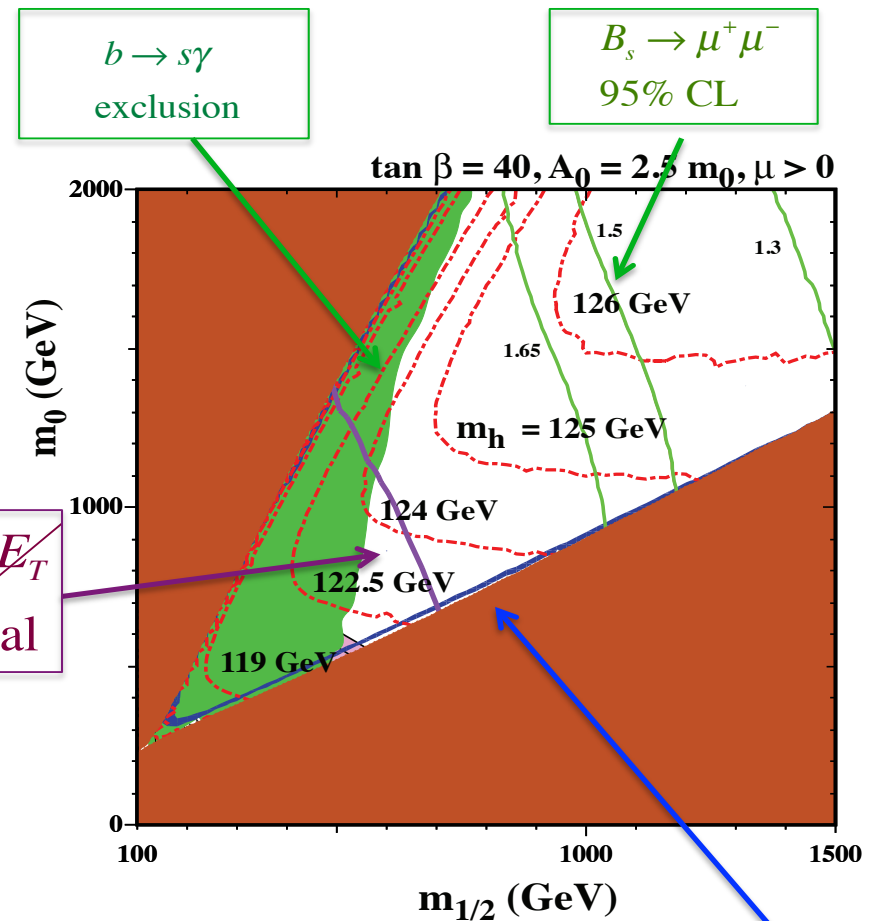
Signal strengths & Standard Model expectations

Constrained MSSM

- m_h contours
- sparticle searches
- $(g-2)_\mu$ anomaly



$g-2$



$b \rightarrow s \gamma$
exclusion

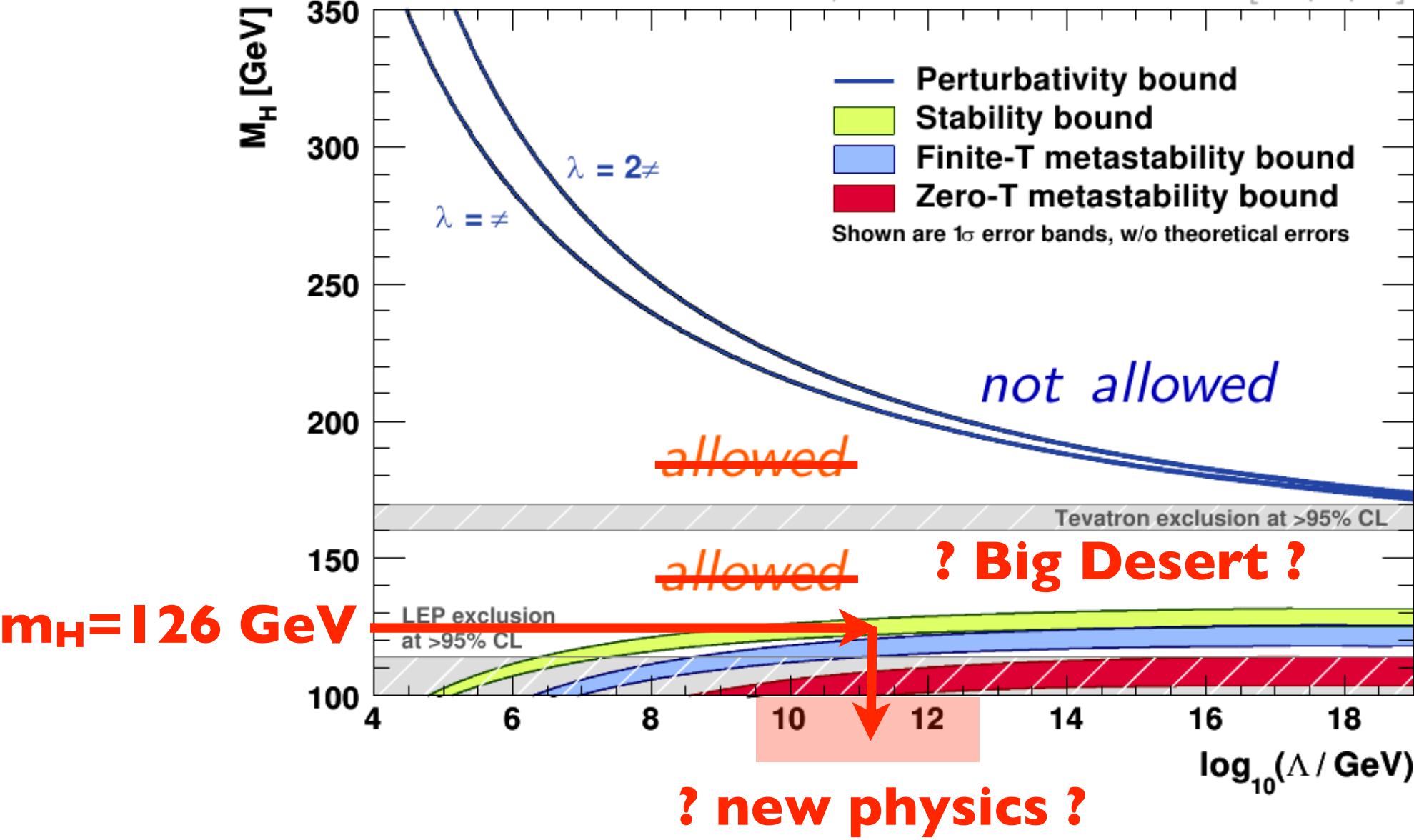
$B_s \rightarrow \mu^+ \mu^-$
95% CL

No E_T
signal

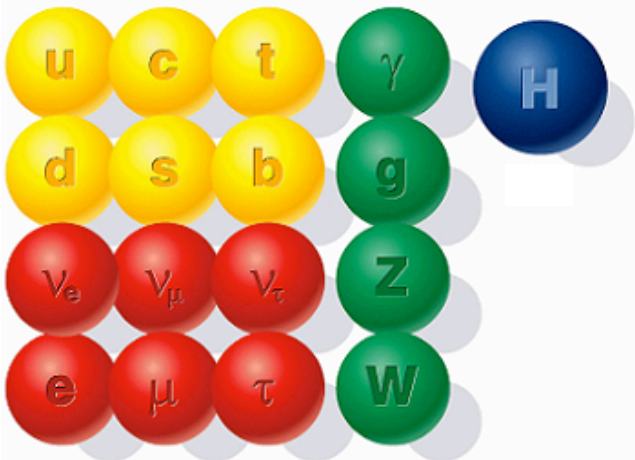
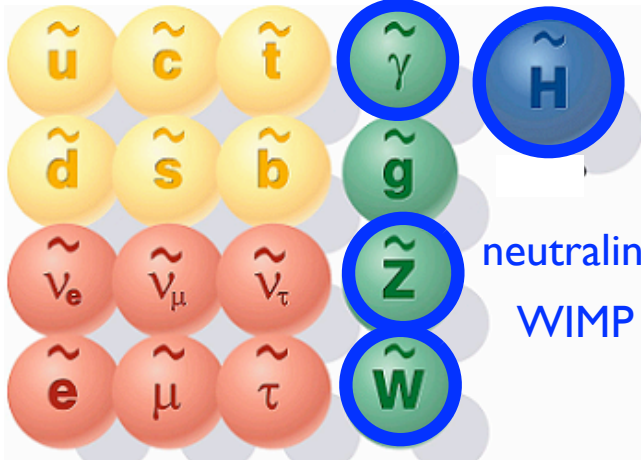
Stau co-annihilation

Standard Model after the Higgs discovery

Ellis et al., arXiv:0906.0954v2 [hep-ph]



Dark Matter Candidates

interactions	standard particles	superpartners
<p>strong & electroweak</p> <p>$\propto (p/M_W)^n$</p> <p>$M_W \sim 10^2 \text{ GeV}$</p> <p>extremely weak</p> <p>$\propto (p/M_{Pl})^n$</p> <p>$M_{Pl} = 2.4 \times 10^{18} \text{ GeV}$</p> <p>$\propto (p/f_{PQ})^n$</p> <p>$f_{PQ} > 10^9 \text{ GeV}$</p>	<p>Standard Model</p>  <p>Gravity</p> <p>G graviton</p> <p>Peccei-Quinn (PQ) Symmetry</p> <p>a axion EWIP</p>	<p>Supersymmetry</p>  <p>neutralino WIMP</p> <p>Supergravity</p> <p>G-tilde gravitino EWIP</p> <p>a-tilde axino EWIP sigma saxion EWIP</p>

New Class → Extremely Weakly Interacting Particles (EWIPs)

Supergravity (N=1, d=4)

$$\begin{aligned}
 \frac{1}{e} \mathcal{L} = & -\frac{M_{\text{P}}^2}{2} R + g_{ij}^* \mathcal{D}_\mu \phi^i \mathcal{D}^\mu \phi^{*j} - \frac{1}{2} g^2 \left[(\text{Re}f)^{-1} \right]^{ab} D_a D_b \\
 & + i g_{ij}^* \bar{\chi}_L^j \gamma^\mu \mathcal{D}_\mu \chi_L^i + \varepsilon^{\mu\nu\rho\sigma} \bar{\psi}_{L\mu} \gamma_\nu \mathcal{D}_\rho \psi_{L\sigma} \\
 & - \frac{1}{4} \text{Re}f_{ab} F_{\mu\nu}^a F^{b,\mu\nu} + \frac{1}{8} \varepsilon^{\mu\nu\rho\sigma} \text{Im}f_{ab} F_{\mu\nu}^a F_{\rho\sigma}^b \\
 & + \frac{i}{2} \text{Re}f_{ab} \bar{\lambda}^a \gamma^\mu \mathcal{D}_\mu \lambda^b - e^{-1} \frac{1}{2} \text{Im}f_{ab} \mathcal{D}_\mu \left[e \bar{\lambda}_R^a \gamma^\mu \lambda_R^b \right] \\
 & + \left[-\sqrt{2} g \partial_i D_a \bar{\lambda}^a \chi_L^i + \frac{1}{4} \sqrt{2} g \left[(\text{Re}f)^{-1} \right]^{ab} \partial_i f_{bc} D_a \bar{\lambda}^c \chi_L^i \right. \\
 & \left. + \frac{i}{16} \sqrt{2} \partial_i f_{ab} \bar{\lambda}^a [\gamma^\mu, \gamma^\nu] \chi_L^i F_{\mu\nu}^b - \frac{1}{2M_{\text{P}}} g D_a \bar{\lambda}_R^a \gamma^\mu \psi_\mu \right. \\
 & \left. - \frac{i}{2M_{\text{P}}} \sqrt{2} g_{ij}^* \mathcal{D}_\mu \phi^{*j} \bar{\psi}_\nu \gamma^\mu \gamma^\nu \chi_L^i + \text{h.c.} \right] \\
 & - \frac{i}{8M_{\text{P}}} \text{Re}f_{ab} \bar{\psi}_\mu \gamma^m, \gamma^n \gamma^\mu \lambda^a F_{mn}^a \quad \text{gravitino} \quad \text{gaugino} \quad \text{gauge boson} \\
 & - e^{K/2M_{\text{P}}^2} \left[\frac{1}{4M_{\text{P}}^2} W^* \bar{\psi}_{R\mu} [\gamma^\mu, \gamma^\nu] \psi_{L\nu} + \frac{1}{2M_{\text{P}}} \sqrt{2} D_i W \bar{\psi}_\mu \gamma^\mu \chi_L^i \right. \\
 & \left. + \frac{1}{2} \mathcal{D}_i D_j W \bar{\chi}_L^c \chi_L^j + \frac{1}{4} g^{ij}^* D_j^* W^* \partial_i f_{ab} \bar{\lambda}_R^a \lambda_L^b + \text{h.c.} \right] \\
 & - e^{K/M_{\text{P}}^2} \left[g^{ij}^* (D_i W) (D_j^* W^*) - 3 \frac{|W|^2}{M_{\text{P}}^2} \right] + \mathcal{O}(M_{\text{P}}^{-2}) .
 \end{aligned}$$

Planck scale

Supersymmetric Hadronic Axion Model

$$\mathcal{L}_{\text{PQ}}^{\text{int}} = -\frac{\sqrt{2}\alpha_s}{8\pi f_{\text{PQ}}} \int d^2\theta A W^b W^b + \text{h.c.}$$

Peccei-Quinn (PQ) scale

saxion axion axino

$$A = (\sigma + ia)/\sqrt{2} + \sqrt{2}\theta\tilde{a} + F_A\theta\theta \quad \text{PQ}$$

$$W^b = \tilde{g}^b + D^b\theta - \sigma^{\mu\nu}\theta G_{\mu\nu}^b + i\theta\theta\sigma^\mu D_\mu\tilde{g}^b \quad \text{field strength}$$

gluino

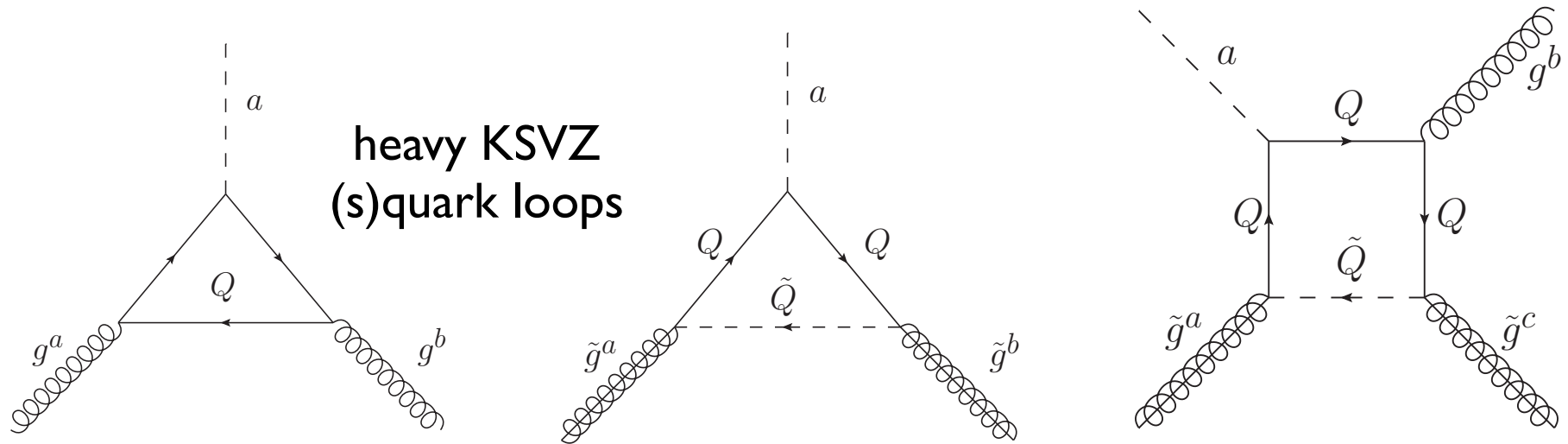
gluon

$$D^b = -g_s \sum_{\tilde{q}} \tilde{q}^* t^b \tilde{q} \quad \text{squark}$$

$$\alpha_s = g_s^2/4\pi$$

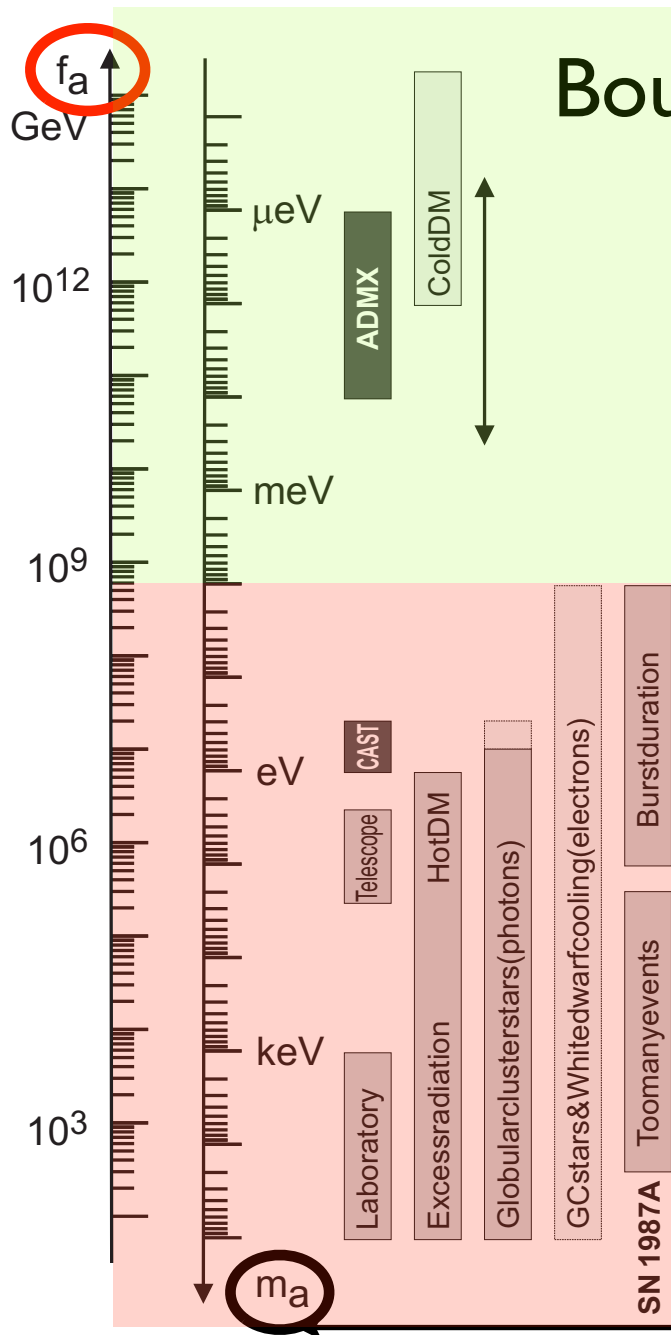
$$\begin{aligned} \mathcal{L}_{\text{PQ}}^{\text{int}} = \frac{\alpha_s}{8\pi f_{\text{PQ}}} & \left[\sigma \left(G^{b\mu\nu} G_{\mu\nu}^b - 2D^b D^b - 2i\tilde{g}_M^b \gamma^\mu D_\mu \tilde{g}_M^b \right) \right. \\ & + a \left(G^{b\mu\nu} \tilde{G}_{\mu\nu}^b + 2\tilde{g}_M^b \gamma^\mu \gamma^5 D_\mu \tilde{g}_M^b \right) \\ & \left. - i\tilde{a}_M \frac{[\gamma^\mu, \gamma^\nu]}{2} \gamma^5 \tilde{g}_M^b G_{\mu\nu}^b + 2\tilde{a}_M D^b \tilde{g}_M^b \right] \end{aligned}$$

Supersymmetric Hadronic (KSVZ) Axion Model



KSVZ [Kim '79; Shifman, Vainshtein, Zakharov '80]

Constraints on the Peccei-Quinn (PQ) scale f_{PQ}



Bounds from Axion Searches

Cosmological Axion Bounds

Astrophysical Axion Bounds

Peccei-Quinn Scale

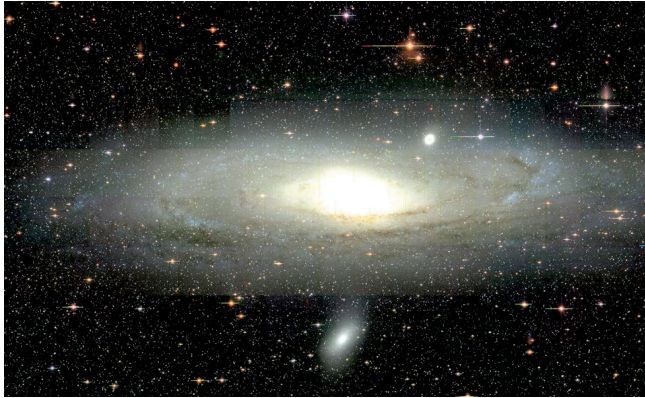
$$f_a \gtrsim 6 \times 10^8 \text{ GeV}$$

Axion Mass

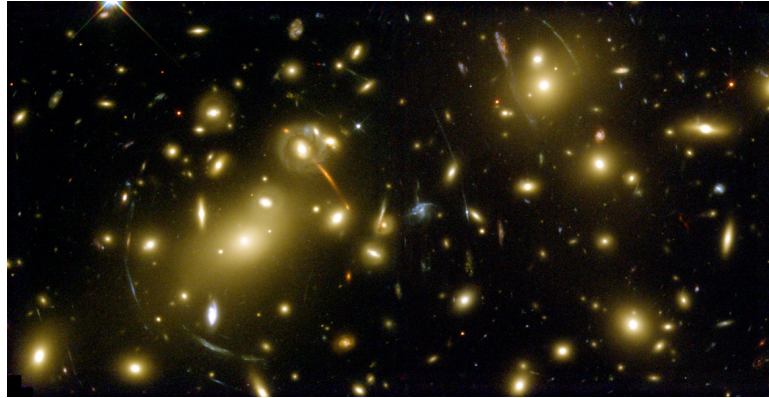
$$m_a \simeq 0.6 \text{ meV} (10^{10} \text{ GeV} / f_{PQ})$$

Extremely Weakly Interacting Particles (EWIPs)

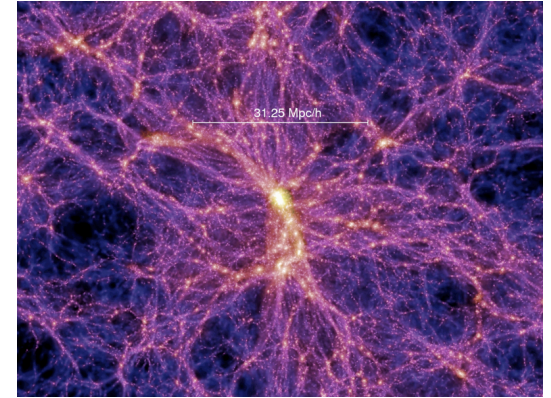
well-motivated candidates for **dark matter**



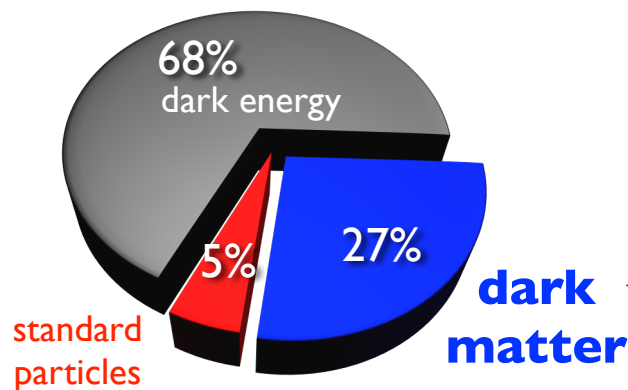
galaxies - rotation velocities



galaxy clusters - gravitational lensing

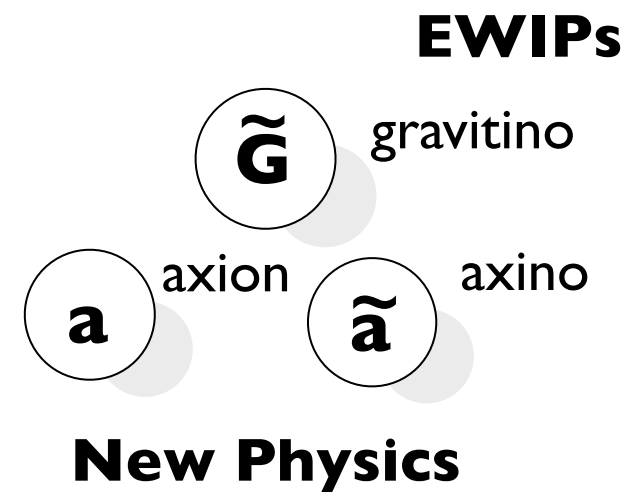


large scale structure



Axions etc.

Origin
???

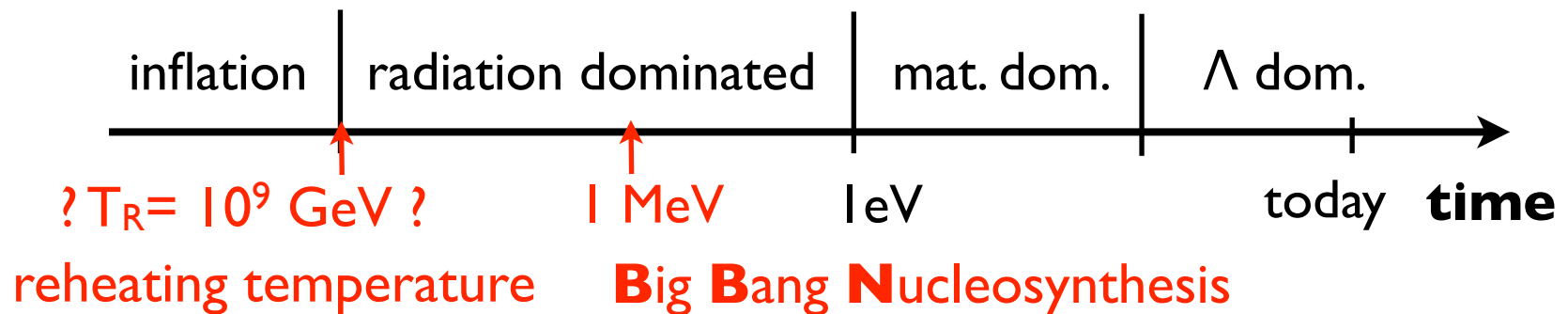
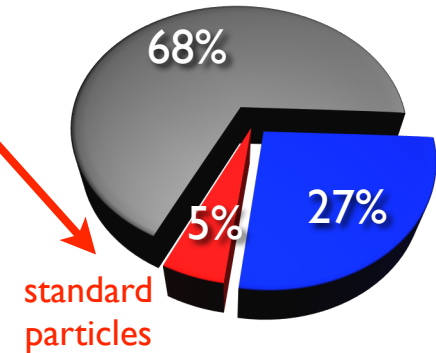


High Reheating Temperature Scenarios

Thermal Leptogenesis

requires $T_R > 10^9 \text{ GeV}$

→ baryon asymmetry



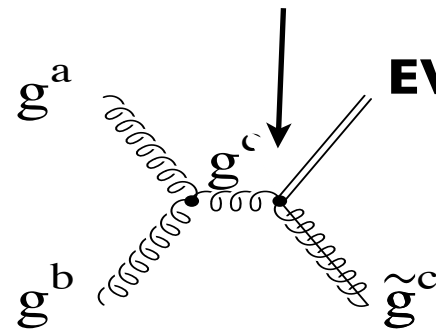
High Reheating Temperature Scenarios

Thermal Leptogenesis

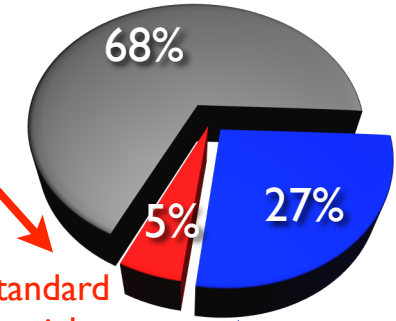
requires $T_R > 10^9 \text{ GeV}$
 \rightarrow baryon asymmetry

Thermal EWIP production

$$\propto (p/M_{\text{Pl}})^n \propto (T_R/M_{\text{Pl}})^n$$

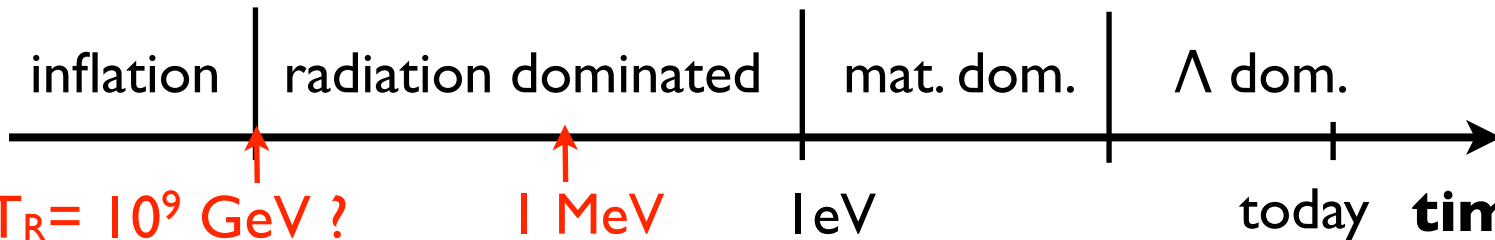


stable & non-relativ.



standard particles

dark matter



reheating temperature

Big Bang Nucleosynthesis

Cosmic Relic Abundances

reheating temp.



decoupling temp. of X



- $T_R > T_D: 1+2 \rightleftharpoons 3+X$

$T > T_D$: X in thermal eq. with the primordial plasma

$T \sim T_D$: X decouples as a **hot thermal relic**

- $T_R > T_D: 1+2 \rightarrow 3+X$

$T_D \gg T$: X is never in th. eq. with the prim. plasma

but **thermally produced** \rightarrow Boltzmann eq.



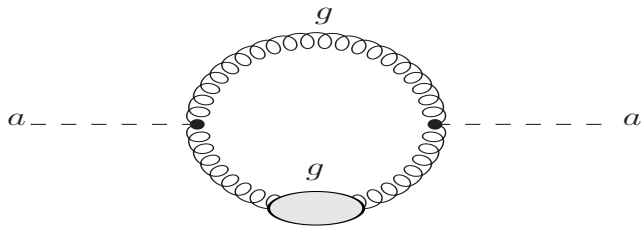
collision term

Calculation of the Collision Term

[cf. Braaten, Yuan, '91]

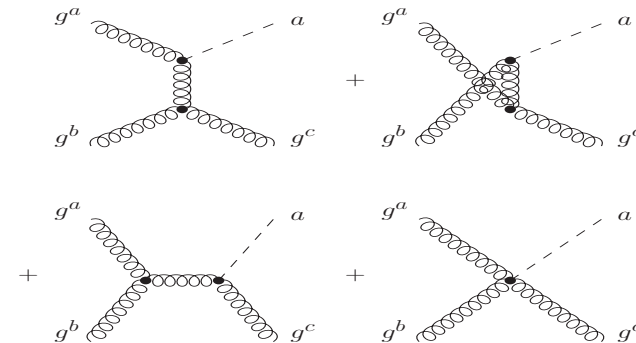
- Thermal Production Rate of X

- Soft Part: $k < k_{\text{cut}}$

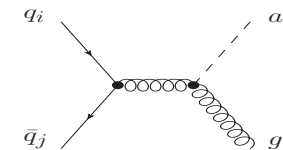


- Hard Part: $k > k_{\text{cut}}$

Process A: $g^a + g^b \rightarrow g^c + a$



Process B: $q_i + \bar{q}_j \rightarrow g^a + a$

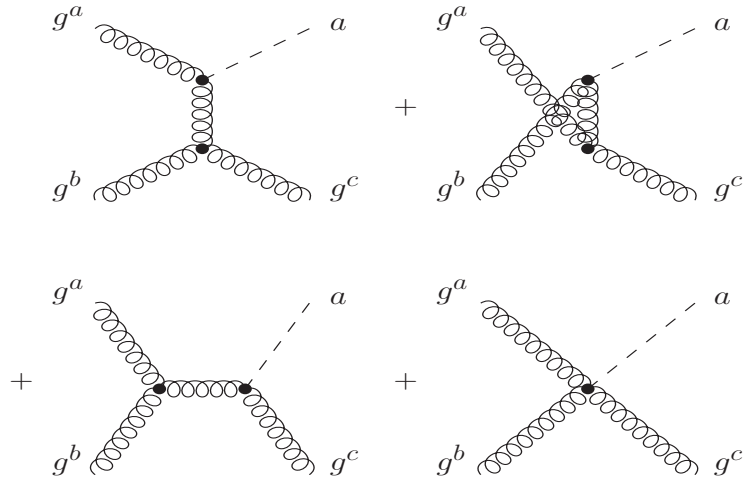


Process C: $q_i + g^a \rightarrow q_j + a$ (crossing of B)

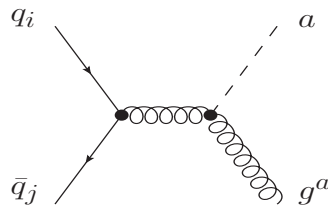
- Hard Thermal Loop (HTL) Resummation

Thermal Axion Production in the Hot QGP

Process A: $g^a + g^b \rightarrow g^c + a$



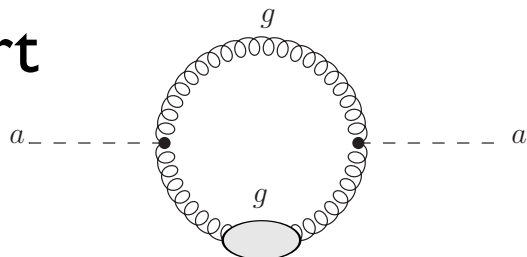
Process B: $q_i + \bar{q}_j \rightarrow g^a + a$



Process C: $q_i + g^a \rightarrow q_j + a$ (crossing of B)

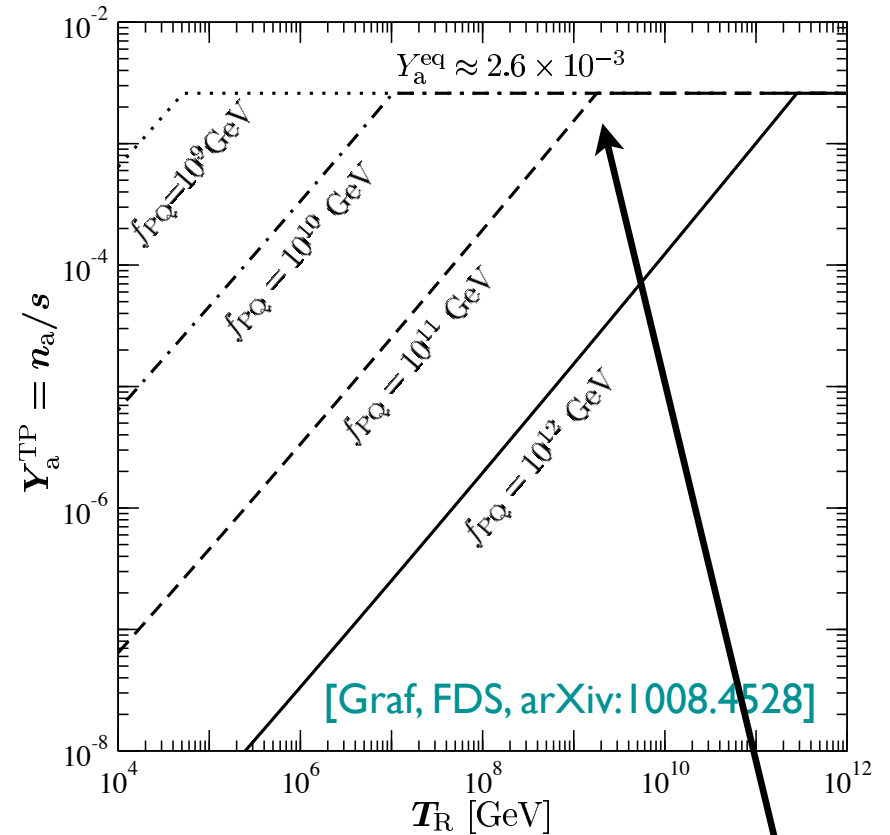
Hard Part

Soft Part



Axions etc.

Axion Yield



Axion Decoupling Temperature

$$T_D \approx 9.6 \times 10^6 \text{ GeV} \left(\frac{f_{\text{PQ}}}{10^{10} \text{ GeV}} \right)^{2.246}$$

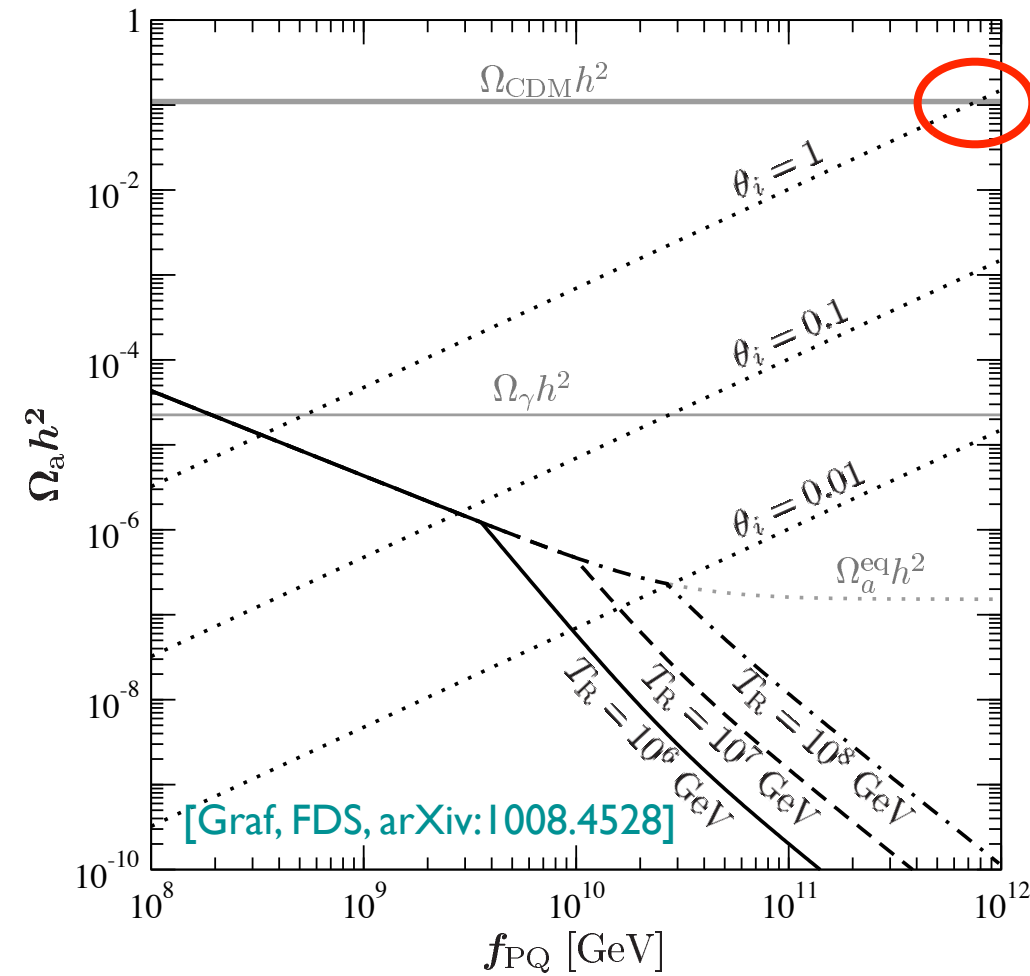
[Masso et al., '02; Sikivie, '08; Graf, FDS, '10]

Axion Dark Matter

Axion Condensate: CDM

$$\Omega_a^{\text{MIS}} h^2 \sim 0.15 \theta_i^2 (f_{\text{PQ}}/10^{12} \text{ GeV})^{7/6}$$

[..., Sikivie, '08; Kim, Carosi, '08, ...]

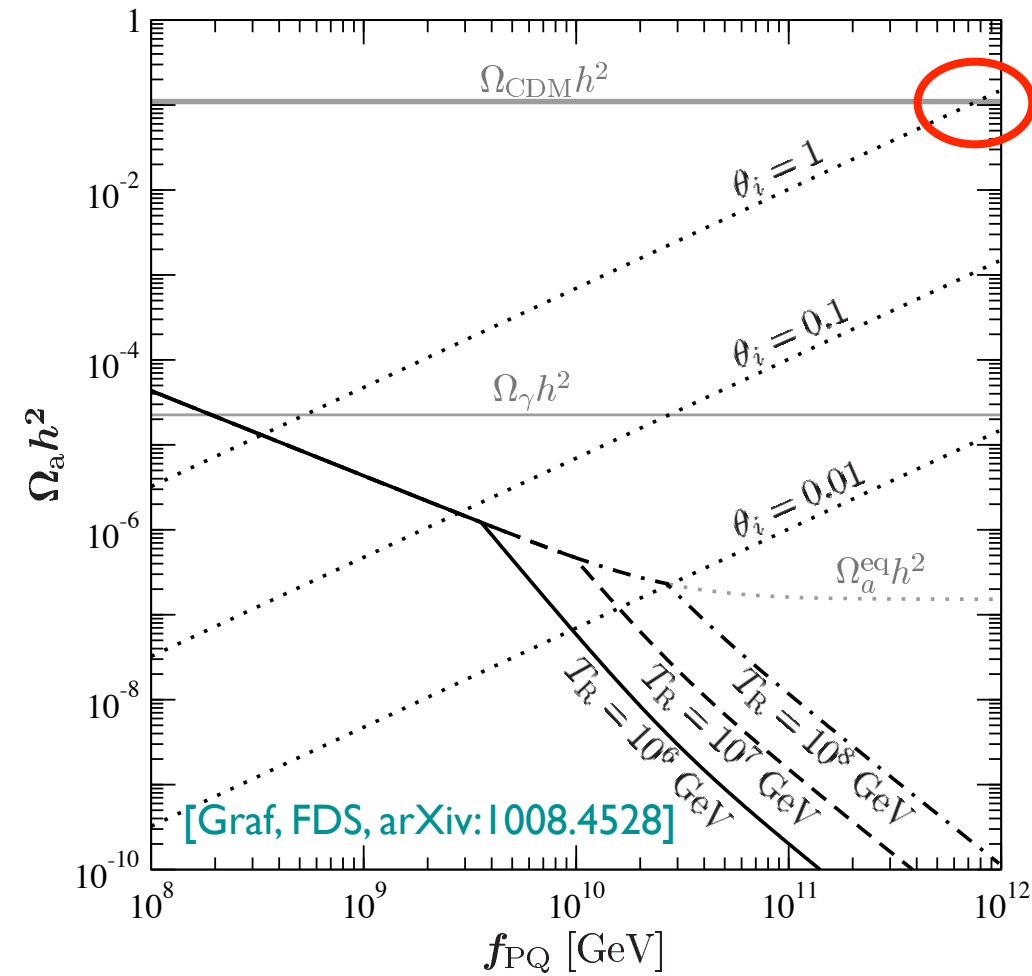


[Graf, FDS, arXiv:1008.4528]

Axion Mass

$$m_a \simeq 0.6 \text{ meV} (10^{10} \text{ GeV} / f_{\text{PQ}})$$

Axion Dark Matter



Axion Condensate: CDM

$$\Omega_a^{\text{MIS}} h^2 \sim 0.15 \theta_i^2 (f_{\text{PQ}}/10^{12} \text{ GeV})^{7/6}$$

[..., Sikivie, '08; Kim, Carosi, '08, ...]

Thermal Axions: WDM/HDM

$$\Omega_a^{\text{TP/eq}} h^2 \simeq \sqrt{\langle p_{a,0} \rangle^2 + m_a^2} Y_a^{\text{TP/eq}} s(T_0) h^2 / \rho_c$$

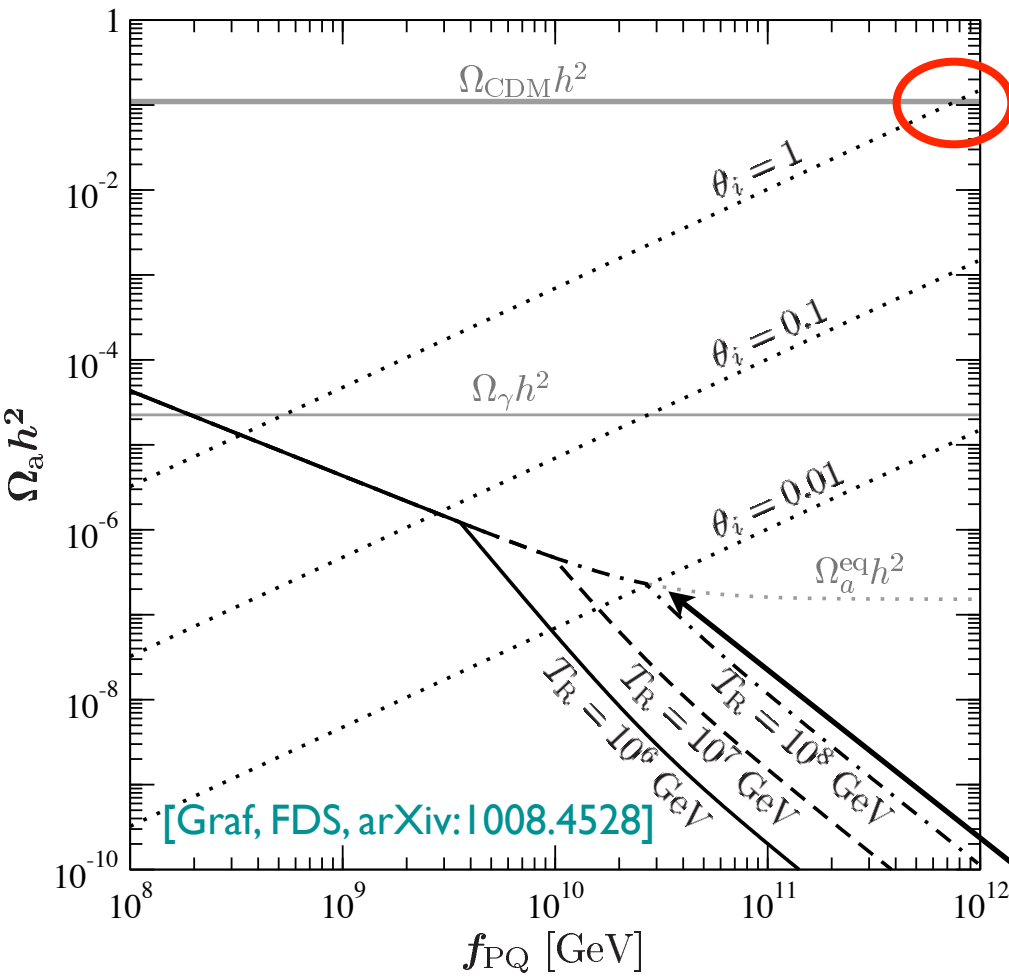
$$18.6 g_s^6 \ln\left(\frac{1.501}{g_s}\right) \left(\frac{10^{10} \text{ GeV}}{f_{\text{PQ}}}\right)^2 \left(\frac{T_{\text{R}}}{10^{10} \text{ GeV}}\right)$$

[Graf, FDS, arXiv:1008.4528]

Axion Mass

$$m_a \simeq 0.6 \text{ meV} (10^{10} \text{ GeV} / f_{\text{PQ}})$$

Axion Dark Matter



[Graf, FDS, arXiv:1008.4528]

Axion Mass

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Axions etc.

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[..., Sikivie, '08; Kim, Carosi, '08, ...]

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$$18.6 g_s^6 \ln\left(\frac{1.501}{g_s}\right) \left(\frac{10^{10} \text{ GeV}}{f_{\text{PQ}}}\right)^2 \left(\frac{T_R}{10^{10} \text{ GeV}}\right)$$

[Graf, FDS, arXiv:1008.4528]

Axion Decoupling Temperature

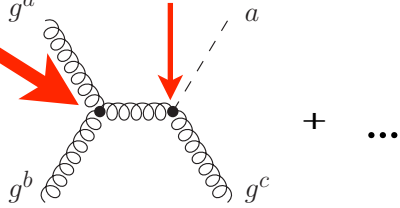
$$T_D \approx 9.6 \times 10^6 \text{ GeV} \left(\frac{f_{\text{PQ}}}{10^{10} \text{ GeV}}\right)^{2.246}$$

[Masso et al., '02; Sikivie, '08; Graf, FDS, '10]

Axion Dark Matter

QCD

Thermal Axion Production



inflation

radiation dominated

mat. dom.

Λ dom.

$$\rho_{\text{rad}} \propto a^{-4}$$

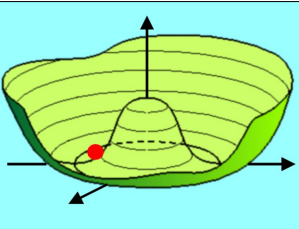
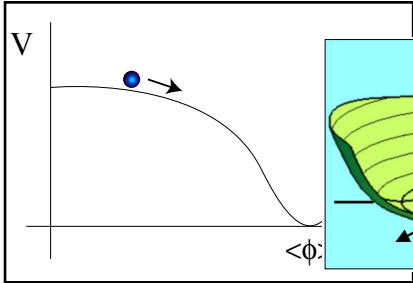
$$\rho_{\text{mat}} \propto a^{-3}$$

$$\rho_{\Lambda} \propto a^0$$

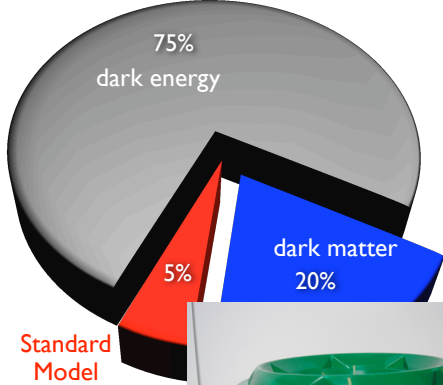
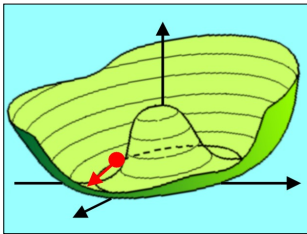
slow roll

reheat phase

$$\rho_{\phi} \propto a^0$$



Axion Condensate



$T_R = ?$
reheating temp.

1 GeV

1 MeV

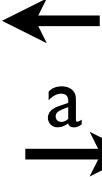
1 eV

$T_0 = 2.73$ K

BBN
CMB

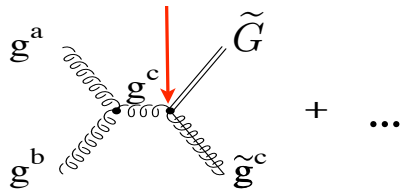
CMB

LHC



Axino Dark Matter

Thermal Axino/
Gravitino Production



inflation

radiation dominated

mat. dom.

Λ dom.

$$\rho_{\text{rad}} \propto a^{-4}$$

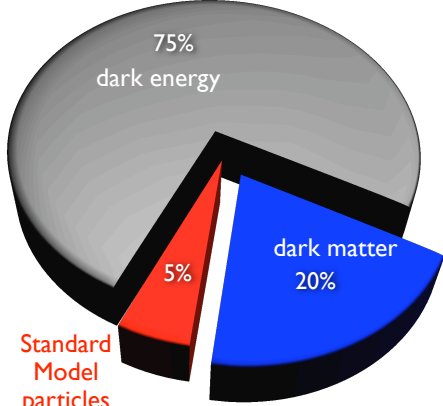
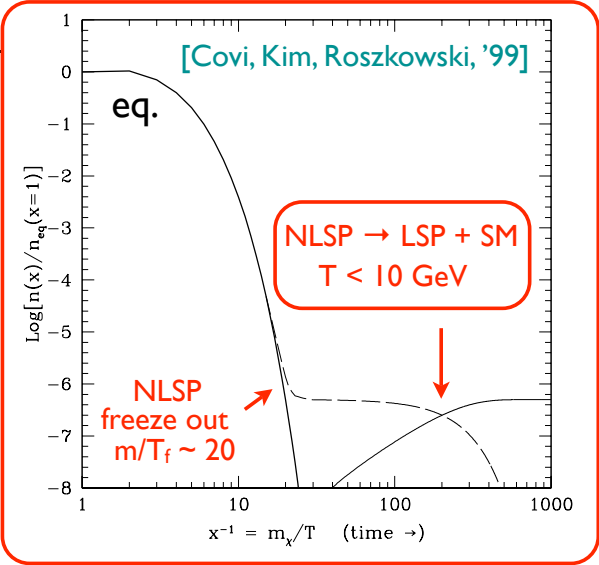
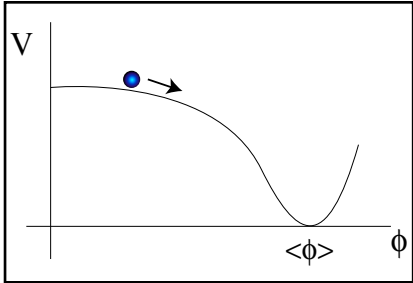
$$\rho_{\text{mat}} \propto a^{-3}$$

$$\rho_{\Lambda} \propto a^0$$

slow roll

reheat phase

$$\rho_{\phi} \propto a^0$$



$T_R = ?$
reheating temp.

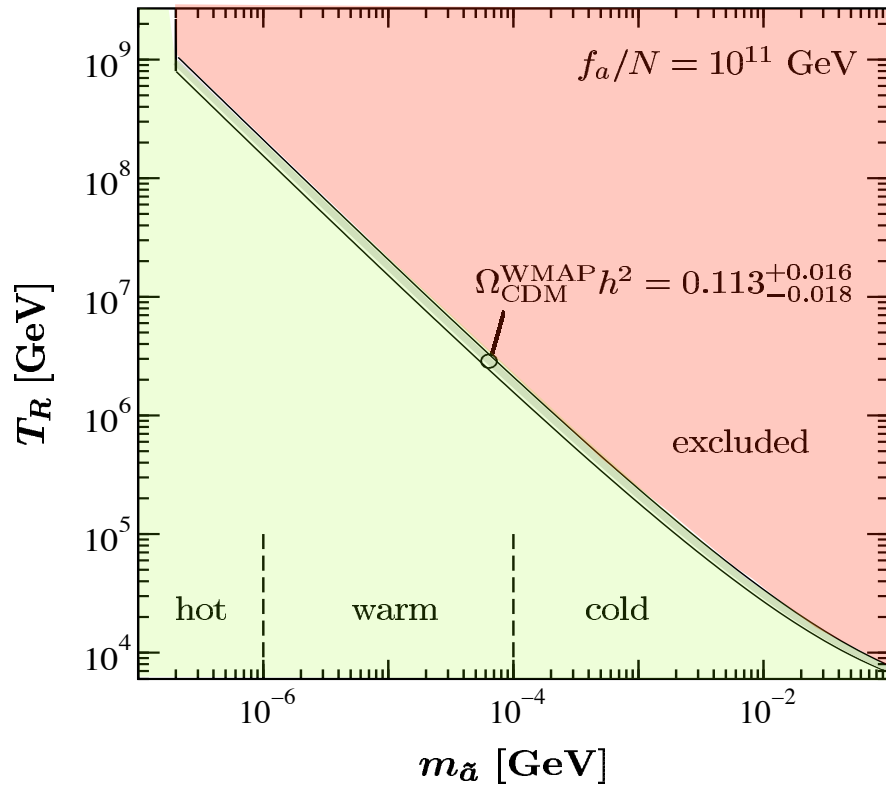
10 GeV WIMP freeze out
1 MeV BBN
leV

$T_0 = 2.73$ K
LHC

t
 T
 a

Axino LSP Case

Thermal \tilde{a} Production



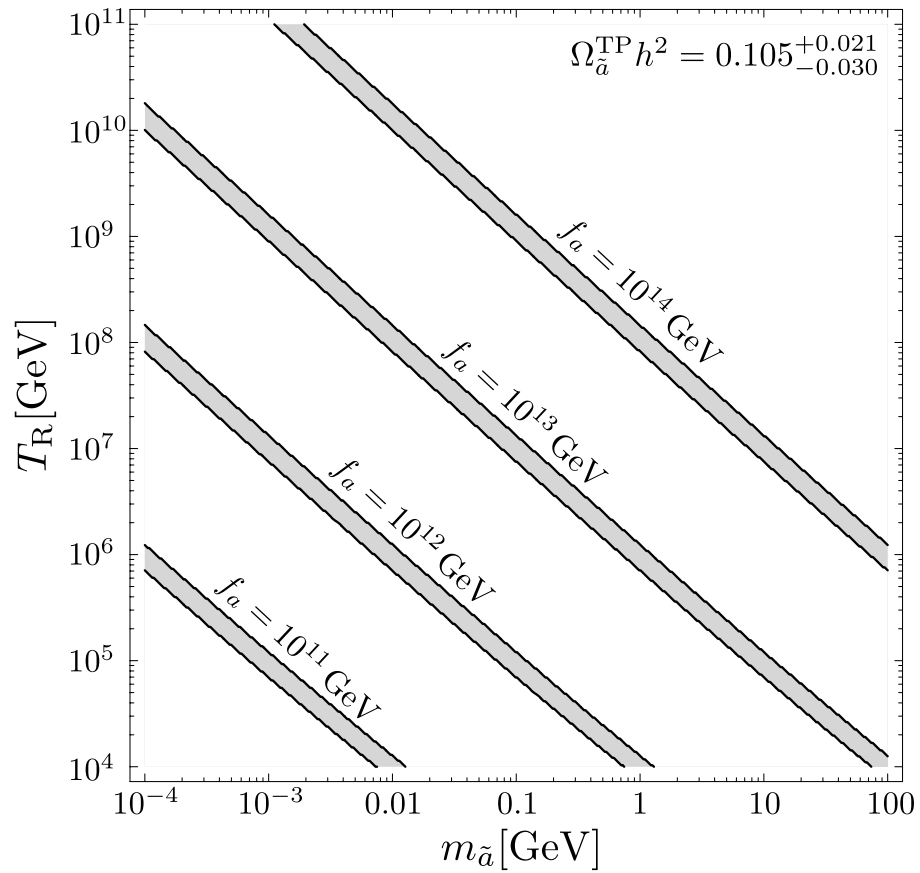
[Brandenburg, FDS, '04]

see also [Covi et al., '01]

and [Strumia, '10]

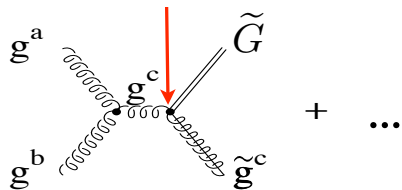
Axino LSP Case

Thermal \tilde{a} Production



[Freitas, FDS, Tajuddin, Wyler, '09]

Thermal Axino/
Gravitino Production



Gravitino Dark Matter

inflation

radiation dominated

mat. dom.

Λ dom.

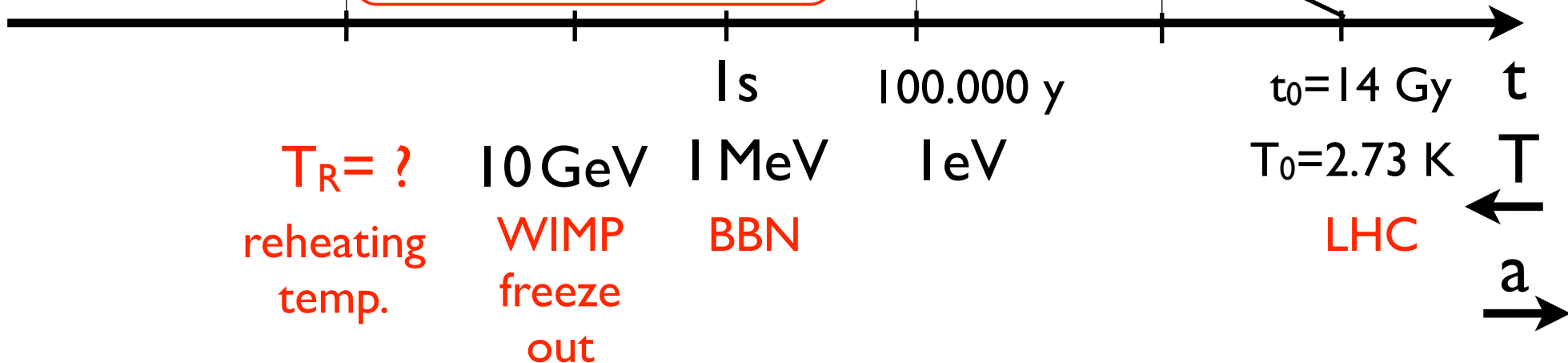
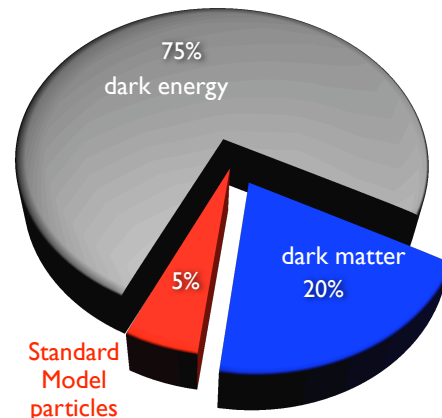
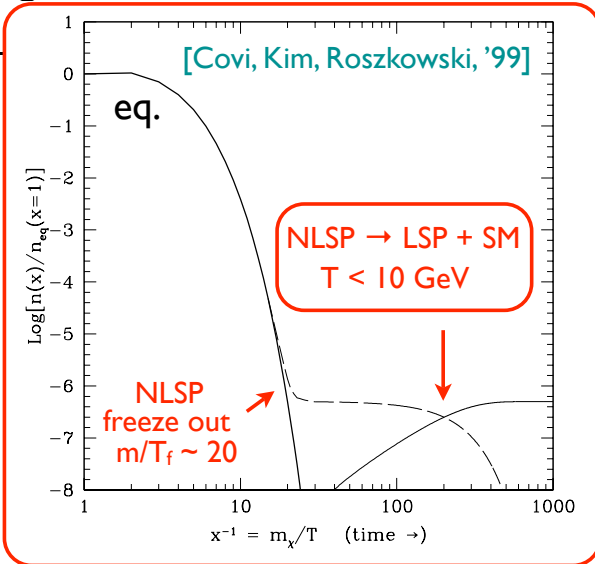
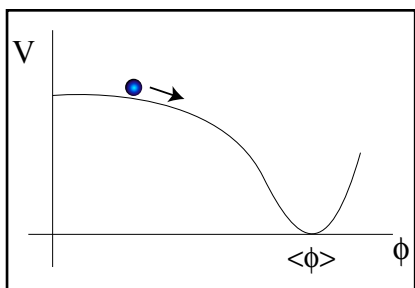
$$\rho_{\text{rad}} \propto a^{-4}$$

$$\rho_{\text{mat}} \propto a^{-3}$$

$$\rho_{\Lambda} \propto a^0$$

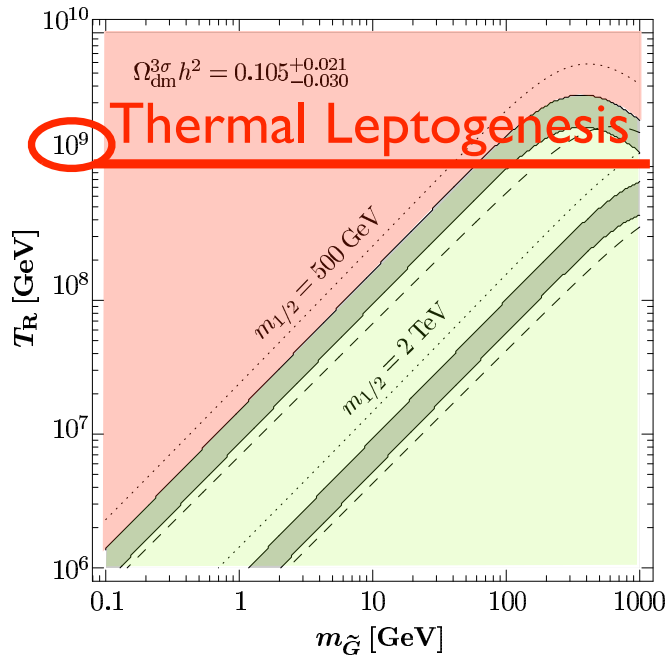
slow roll
 $\rho_{\phi} \propto a^0$

reheat phase

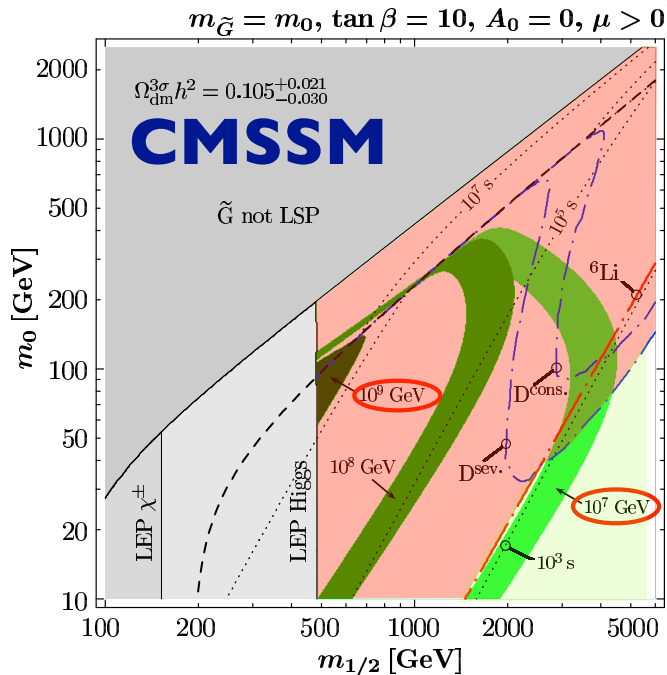


Upper Limits on the Reheating Temperature T_R

Stable Gravitino LSP



Ω_{DM} constraint
for gravitino DM
[Pradler, FDS, '07]

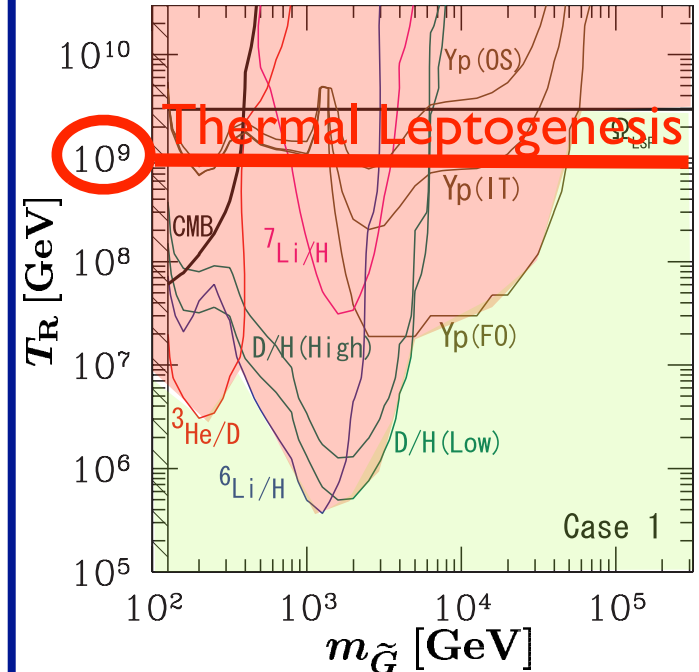


+ BBN constraints

$$T_R \lesssim 4.9 \times 10^7 \text{ GeV} \left(\frac{m_{\tilde{G}}}{10 \text{ GeV}} \right)^{1/5}$$

[Pradler, FDS, arXiv:0710.2213]

Unstable Gravitino



BBN constraints

+ Ω_{DM} constraint
for neutralino DM

[Kohri, Moroi, Yotsuyanagi, '05]

Thermal Leptogenesis requires $T > 10^9$ GeV

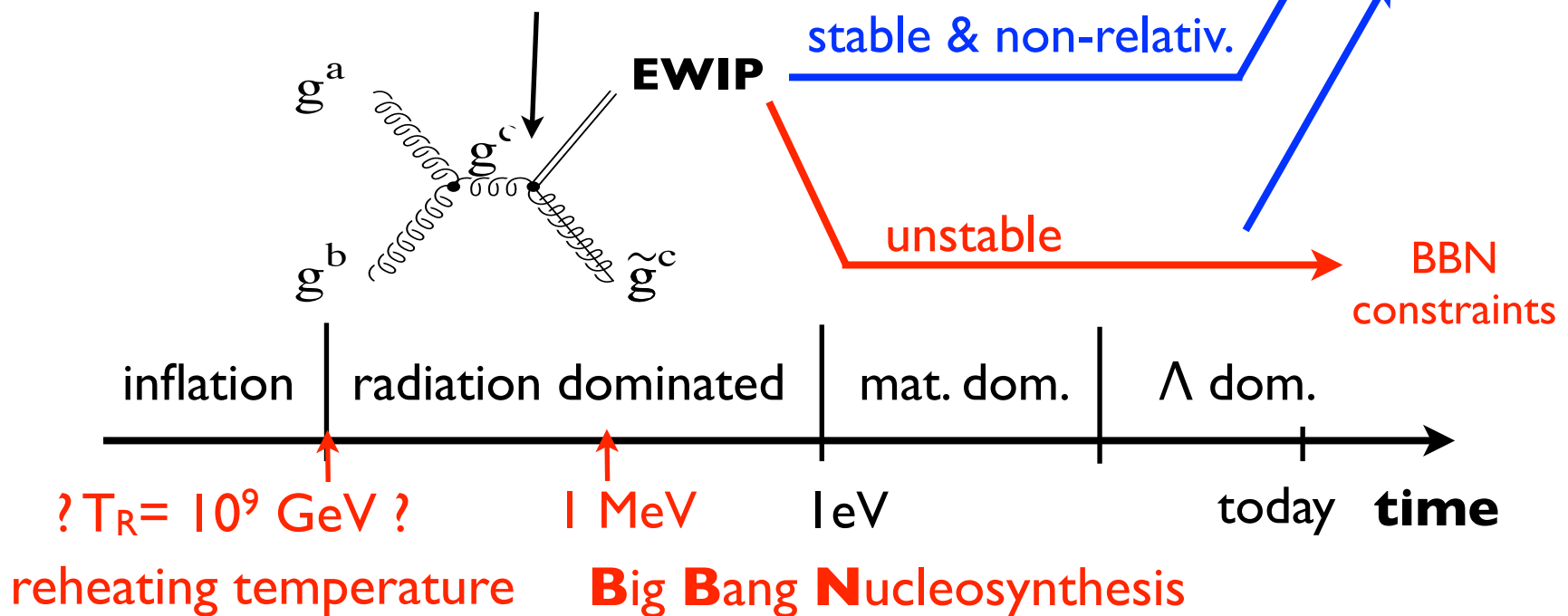
High Reheating Temperature Scenarios

Thermal Leptogenesis

requires $T_R > 10^9 \text{ GeV}$
 → baryon asymmetry

Thermal EWIP production

$$\propto (p/M_{\text{Pl}})^n \propto (T_R/M_{\text{Pl}})^n$$



Dark Radiation

- Radiation content of the Universe at BBN and later

$$\rho_{\text{rad}} = \left[1 + \frac{7}{8} (N_\nu + \Delta N_{\text{eff}}) \left(\frac{T}{T_\nu} \right)^4 \right] \rho_\gamma$$

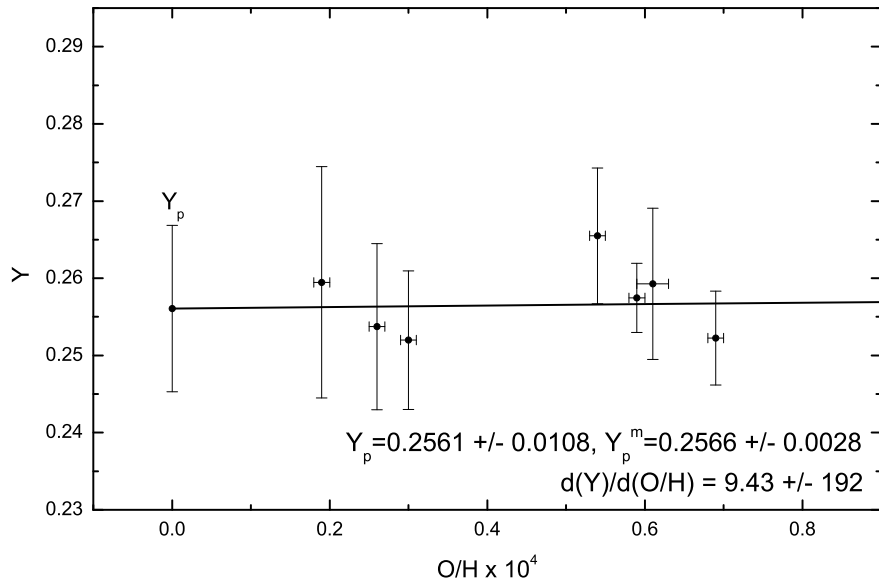
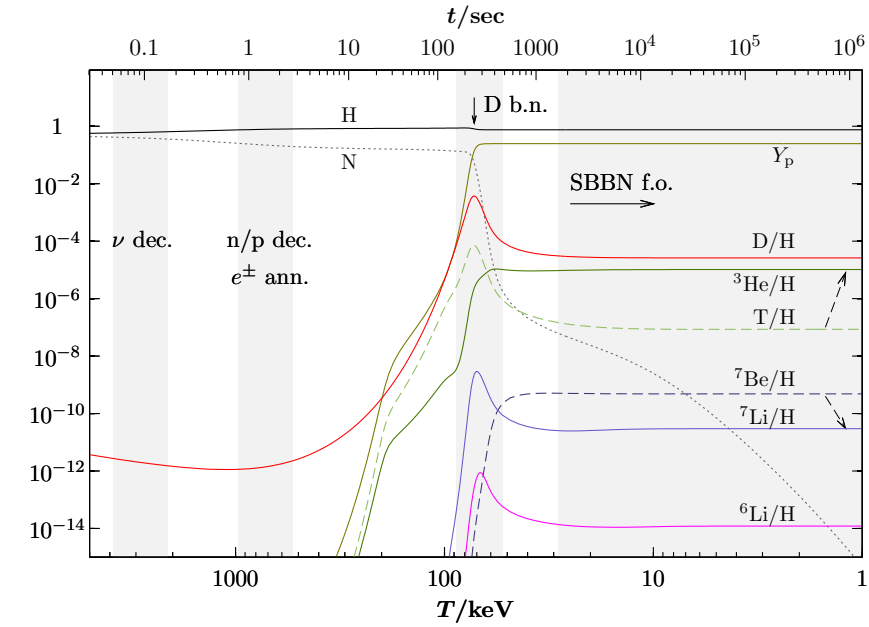
↑
number of neutrinos
↑
neutrino temperature
←
photon energy density

Parametrizes contribution of additional relativistic species

- More radiation → faster expansion → more efficient BBN of ^4He
- More radiation → later mat-rad eq → visible in CMB + LSS

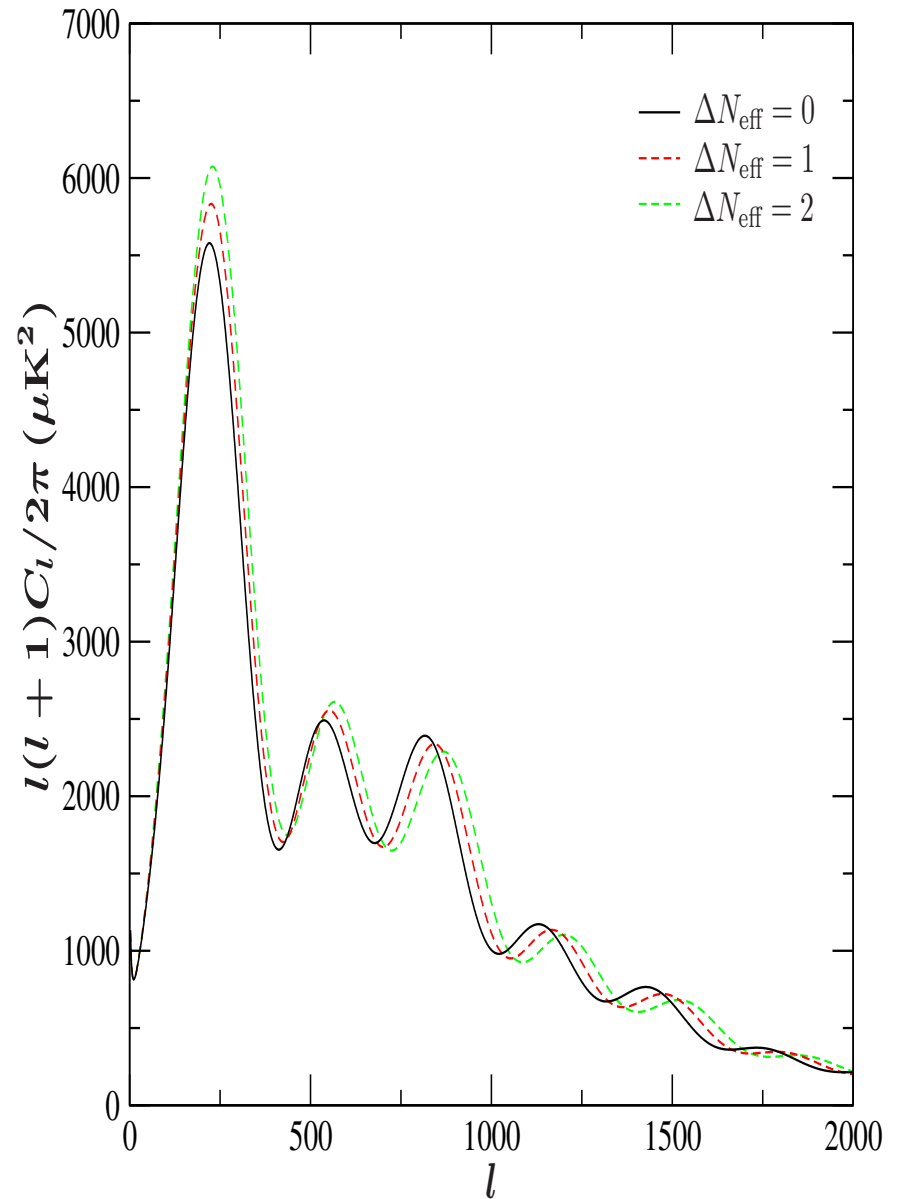
Data	p.m./mean	upper limit	
Y_p^{IT} [1] + $[\text{D}/\text{H}]_p$ [49]	0.76	< 1.97 (3σ)	↕ BBN
Y_p^{Av} [2] + $[\text{D}/\text{H}]_p$ [49]	0.77	< 3.53 (3σ)	
CMB + HPS + HST [6]	1.73	< 3.59 (2σ)	↕ CMB + LSS
Planck+WP+highL+BAO [8]	0.25	< 0.79 (2σ)	
Planck+WP+highL+ H_0 +BAO [8]	0.47	< 0.95 (2σ)	

BBN



Axions etc.

CMB + LSS



43

Steffen

Planck 2013 results XVI: Cosmological Parameters

Combining *Planck*, *WMAP* polarization and the high- ℓ experiments gives

$$N_{\text{eff}} = 3.36_{-0.64}^{+0.68} \quad (95\%; \text{Planck+WP+highL}). \quad (74)$$

The marginalized posterior distribution is given in Fig. 27 (black curve).

Increasing N_{eff} at fixed θ_* and z_{eq} necessarily raises the expansion rate at low redshifts too. Combining CMB with distance measurements can therefore improve constraints (see Fig. 27) although for the BAO observable $r_{\text{drag}}/D_V(z)$ the reduction in both r_{drag} and $D_V(z)$ with increasing N_{eff} partly cancel. With the BAO data of Sect. 5.2, the N_{eff} constraint is tightened to

$$N_{\text{eff}} = 3.30_{-0.51}^{+0.54} \quad (95\%; \text{Planck+WP+highL+BAO}). \quad (75)$$

Our constraints from CMB alone and CMB+BAO are compatible with the standard value $N_{\text{eff}} = 3.046$ at the 1σ level, giving no evidence for extra relativistic degrees of freedom.

Since N_{eff} is positively correlated with H_0 , the tension between the *Planck* data and direct measurements of H_0 in the base Λ CDM model (Sect. 5.3) can be reduced at the expense of high N_{eff} . The marginalized constraint is

$$N_{\text{eff}} = 3.62_{-0.48}^{+0.50} \quad (95\%; \text{Planck+WP+highL+H}_0). \quad (76)$$

For this data combination, the χ^2 for the best-fitting model allowing N_{eff} to vary is lower by 5.0 than for the base $N_{\text{eff}} = 3.046$ model. The H_0 fit is much better, with $\Delta\chi^2 = -4.0$, but there is no strong preference either way from the CMB. The low- ℓ temperature power spectrum does mildly favour the high N_{eff} model ($\Delta\chi^2 = -1.6$) since N_{eff} is positively correlated with n_s (see Fig. 24) and increasing n_s reduces power on large scales. The rest of the *Planck* power spectrum is agnostic ($\Delta\chi^2 = -0.5$), while the high- ℓ experiments mildly disfavour high N_{eff} in our fits ($\Delta\chi^2 = 1.3$). Further including the BAO data pulls the central value downwards by around 0.5σ (see Fig. 27):

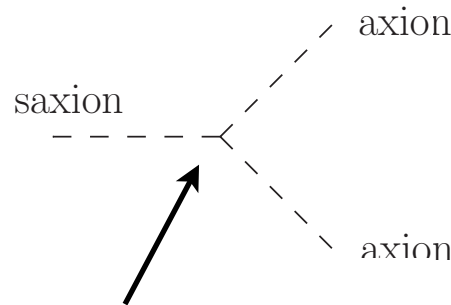
$$N_{\text{eff}} = 3.52_{-0.45}^{+0.48} \quad (95\%; \text{Planck+WP+highL+H}_0\text{+BAO}). \quad (77)$$

The χ^2 at the best-fit for this data combination ($N_{\text{eff}} = 3.37$) is lower by 3.6 than the best-fitting $N_{\text{eff}} = 3.046$ model. While the high N_{eff} best-fit is preferred by *Planck*+WP ($\Delta\chi^2 = -3.3$) and the H_0 data ($\Delta\chi^2 = -2.8$ giving an acceptable $\chi^2 = 2.4$ for this data point), it is disfavoured by the high- ℓ CMB data ($\Delta\chi^2 = 2.0$) and slightly by BAO ($\Delta\chi^2 = 0.4$). We conclude that the tension between direct H_0 measurements and the CMB and BAO data in the base Λ CDM can be relieved at the cost of additional neutrino-like physics, but there is no strong preference for this extension from the CMB damping tail.

$$\Delta N_{\text{eff}} = 3.62 + 0.5 - 3.046 = 1.074 @ 95\% \text{ CL}$$

Dark Radiation → SUSY + PQ

- Axions from decays of thermal saxions → extra radiation



$$\mathcal{L}_{\text{PQ}}^{\text{kin}} = \left(1 + \frac{\sqrt{2}x}{v_{\text{PQ}}} \sigma \right) \left[\frac{1}{2} \partial^\mu a \partial_\mu a + \frac{1}{2} \partial^\mu \sigma \partial_\mu \sigma + i \bar{a} \gamma^\mu \partial_\mu \tilde{a} \right] + \dots$$

$$x = \sum_i \frac{q_i^3 v_i^2}{v_{\text{PQ}}^2}, \quad v_{\text{PQ}} = \sqrt{\sum_i v_i^2 q_i^2}, \quad f_{\text{PQ}} = \sqrt{2} v_{\text{PQ}}$$

- $\Delta N_{\text{eff}} \propto \left(\frac{100 \text{ GeV}}{m_\sigma} \right)^{1/2} \left(\frac{f_{\text{PQ}}}{10^{11} \text{ GeV}} \right) \left(\frac{Y_\sigma^{\text{eq/TP}}}{10^{-3}} \right)$

[Graf, FDS, '12]

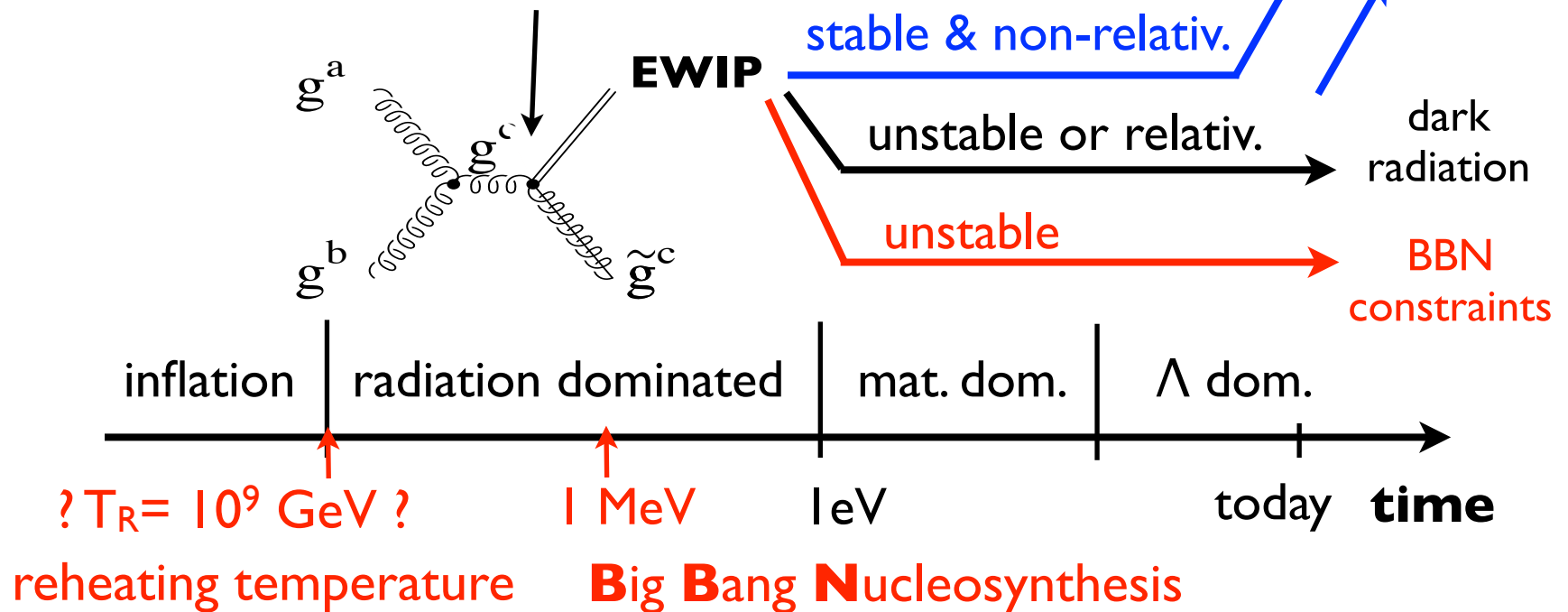
High Reheating Temperature Scenarios

Thermal Leptogenesis

requires $T_R > 10^9 \text{ GeV}$
 \rightarrow baryon asymmetry

Thermal EWIP production

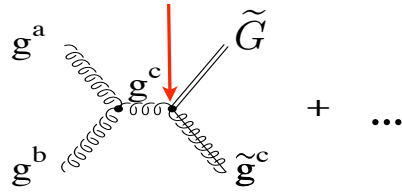
$$\propto (p/M_{\text{Pl}})^n \propto (T_R/M_{\text{Pl}})^n$$



Gravitino Dark Matter Scenario

- **Gravitino is the stable LSP (R-parity conservation is assumed)**
- **Axino is heavy & unstable (decays prior to LOSP decoupling)**
- **Saxion decays into axion dark radiation**
- **Axion contributes to dark radiation and dark matter**
- **Sneutrino NLSP case allows for thermal leptogenesis**

Thermal Axino/
Gravitino Production



Gravitino Dark Matter

inflation

radiation dominated

mat. dom.

Λ dom.

$$\rho_{\text{rad}} \propto a^{-4}$$

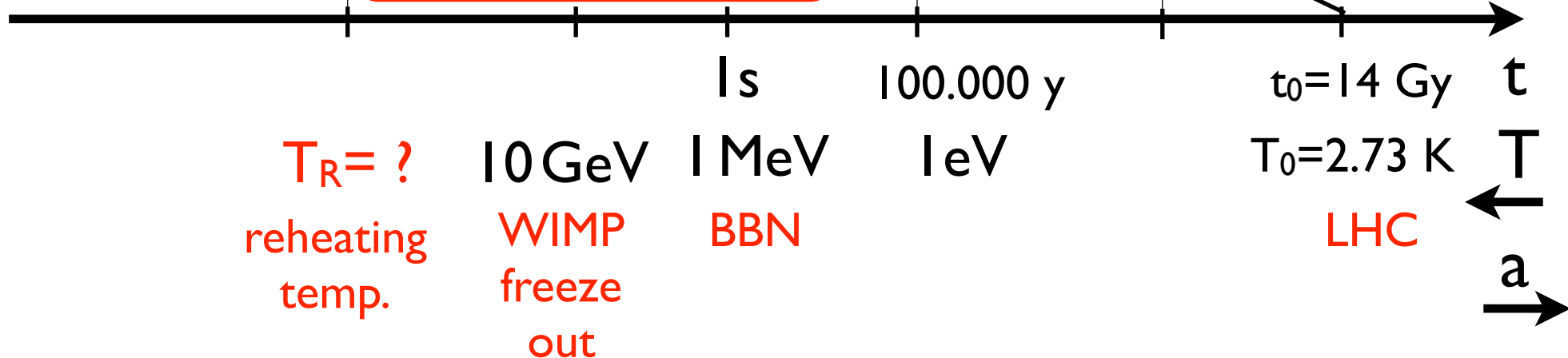
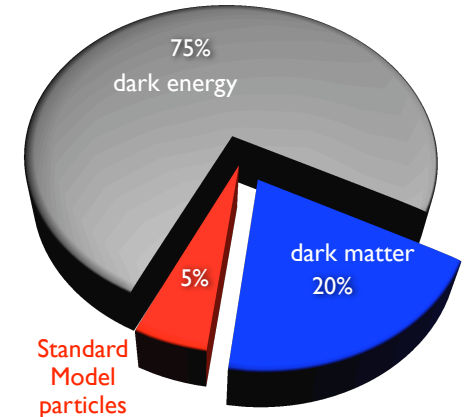
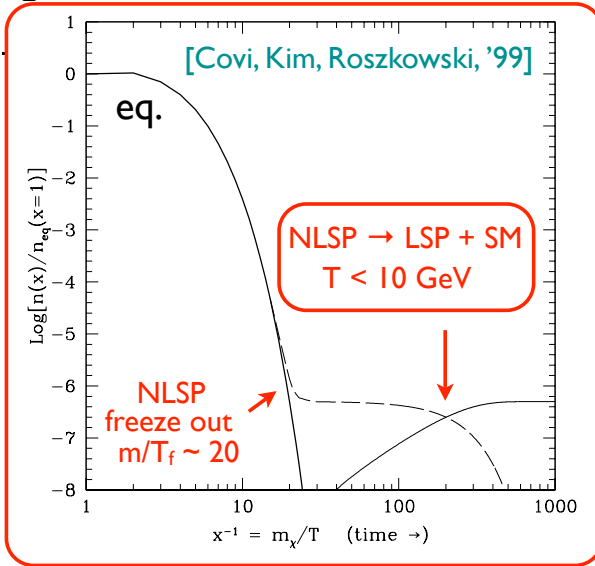
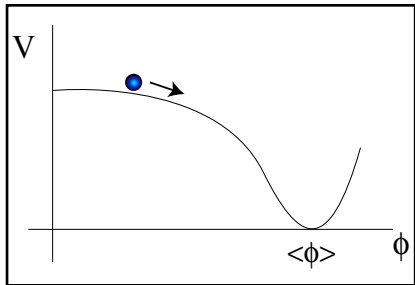
$$\rho_{\text{mat}} \propto a^{-3}$$

$$\rho_{\Lambda} \propto a^0$$

slow roll

reheat phase

$$\rho_{\phi} \propto a^0$$



$T_R = ?$
reheating
temp.

10 GeV
WIMP
freeze
out

1 MeV
BBN

1 eV

$T_0 = 2.73$ K
LHC

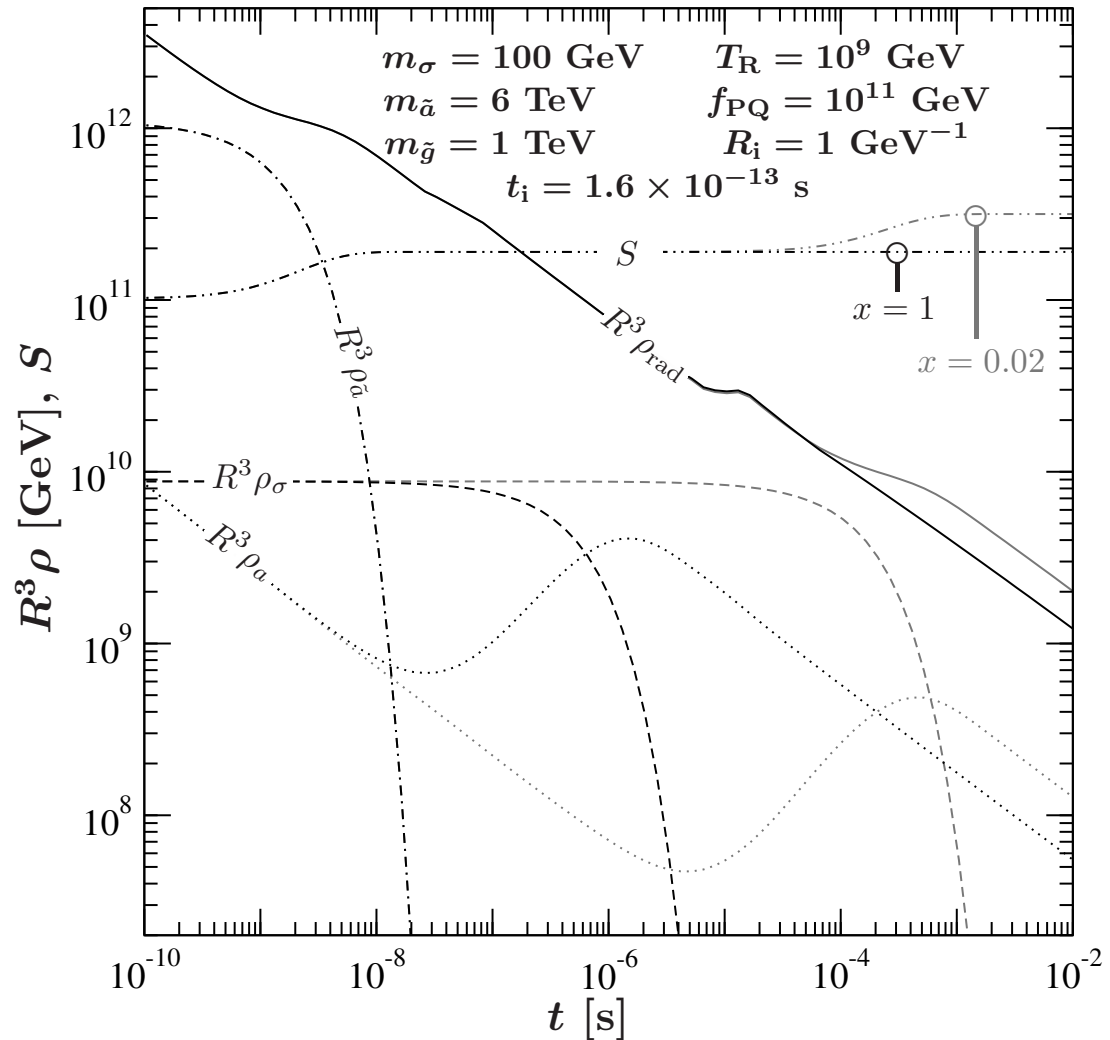
Gravitino Dark Matter Scenario

axino \rightarrow gluon + gluino

axino \rightarrow axion + gravitino

saxion \rightarrow 2 axions

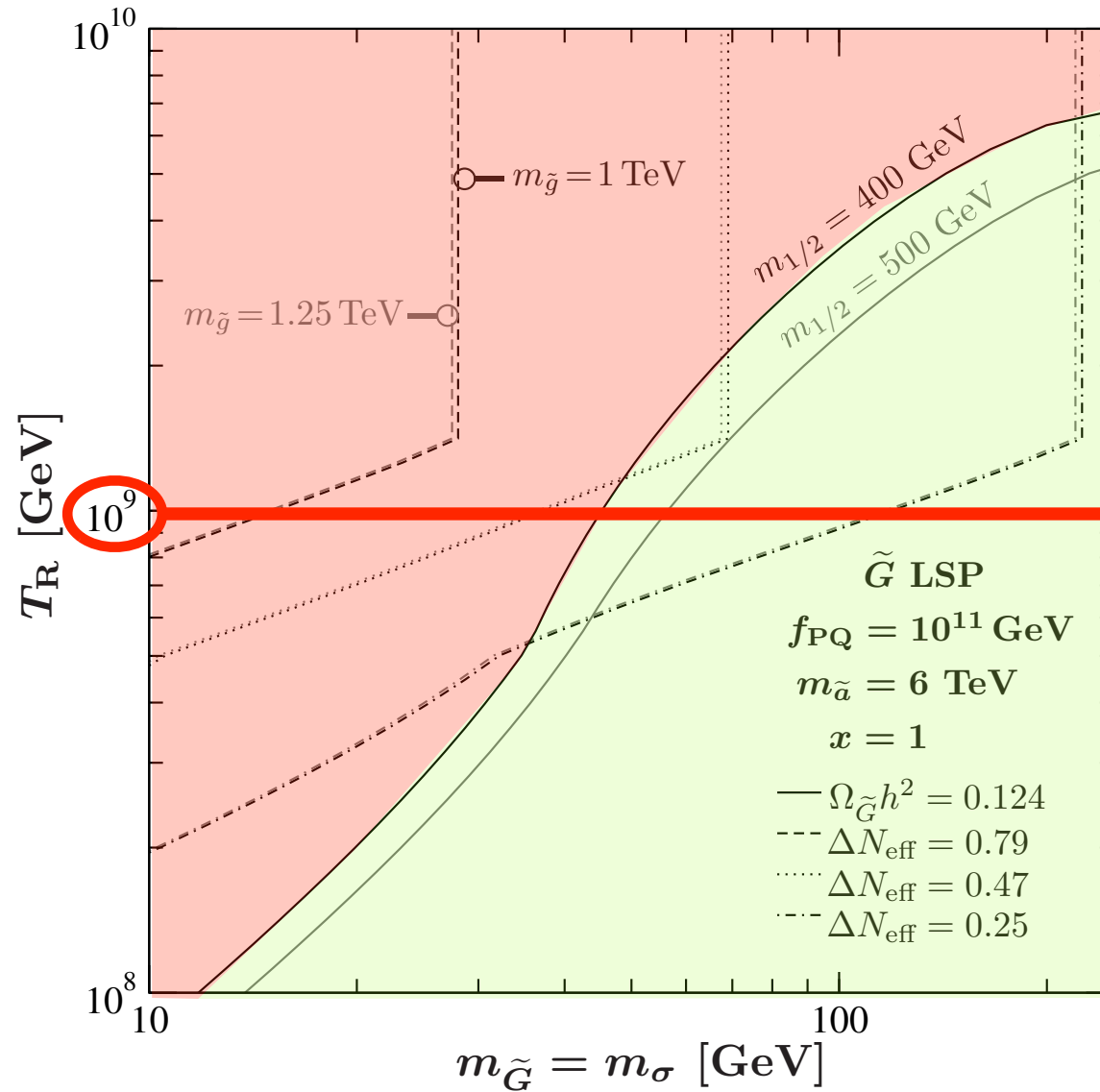
saxion \rightarrow 2 gluons



$$\mathcal{L}_{\tilde{G}\tilde{a}a} = \frac{1}{2} \left(\partial_\mu \sigma \bar{\Psi}_{M\nu} \gamma^\mu \gamma^\nu \tilde{a}_M + i \partial_\mu a \bar{\Psi}_{M\nu} \gamma^\mu \gamma^\nu \gamma^5 \tilde{a}_M \right)$$

[Graf, FDS, I302.2|43]

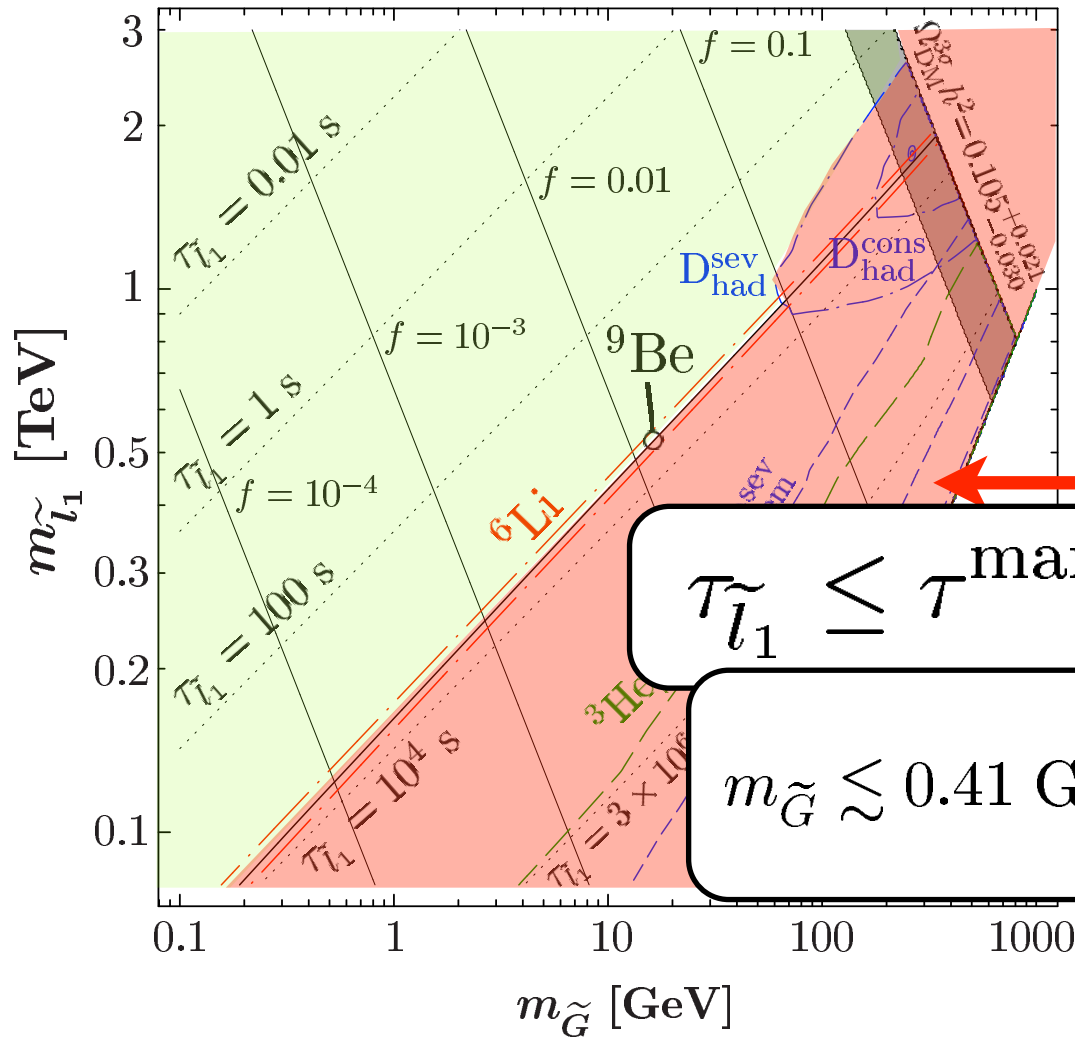
Gravitino Dark Matter Scenario



Gravitino LSP Case with a Charged Slepton NLSP

(C)BBN Constraints

[Pospelov, Pradler, FDS, '08]



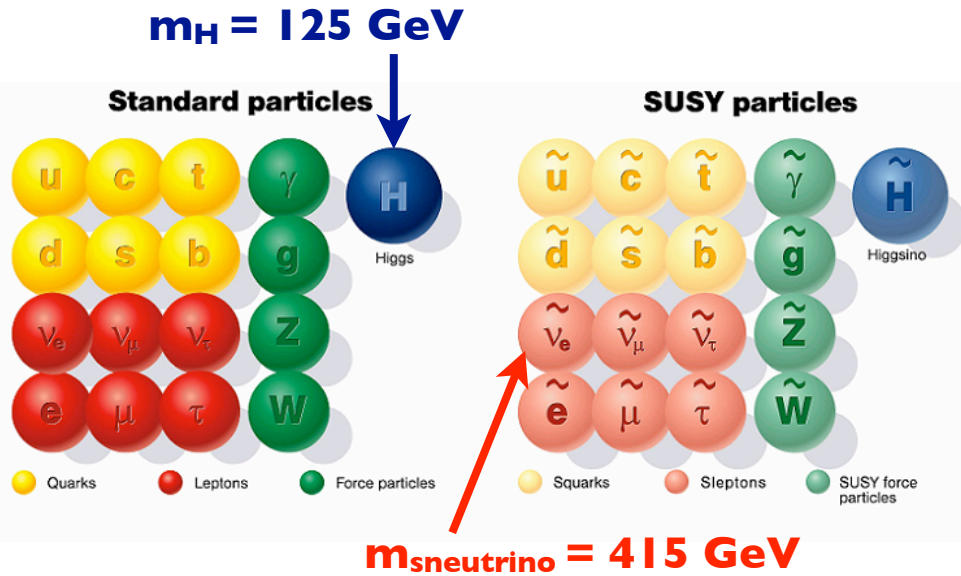
**disfavored
by
cosmological
constraints**

$$\tau_{\tilde{l}_1} \leq \tau^{\text{max}}$$

$$m_{\tilde{G}} \lesssim 0.41 \text{ GeV} \left(\frac{\tau_{\tilde{l}_1}}{10^4 \text{ s}} \right)^{\frac{1}{2}} \left(\frac{m_{\tilde{l}_1}}{100 \text{ GeV}} \right)^{\frac{5}{2}} \equiv m_{\tilde{G}}^{\text{max}}$$

Particle Physics

- 2016: Large Hadron Collider (14 TeV) **neutrino discovery** at ATLAS & CMS



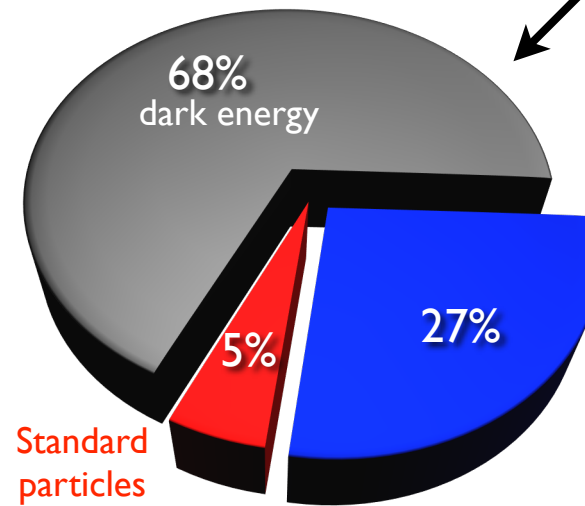
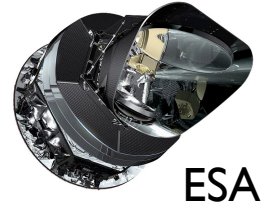
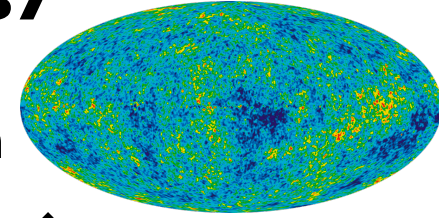
- Intrinsic fine tuning problems

- ✓ Hierarchy Problem ($m_H \ll M_{\text{Planck}}$)
supersymmetry
- ✓ Strong CP Problem ($\Theta_{\text{QCD}} \ll 1$)
Peccei-Quinn symmetry
- ✓ Small Neutrino Masses ($m_\nu \ll m_H$)
See-saw mechanism

Axions etc.

Cosmology

- 2013: Planck sky map of the CMB radiation



gravitino EWIP
dark matter



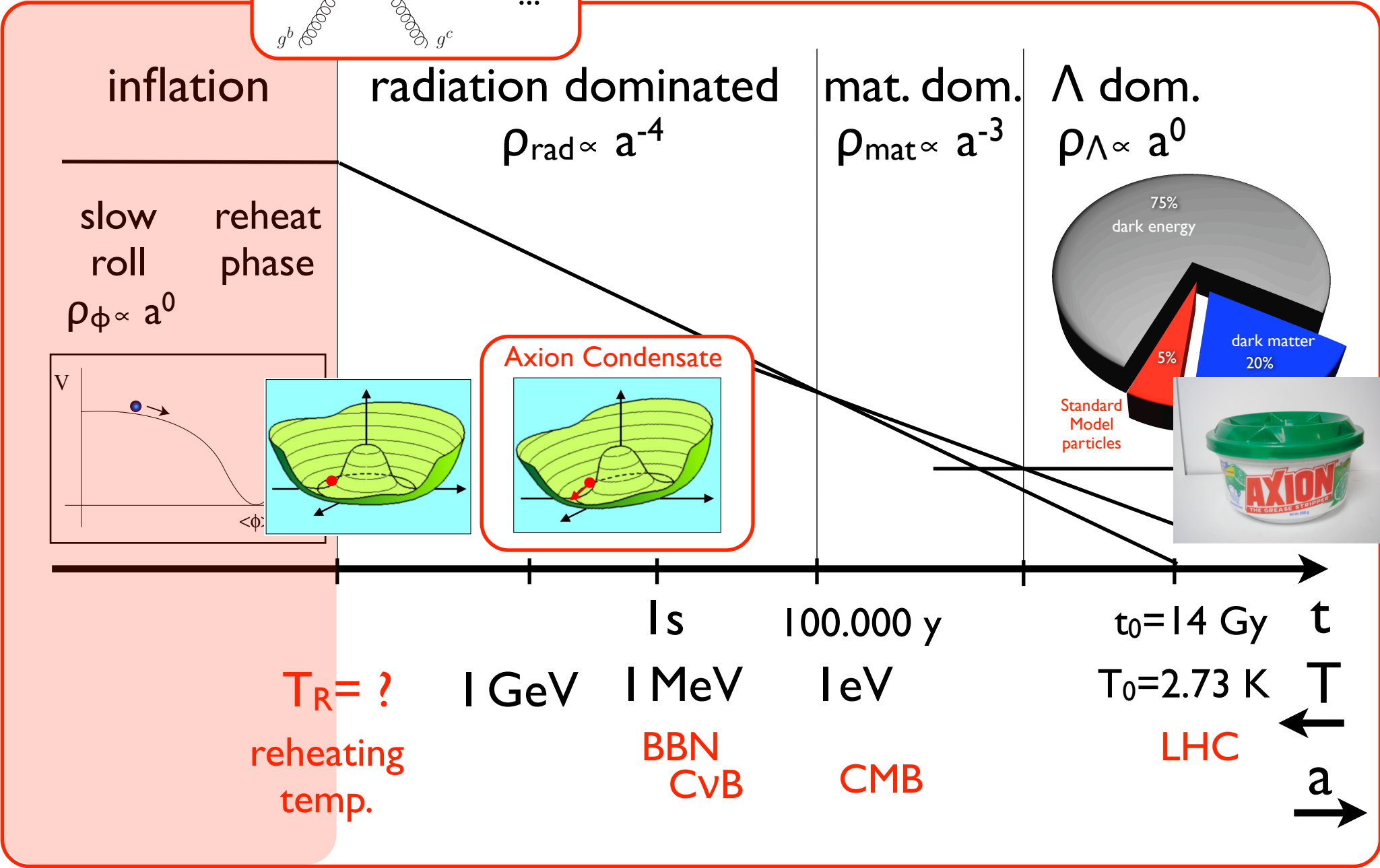
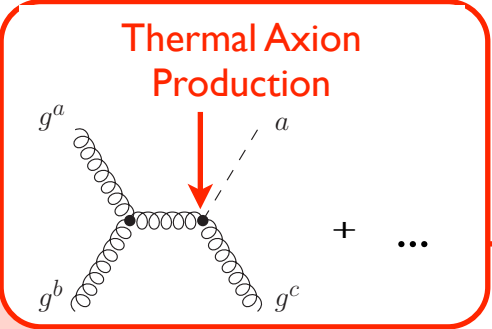
- Cosmological puzzles

- ✓ Matter-Antimatter Asymmetry
thermal leptogenesis
- ✓ Particle Identity & Origin of Dark Matter
thermally produced gravitinos
- ? Dark Energy = Cosmological Constant

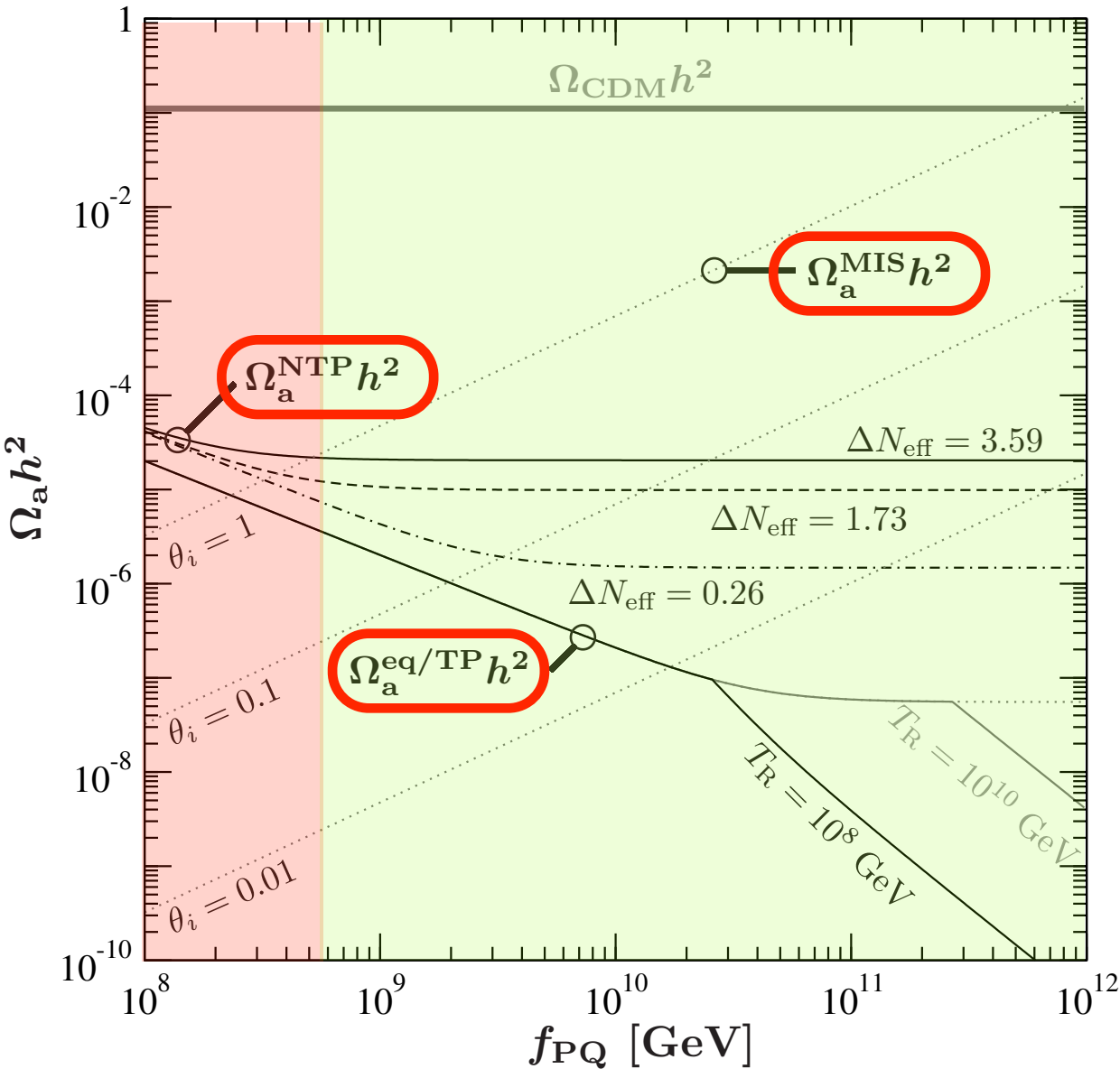
Axion Dark Matter Scenario

- **Axino is a very light stable LSP (R-parity conservation)**
- **Gravitino is the NLSP and decays into axion & axino dark radiation**
- **Saxion decays into axion dark radiation**
- **Axion contributes to dark radiation and dark matter**
- **Stau NLSP case & thermal leptogenesis is possible**

Axion Dark Matter



Lee-Weinberg Curve for Axions



Axion Condensate: CDM

$$\Omega_a^{\text{MIS}} h^2 \sim 0.15 \theta_i^2 (f_{\text{PQ}}/10^{12} \text{ GeV})^{7/6}$$

[... ; Sikivie, '08; Kim, Carosi, '08; ...]

\geq Three Axion Populations

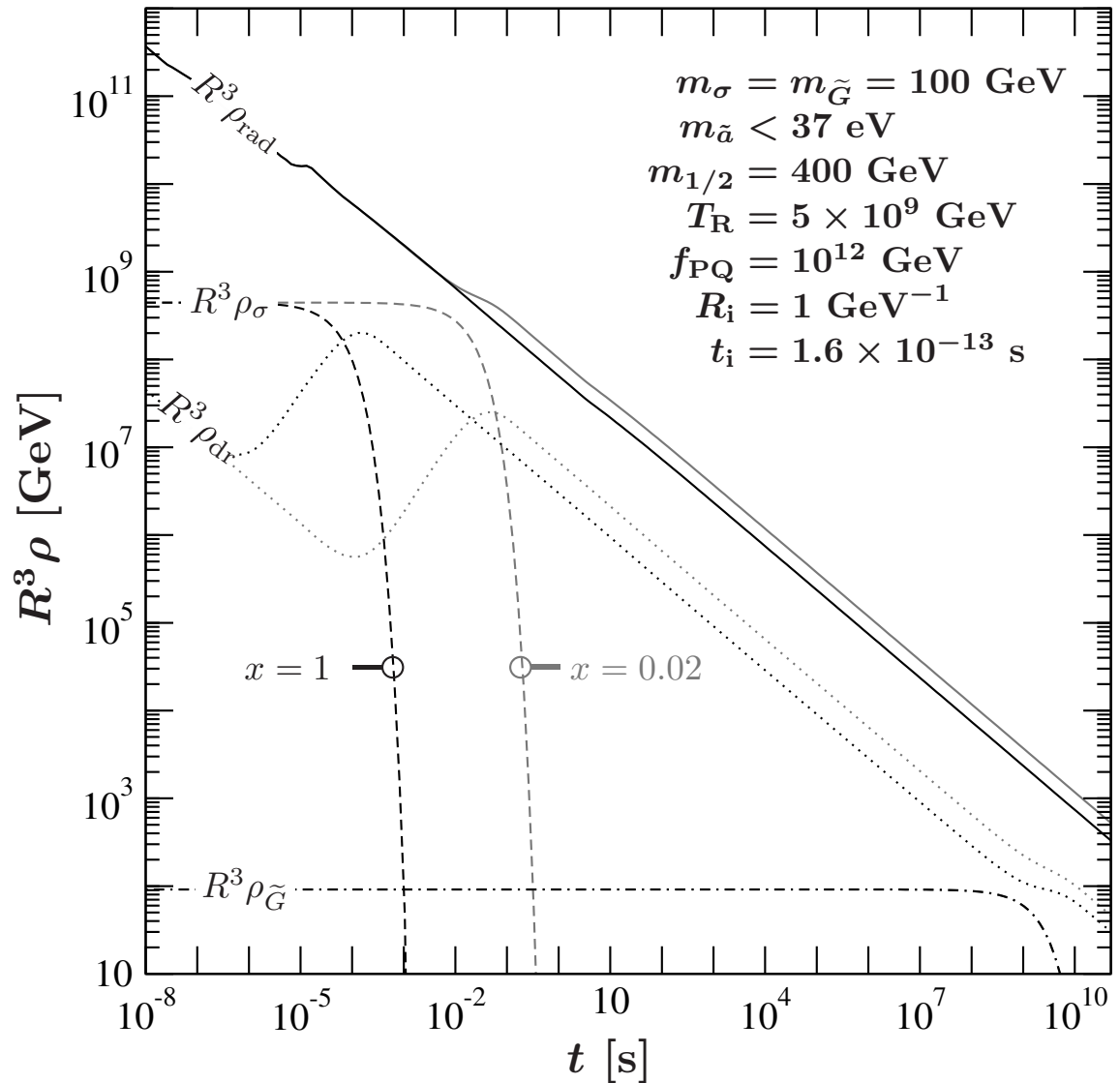
[Graf, Steffen, '12]

Axion Dark Matter Scenario

saxion \rightarrow 2 gluons

saxion \rightarrow 2 axions

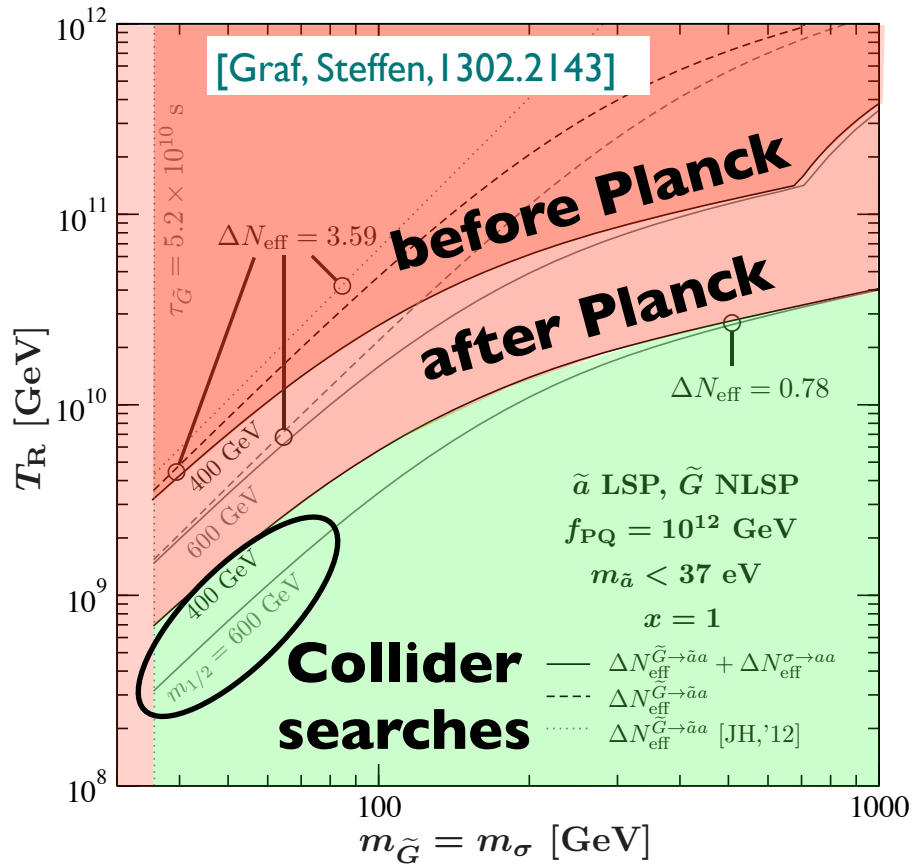
gravitino \rightarrow axion + axino



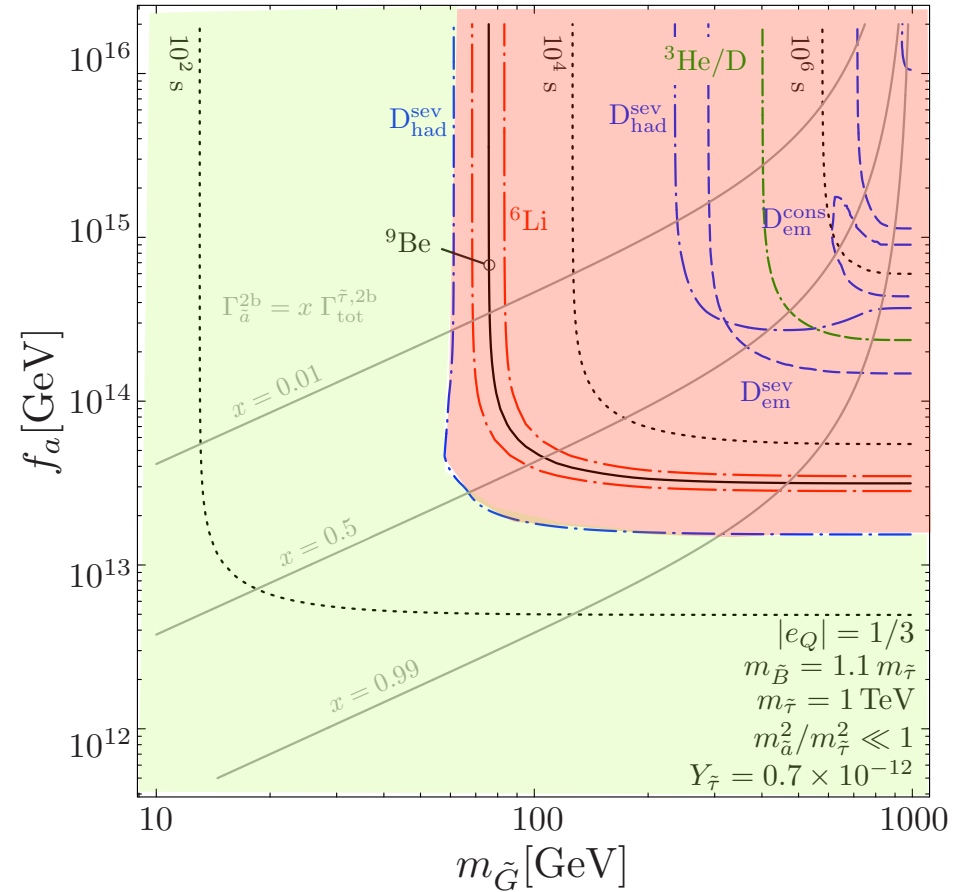
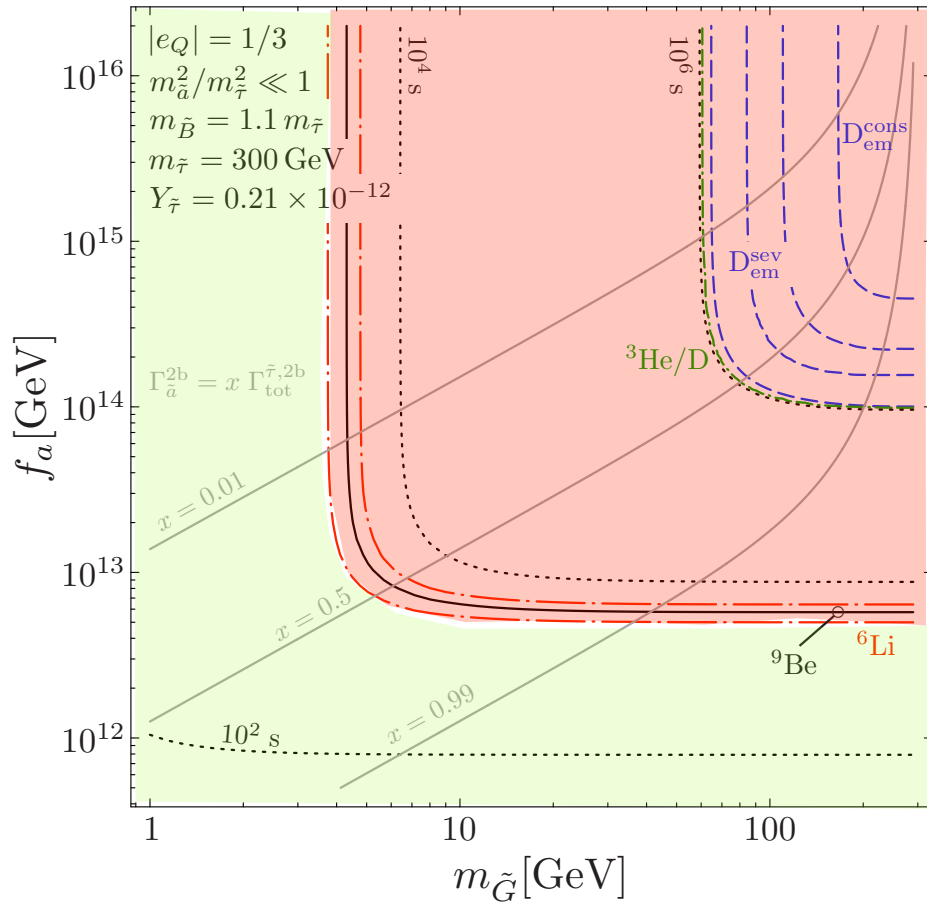
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[Graf, FDS, I302.2|43]

Axion Dark Matter Scenario

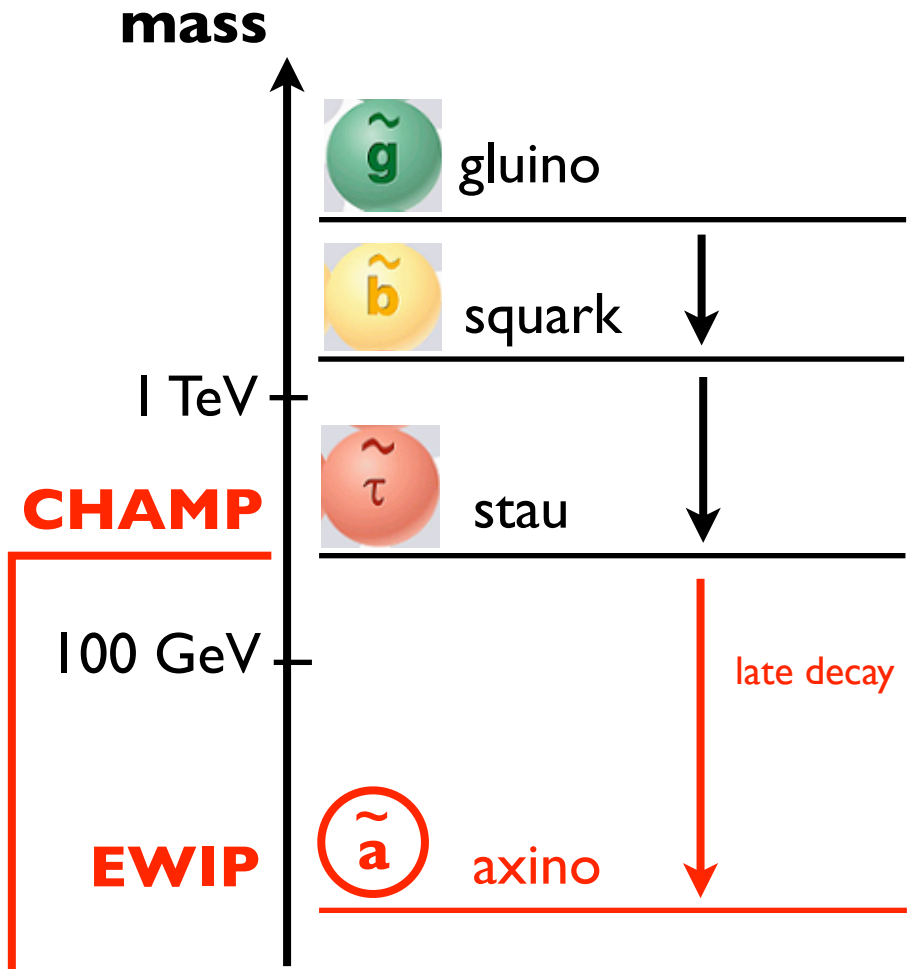
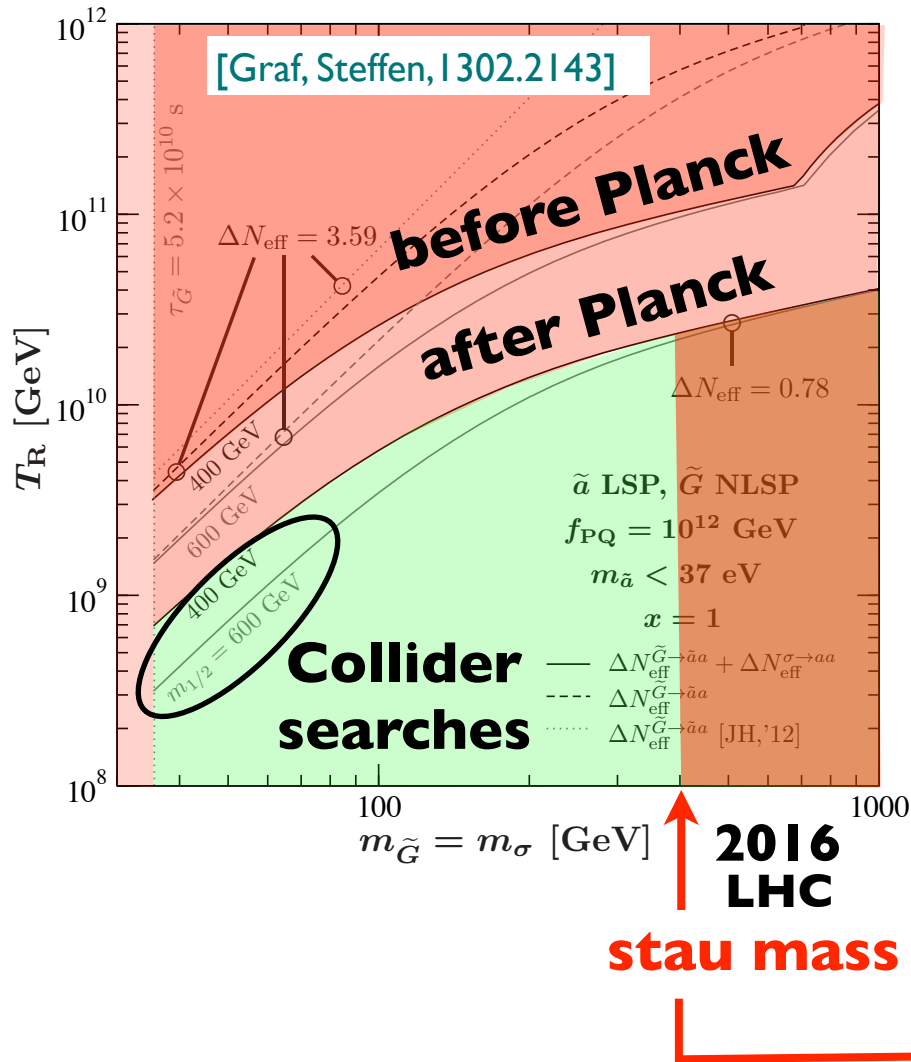


$m_{\text{maxino}}, m_{\text{gravitino}} < m_{\text{stau}}$

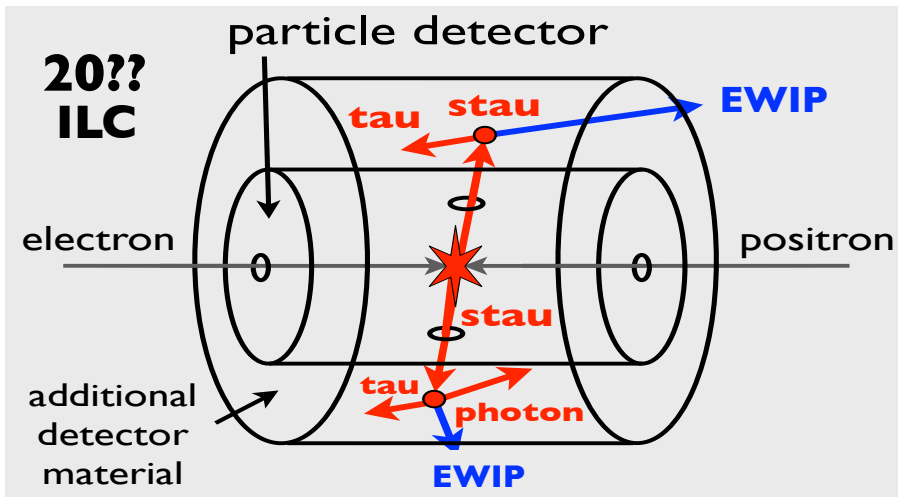
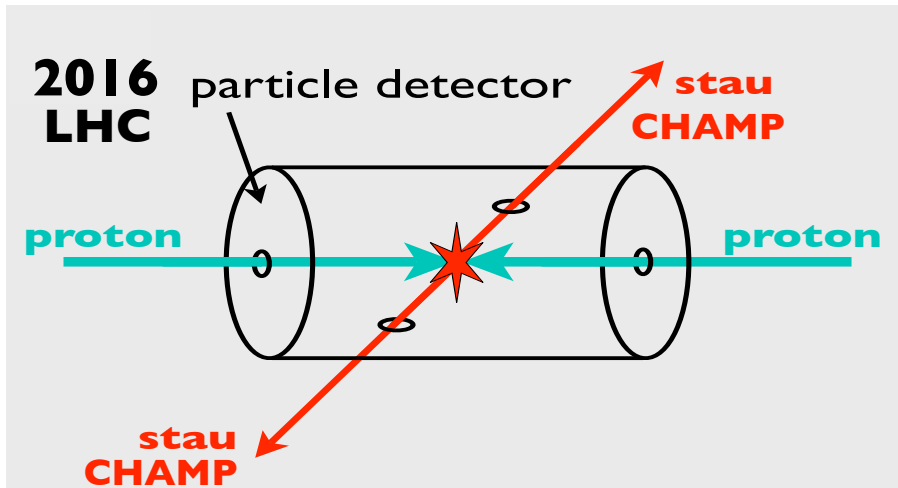


[Freitas, Tajuddin, FDS, Wyler, '09]

Axion Dark Matter Scenario



CHARGED Massive Particles (CHAMPs)



Key questions on CHAMP properties

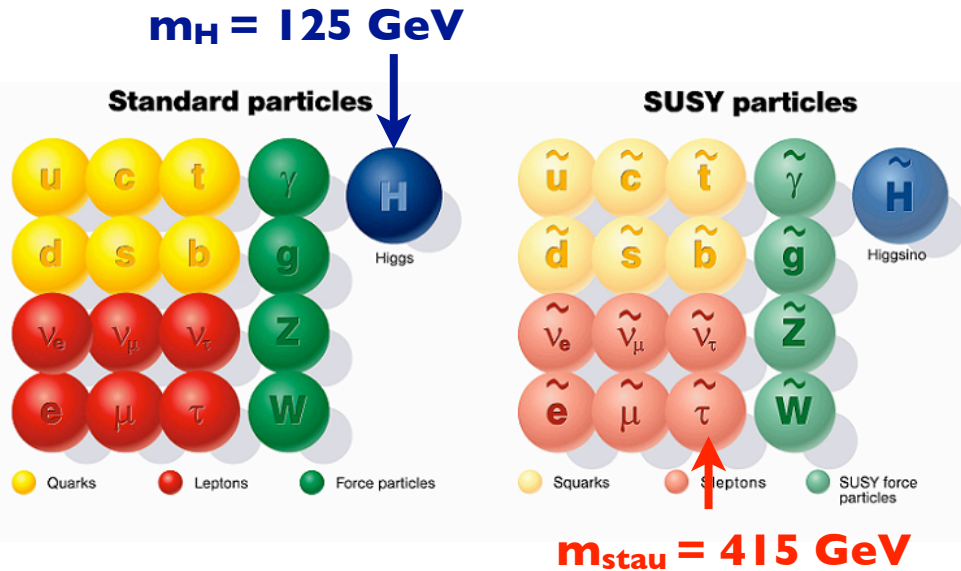
**Stable? Lifetime?
Decay products?**



New detector concepts
→ **stop/collect CHAMPs**
→ **study CHAMP decays**

Particle Physics

- 2016: Large Hadron Collider (14 TeV) **stau discovery** at ATLAS & CMS

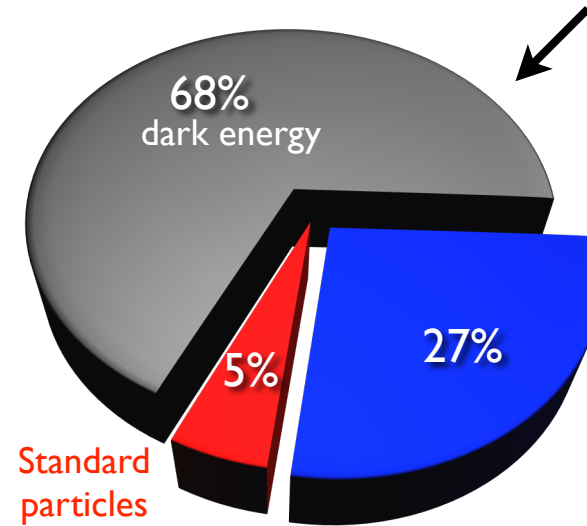
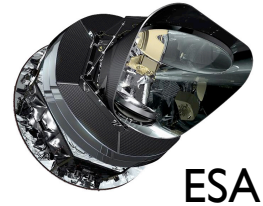
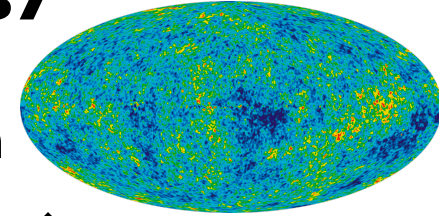


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axion EWIP
dark matter



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axion condensate
- ? Dark Energy = Cosmological Constant

Conclusions

- **Cosmological observations still call for new physics**
- **Hierarchy problem & strong CP problem → SUSY axion models**
- **EWIPs are a new well-motivated class of particles**
- **EWIP can explain dark matter and dark radiation (?)**
- **High-reheating temperature scenarios → thermal leptogenesis**
 - **Gravitino dark matter with sneutino LOSPs**
 - **Axion dark matter with stau LOSPs**
- **Various cosmological aspects & promising collider projects**