Axions in Astrophysics and Cosmology.

Andreas Ringwald

Annual Retreat of GRK and PRISMA Kloster Johannisberg, D 21 September 2017





> Axion field is born after PQ symmetry breaking, $T \lesssim T_c^{PQ} \sim v_{PQ} = N_{DW} f_A$

Unbroken Symmetry

Broken Symmetry



[Peking University]



> Axion field becomes massive near QCD phase transition, $T_c^{\text{QCD}} \sim 200 \,\text{MeV}$





> DM from vacuum realignment:

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,....]

- In early universe, axion frozen at random initial value
- Later, field feels pull of mass towards zero and oscillates around it





[Raffelt]



DM from vacuum realignment: >

[Preskill, Wise, Wilczek 83; Abbott, Sikivie 83; Dine, Fischler 83,....]

- In early universe, axion frozen at random initial value
- Later, field feels pull of mass towards zero and oscillates around it
- Spatially uniform oscillating classical field = coherent state of many, extremely non-relativistic particles = CDM



> DM from vacuum realignment:

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,....]

- In early universe, axion frozen at random initial value
- Later, field feels pull of mass towards zero and oscillates around it
- Spatially uniform oscillating classical field = coherent state of many, extremely non-relativistic particles = CDM
- Crucial QCD input for prediction of axion DM abundance:
 - Equation of state at temperatures around 1 GeV: determines H(T)
 - Topological susceptibility:

 $\chi(T)\equiv \int d^4x \langle q(x)q(0)\rangle_T$ determines $m^2_A(T)=\chi(T)/f^2_A$



> DM from vacuum realignment:

[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83,....]

- In early universe, axion frozen at random initial value
- Later, field feels pull of mass towards zero and oscillates around it
- Spatially uniform oscillating classical field = coherent state of many, extremely non-relativistic particles = CDM
- Crucial QCD input for prediction of axion DM abundance:
 - Equation of state at temperatures around 1 GeV: determines H(T)
 - Topological susceptibility:

 $\chi(T)\equiv \int d^4x \langle q(x)q(0)\rangle_T$ determines $m^2_A(T)=\chi(T)/f^2_A$





- If PQ symmetry broken during inflation and not restored afterwards (pre-inflationary PQ breaking scenario)
 - Axion CDM density depends on single initial angle during inflation and f_A

f_A [GeV]





Pre-inflationary PQ symmetry breaking scenario

- If PQ symmetry broken during inflation and not restored afterwards (pre-inflationary PQ breaking scenario)
 - Axion CDM density depends on single initial angle during inflation and *f_A*
 - Axion is present during inflation and creates isocurvature fluctuations which are not erased after inflation



[Wilczek,Turner `91; Beltran et al. 06; Hertzberg,Tegmark,Wilczek 08; Visinelli,Gondolo 09; Hamann et al. 09; **Wantz,Shellard 09**]



- If Peccei-Quinn symmetry restored after inflation (post-inflationary PQ breaking scenario)
 - Vacuum realignment contribution depends on spatially averaged initial misalignment angle and f_A

$$\Omega_A^{\rm real} h^2 \approx 0.12 \, \left(\frac{30 \, \mu {\rm eV}}{m_A}\right)^{1.165}$$





[Saikawa]

- If Peccei-Quinn symmetry restored after inflation (post-inflationary PQ breaking scenario)
 - Vacuum realignment contribution depends on spatially averaged initial misalignment angle and *f_A*

$$\Omega_A^{\rm real} h^2 \approx 0.12 \, \left(\frac{30 \, \mu {\rm eV}}{m_A}\right)^{1.165}$$

 Additional contributions arise from decay of topological defects (axion strings and axion domain walls)





[DAMTP]



- If Peccei-Quinn symmetry restored after inflation (post-inflationary PQ breaking scenario)
 - Vacuum realignment contribution depends on spatially averaged initial misalignment angle and *f_A*

$$\Omega_A^{\rm real} h^2 \approx 0.12 \, \left(\frac{30 \, \mu {\rm eV}}{m_A}\right)^{1.165}$$

- Additional contributions arise from decay of topological defects (axion strings and axion domain walls)
- Latter determined by simulations of cosmic string and wall networks
 [Hiramatsu et al. 12; Kawasaki et al. 15]
- May explain dark matter for

 $30\,\mu\mathrm{eV} \lesssim m_A \lesssim 10\,\mathrm{meV}$





[Hiramatsu et al. 12]



Axion Dark Matter Direct Detection Experiments

Upcoming generation of axion dark matter direct detection experiments can probe entire mass range:



> $|\sigma| = \rho/\sqrt{2}$ or mixture with Higgs modulus may play role of inflaton, if it has non-minimal coupling to gravity,

$$S \supset -\int d^4x \sqrt{-g} \left[\frac{M^2}{2} + \xi_\sigma \,\sigma^*\sigma\right] R$$

[Fairbairn,Hogan,Marsh `14]



[Ballesteros, Redondo, AR, Tamarit `16]



 $> |\sigma| = \rho/\sqrt{2}$ or mixture with Higgs 10^{0} modulus may play role of inflaton, Planck + WP + BACif it has non-minimal coupling to 0.25 N = 50gravity, 0.20 $S \supset -\int d^4x \sqrt{-g} \left[\frac{M^2}{2} + \xi_\sigma \,\sigma^*\sigma\right] R$ 10^{-1} N = 600.15 ٤ [Fairbairn, Hogan, Marsh `14] 0.1010⁻² CMB observables $A_s = (2.20 \pm 0.08) \times 10^{-9}$, 0.05 $n_s = 0.967 \pm 0.004$, '10⁻³ 0.00 r < 0.070.97 0.98 0.95 0.96 0.99 0.94 fit by n_{s} [Fairbairn, Hogan, Marsh `14] $\xi \simeq 2 \times 10^5 \sqrt{\lambda} \gtrsim 10^{-3}$



PQ symmetry restored after inflation already in preheating stage when PQ field undergoes Hubble damped oscillations in quartic potential

[Ballesteros, Redondo, AR, Tamarit `16]







PQ symmetry restored after inflation already in preheating stage when PQ field undergoes Hubble damped oscillations in quartic potential

[Ballesteros, Redondo, AR, Tamarit `16]

CDM abundance fixes decay constant (mass) of axion





Unifying Inflation, Dark Matter, and Seesaw with PQ Field

- > Augmenting axion models with three SM singlet neutrinos, getting their Majorana masses also through the vev $v_{\sigma} = N_{\rm DW} f_A$
 - no strong CP problem
 - dark matter
 - inflation
 - neutrino masses and mixing
 - baryogenesis via leptogenesis

[Dias et al. `14; Ballesteros et al. `16]





Unifying Inflation, Dark Matter, and Seesaw with PQ Field

- > Augmenting axion models with three SM singlet neutrinos, getting their Majorana masses also through the vev $v_{\sigma} = N_{\rm DW} f_A$
 - no strong CP problem
 - dark matter
 - inflation
 - neutrino masses and mixing
 - baryogenesis via leptogenesis

[Dias et al. `14; Ballesteros et al. `16]

Complete and consistent history of the universe from inflation to now



[desy.de]



Evolution of stars (Main Sequence – Red-Giant (RG) – Helium Burning (HB) – White Dwarf (WD)) sensitive to additional energy losses



[Copyright Addison Wesley]





- RG cooling rate: Brightness of tip of RG branch in color-magnitude diagram of globular cluster [Viaux et al. 13]
- HB cooling rate: Number of HB stars vs. number of RGs in colormagnitude diagram of globular cluster [Ayala et al. 14]



[Giannotti `16]



- RG cooling rate: Brightness of tip of RG branch in color-magnitude diagram of globular cluster [Viaux et al. 13]
- HB cooling rate: Number of HB stars vs. number of RGs in colormagnitude diagram of globular cluster [Ayala et al. 14]





[Giannotti 15]



- RG cooling rate: Brightness of tip of RG branch in color-magnitude diagram of globular cluster [Viaux et al. 13]
- HB cooling rate: Number of HB stars vs. number of RGs in colormagnitude diagram of globular cluster [Ayala et al. 14]
- > WD cooling rate:
 - Period decrease of variable
 WDs [Kepler et al. 91,...]











- RG cooling rate: Brightness of tip of RG branch in color-magnitude diagram of globular cluster [Viaux et al. 13]
- HB cooling rate: Number of HB stars vs. number of RGs in colormagnitude diagram of globular cluster [Ayala et al. 14]
- > WD cooling rate:
 - Period decrease of variable
 WDs [Kepler et al. 91,...]
 - White dwarf luminosity function (WDLF) [Isern et al. 08-12]





RG cooling rate: Brightness of tip of RG branch in color-magnitude diagram of globular cluster [Viaux et al. 13]

- Period decrease of variable
 WDs [Kepler et al. 91,...]
- White dwarf luminosity function (WDLF) [Isern et al. 08-12]

$$g_{ae} \equiv \frac{m_e}{f_a} |C_{ae}| < 4.3 \times 10^{-13}$$

$$g_{a\gamma} \equiv \frac{\alpha}{2\pi f_a} |C_{a\gamma}| < 6.6 \times 10^{-11} \,\mathrm{GeV}^{-1}$$

$$g_{ae} \equiv \frac{m_e}{f_a} |C_{ae}| < 2.8 \times 10^{-13}$$



However, practically every stellar systems seems to be cooling a bit faster than predicted by models based on SM:



[Giannotti, Irastorza, Redondo, AR (2015); Giannotti, Irastorza, Redondo, AR 17]



Excessive energy losses of HBs, RG, WDs can be explained at one stroke by production of axion/ALP with coupling to photons and electrons:



$$g_{a\gamma} = C_{a\gamma} \alpha / (2\pi f_a)$$

 $\gamma + Ze \rightarrow Ze + a$
 $g_{ai} = C_{ai} m_i / f_a$
 $e + Ze \rightarrow Ze + e + a$



Excessive energy losses of HBs, RG, WDs can be explained at one stroke by production of axion/ALP with coupling to photons and electrons:





- > Axion/ALP so weakly interacting, that it passes any barrier
- Light-shining-through a wall:



> Oscillation probability:

$$P(a \leftrightarrow \gamma) = 4 \frac{(g_{a\gamma}\omega B)^2}{m_a^4} \sin^2\left(\frac{m_a^2}{4\omega}L_B\right)$$



> Any Light Particle Search (ALPS) at DESY (in coll. with AEI, UHH)







Currently best limits from LSW: ALPS (DESY) and OSQAR (CERN)





ALPS II at DESY (in coll. with AEI, UFL, U Mainz)



> ALPS II

- First pure laboratory experiment to surpass stellar bounds
- Can probe part of parameter space relevant for astro hints





> Appreciable flux of solar ALPs produced by Primakoff process in core:



Even higher flux of solar ALPs produced by Atomic Recombination, Bremsstrahlung and Compton (ABC) processes:







[Irastorza `16]



- Most sensitive until now: CERN Axion Solar Telescope (CAST)
 - Superconducting LHC dipole magnet
 - X-ray detectors
 - Use of buffer gas to extend sensitivity to higher masses (axion band)





Proposed successor: International Axion Observatory (IAXO)

- Dedicated superconducting toroidal magnet with much bigger aperture than CAST
- Extensive use of X-ray optics
- Low background X-ray detectors





[Armengaud et al (IAXO CDR) 1401.3233]



> IAXO

- Covers most of parameter space relevant for astro hints
- Able to probe meV mass axion





Neutron star in Cas A:



- Measured surface temperature reveals unusually fast cooling rate
- Hint on extra cooling by axion/ALP due to nucleon bremsstrahlung

$$N+N \to N+N+a$$

Required coupling to neutron:

$$g_{an}^2 \sim (1.4 \pm 0.5) \times 10^{-19}$$



[Leinson 14]



> SN 1987A:





 Emission of axion/ALPs would take away energy from neutrino burst and shorten it

 $g_{ap}^2 + g_{an}^2 < 3.6 \times 10^{-19}$

[Raffelt 08; Fischer et al. 16; Giannotti et al. (in prep.)]

 Not very solid bound: sparse data and interaction difficult to model



Excessive energy losses of HBs, RG, WDs, NS can be explained at one stroke by production of axion with coupling to photons, electrons, nucleons, e.g. for DFSZ axion model:



$$f_a = rac{v_{
m PQ}}{6},$$

$$C_{ae}^{
m DFSZ \ I} = rac{1}{3} \sin^2 eta \,, \qquad C_{ae}^{
m DFSZ \ II} = rac{1}{3} (1 - \sin^2 eta)$$

$$\begin{split} C_{a\gamma}^{\rm DFSZ\ I} &= \frac{8}{3} - 1.92(4)\,, \qquad C_{a\gamma}^{\rm DFSZ\ II} = \frac{2}{3} - 1.92(4)\,, \\ C_{Ap} &= -0.435\sin^2\beta + (-0.182\pm 0.025)\,\,, \\ C_{An} &= 0.414\sin^2\beta + (-0.16\pm 0.025)\,\,. \end{split}$$





[Giannotti, Irastorza, Redondo, AR, Saikawa 17]

Excessive energy losses of HBs, RG, WDs, NS can be explained at one stroke by production of axion with coupling to photons, electrons, nucleons, e.g. for DFSZ axion model:



$$f_a=rac{v_{
m PQ}}{6},$$

$$C_{ae}^{\mathrm{DFSZ\ I}} = \frac{1}{3} \sin^2 \beta, \qquad C_{ae}^{\mathrm{DFSZ\ II}} = \frac{1}{3} (1 - \sin^2 \beta)$$

$$\begin{split} C_{a\gamma}^{\rm DFSZ \ I} &= \frac{8}{3} - 1.92(4) \,, \qquad C_{a\gamma}^{\rm DFSZ \ II} = \frac{2}{3} - 1.92(4) \,, \\ C_{Ap} &= -0.435 \sin^2\beta + (-0.182 \pm 0.025) \,\,, \\ C_{An} &= 0.414 \sin^2\beta + (-0.16 \pm 0.025) \,\,. \end{split}$$





[Giannotti,Irastorza,Redondo,AR,Saikawa 17]

Excessive energy losses of HBs, RG, WDs, NS can be explained at one stroke by production of axion with coupling to photons, electrons, nucleons, e.g. for KSVZ axion/majoron model (as in SMASH):



 $\mathcal{L}_Y \supset -\overline{L}_i F_{ij} N_j H - \frac{1}{2} \overline{N}_i^c Y_{ij} N_j \sigma + \text{h.c.}$ $C_{ae} \simeq \frac{1}{8\pi^2 N} \left(\kappa_{ee} - \frac{1}{2} \operatorname{tr} \kappa \right)$ $C_{aq} \simeq \frac{1}{8\pi^2 N} T_3^q \operatorname{tr} \kappa$ $\kappa \equiv \frac{m_D m_D^{\dagger}}{\kappa^2} = \frac{F F^{\dagger}}{2}$



meV mass DFSZ axion or KSVZ axion/majoron could at the same time explain stellar energy losses and dark matter!

Model	$N_{\rm DW}$	Global fit includes	$f_a[10^8{ m GeV}]$	m_a [meV]	g	Upper limit on Δ
DFSZ I	6	HB,RGB,WD	0.77	74	0.033-1.8	1.1 - 1.5
	6	HB,RGB,WD,NS,SN	9.9	5.8	$(0.72-6.2) \times 10^{-8}$	$(0.26-2.4) \times 10^{-2}$
DFSZ II	6	HB,RGB,WD	1.2	47	$(0.23-5.7) \times 10^{-2}$	0.52 - 1.5
	6	HB,RGB,WD,NS,SN	9.1	6.3	$(0.12-1.1) \times 10^{-7}$	$(0.32-2.9) \times 10^{-2}$
	2	HB,RGB,WD	1.2	46	$(0.031-2.8) \times 10^2$	0.74 - 1.6
	4		1.2	46	0.064 - 4.1	0.21 - 2.3
KSVZ A/J	6		0.86	66	$(0.17-7.1) \times 10^{-1}$	0.93 - 1.5
$R_Q = (3, 1, -\frac{1}{3})$	2	HB,RGB,WD,NS,SN	4.1	14	$(0.19-7.6) \times 10^{-2}$	$(0.39-9.9) \times 10^{-1}$
	4		6.1	9.2	$(0.049-1.0) \times 10^{-4}$	$(0.31-6.6) \times 10^{-2}$
	6		8.0	7.1	$(0.26-2.5) \times 10^{-7}$	$(0.44-4.2) \times 10^{-2}$

[Giannotti, Irastorza, Redondo, AR, Saikawa 17]



Direct axion dark matter detection in this mass range challenging!



Summary

> Axion/ALPs may explain puzzles from astrophysics and cosmology:

- Dark matter
- Stellar energy losses
- Relevant parameter range will be probed decisively by next generation experiments
- > Stay tuned!

